

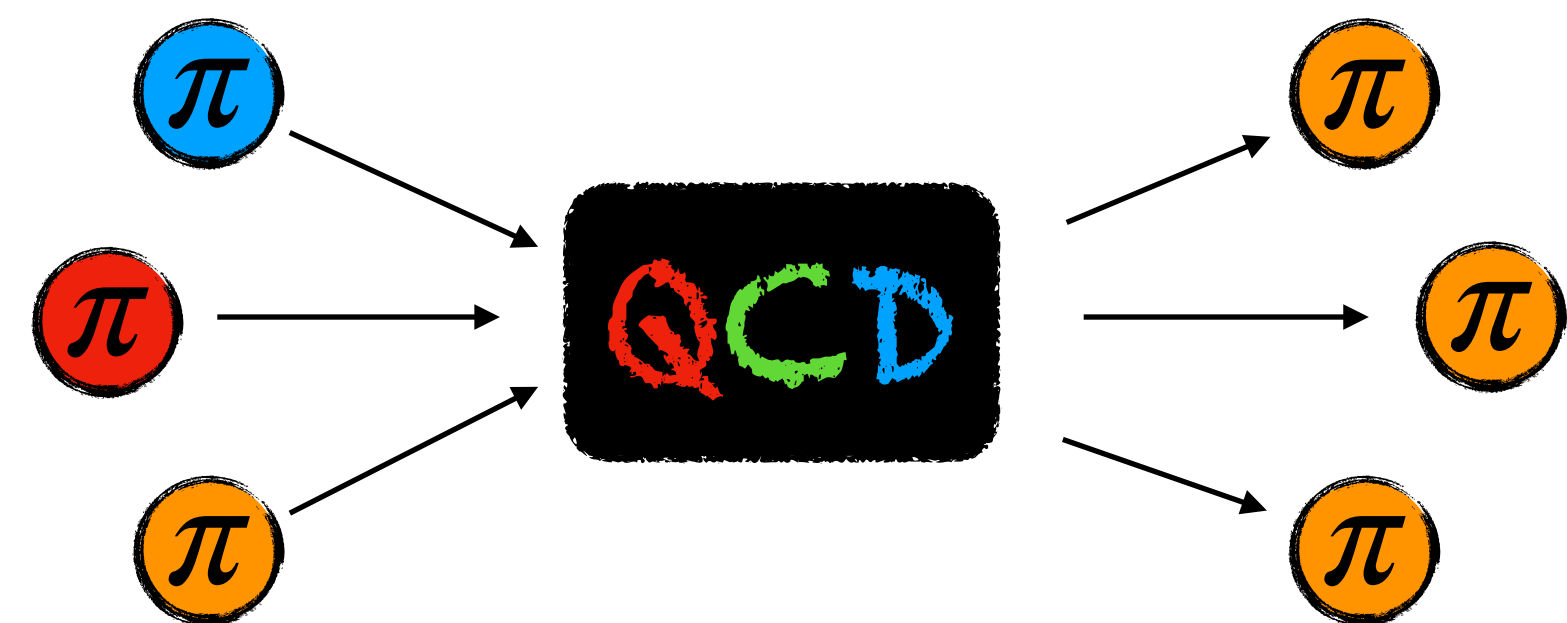
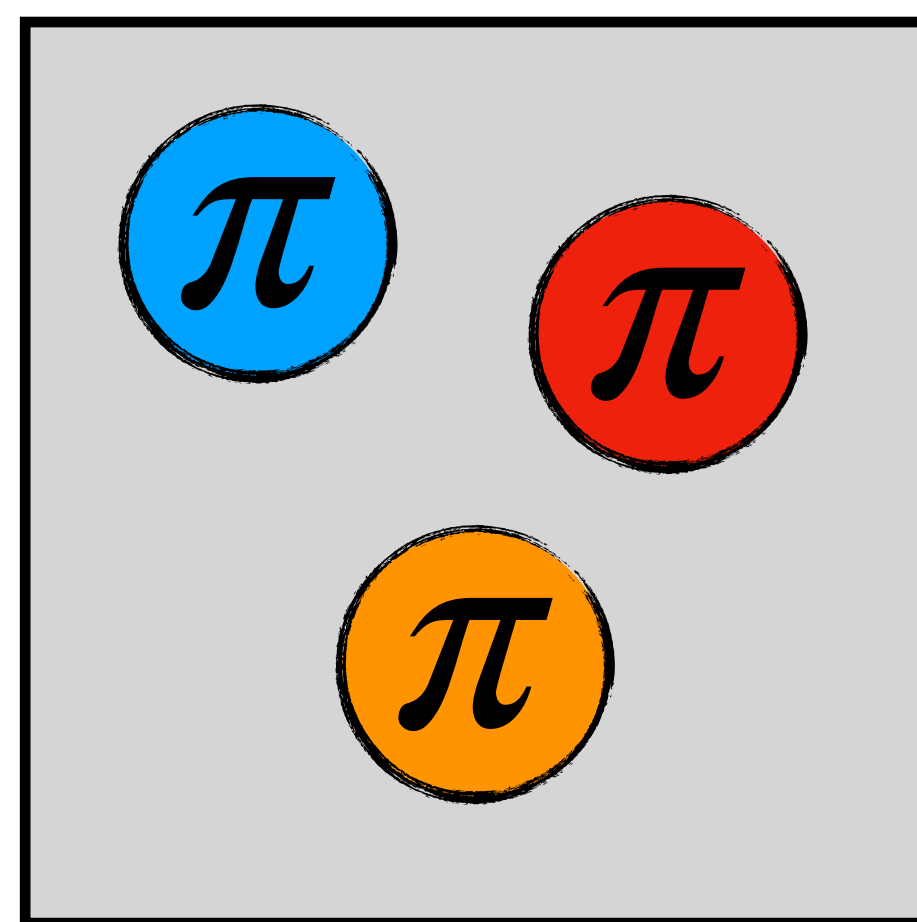
Three-particle scattering amplitudes from Lattice QCD

Fernando Romero-López

University of Valencia

fernando.romero@uv.es

Valencia, 15th Dec



Acknowledgements

IFIC people

Jorge Baeza-Ballesteros
Pilar Hernández

Three-particle people

Tyler Blanton
Raúl Briceño
Drew Hanlon
Max Hansen
Ben Hörz
Steve Sharpe

Bonn Lattice Group

Mathias Fischer
Bartek Kostrzewa
Liuming Liu
Akaki Rusetsky
Nikolas Schlage
Martin Ueding
Carsten Urbach

VNIVERSITAT
ID VALÈNCIA



invisiblesPlus

Introduction

Quantum Chromodynamics

Quantum chromodynamics is conceptually simple. Its realization in nature, however, is usually very complex. But not always.

Frank Wilczek

Quantum Chromodynamics

Quantum chromodynamics is conceptually simple. Its realization in nature, however, is usually very complex. But not always.

Frank Wilczek

- There are still many puzzling hadrons out there
 - XYZ, charmonium, bottomonium, Roper...

Quantum Chromodynamics

Quantum chromodynamics is conceptually simple. Its realization in nature, however, is usually very complex. But not always.

Frank Wilczek

- There are still many puzzling hadrons out there
 - XYZ, charmonium, bottomonium, Roper...
- Nonperturbative QCD dynamics present in relevant processes:
 - CP violating decays: $K \rightarrow \pi\pi$, $K \rightarrow \pi\pi\pi$, $D \rightarrow \pi\pi$, $K\bar{K}$, $(\pi\pi\pi)$, ...

Quantum Chromodynamics

Quantum chromodynamics is conceptually simple. Its realization in nature, however, is usually very complex. But not always.

Frank Wilczek

- There are still many puzzling hadrons out there
 - XYZ, charmonium, bottomonium, Roper...
- Nonperturbative QCD dynamics present in relevant processes:
 - CP violating decays: $K \rightarrow \pi\pi$, $K \rightarrow \pi\pi\pi$, $D \rightarrow \pi\pi$, $K\bar{K}$, $(\pi\pi\pi)$, ...
- First-principles nuclear interactions.

Towards the QCD S-Matrix

The S-Matrix contains the physical information of the theory:

$$S_{ab}(E) \equiv \langle \text{out} | \mathcal{S} | \text{in} \rangle$$

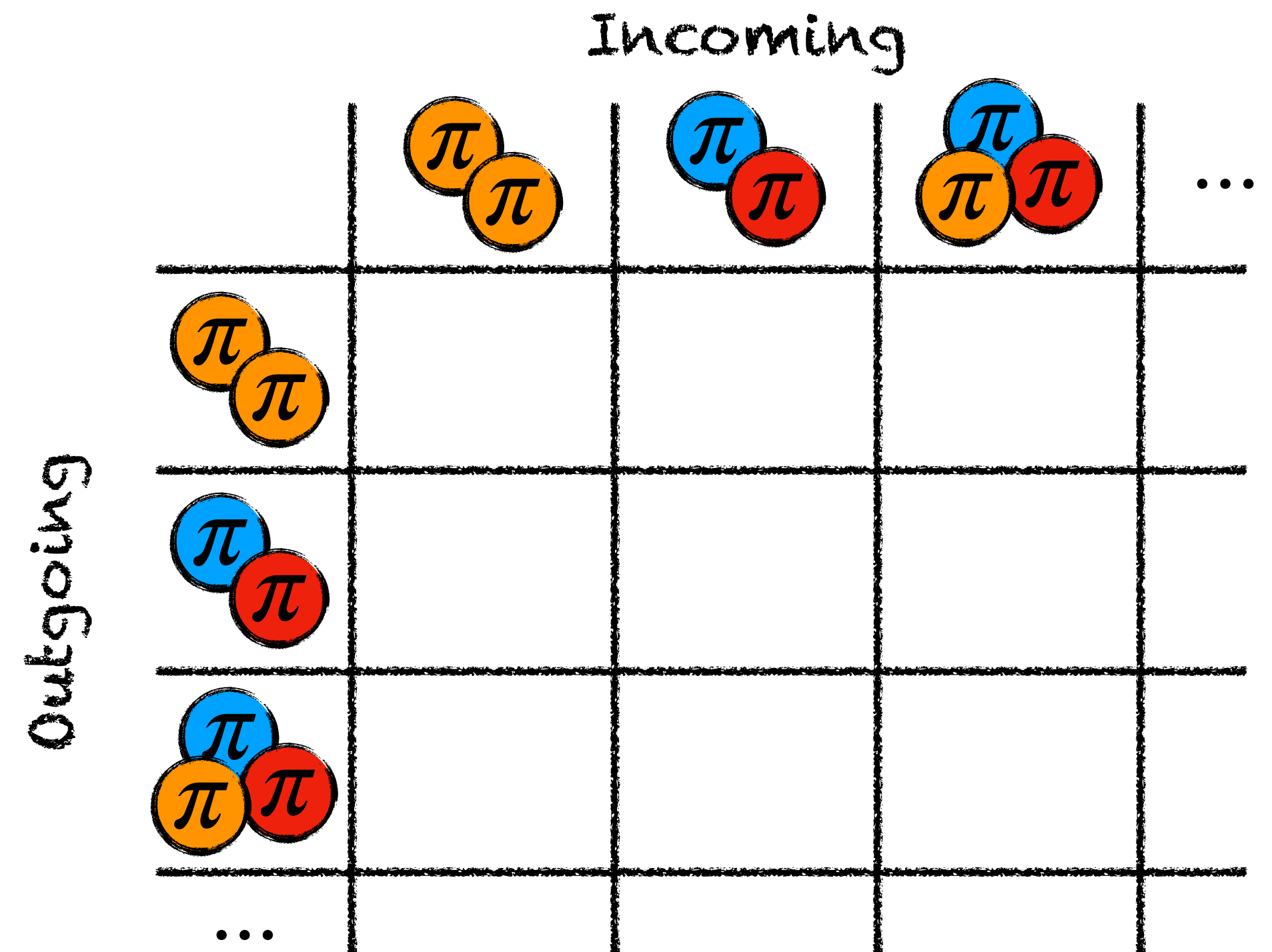
Lattice QCD \longrightarrow QCD S-matrix

Towards the QCD S-Matrix

The S-Matrix contains the physical information of the theory:

$$S_{ab}(E) \equiv \langle \text{out} | \mathcal{S} | \text{in} \rangle$$

Lattice QCD \longrightarrow QCD S-matrix

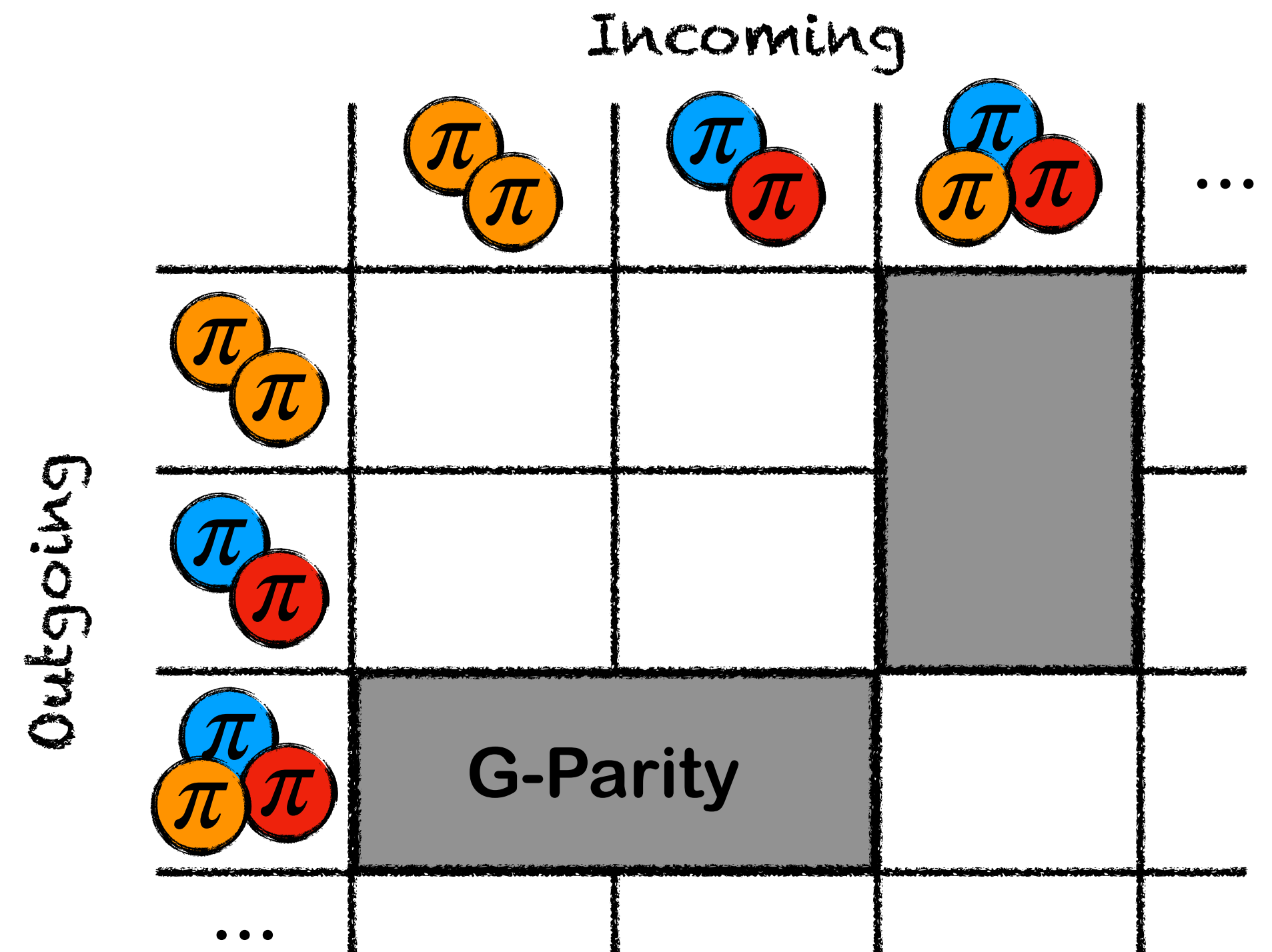


Towards the QCD S-Matrix

The S-Matrix contains the physical information of the theory:

$$S_{ab}(E) \equiv \langle \text{out} | \mathcal{S} | \text{in} \rangle$$

Lattice QCD \longrightarrow QCD S-matrix

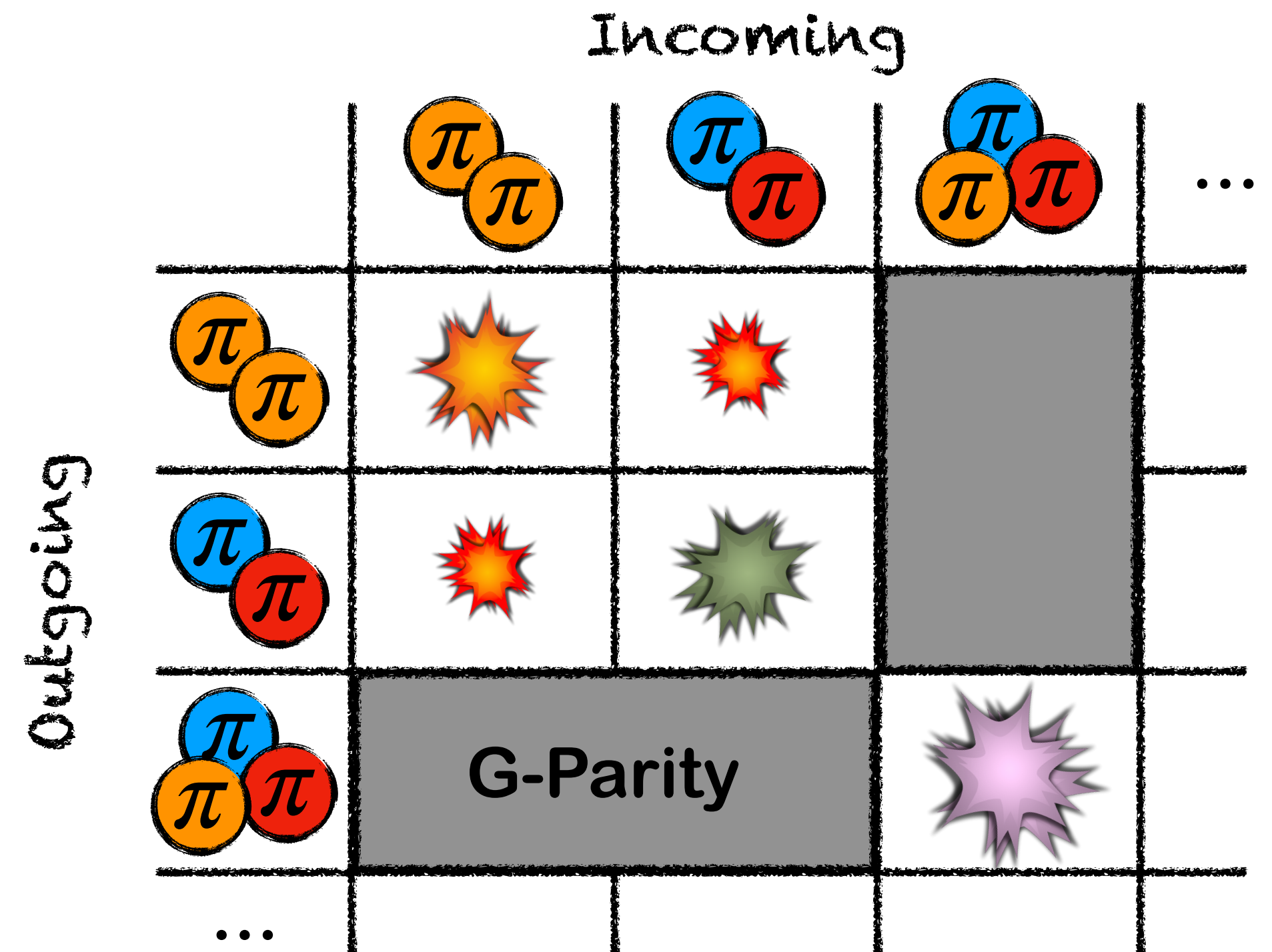


Towards the QCD S-Matrix

The S-Matrix contains the physical information of the theory:

$$S_{ab}(E) \equiv \langle \text{out} | \mathcal{S} | \text{in} \rangle$$

Lattice QCD \longrightarrow QCD S-matrix



Towards the QCD S-Matrix

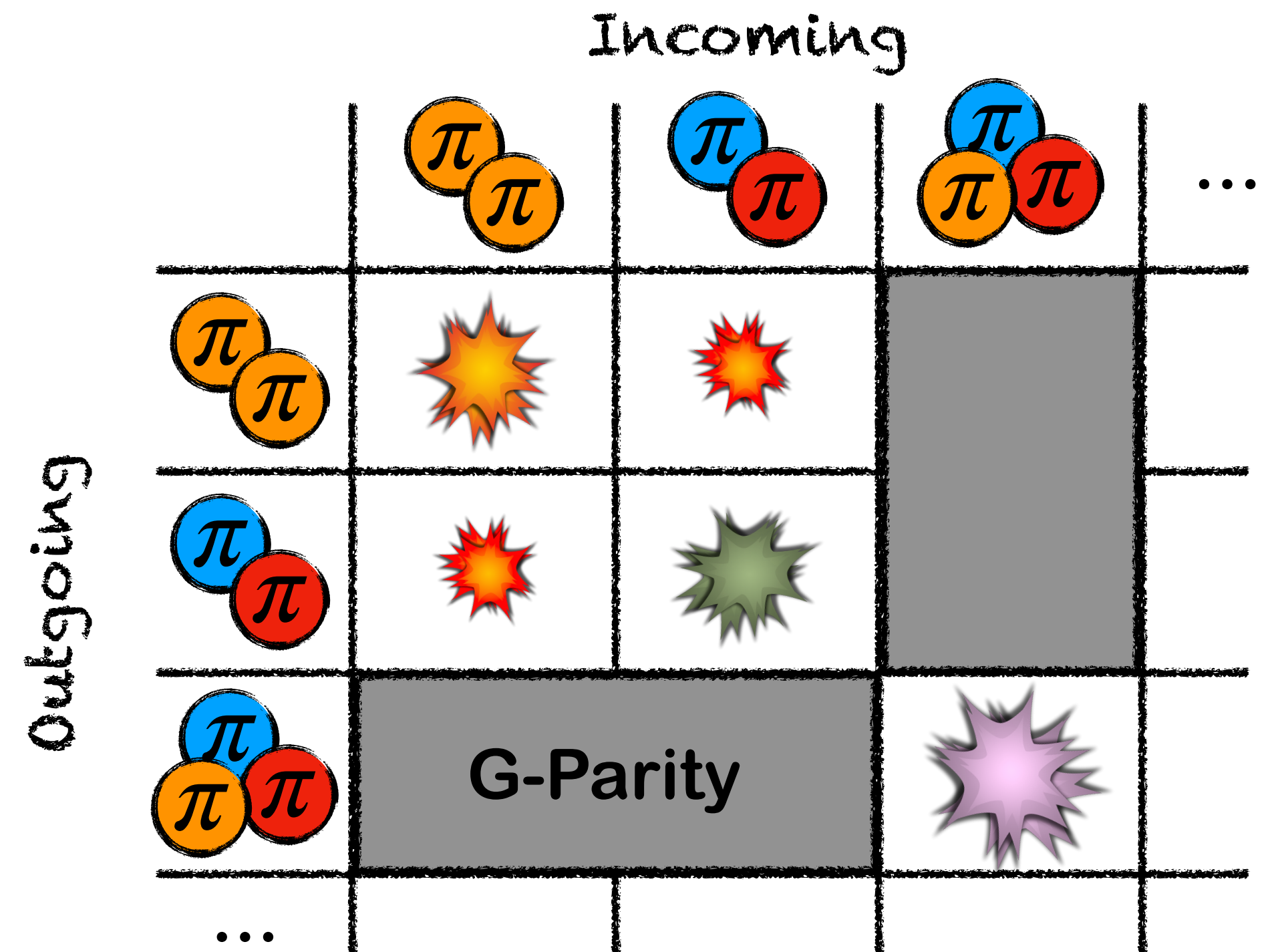
The S-Matrix contains the physical information of the theory:

$$S_{ab}(E) \equiv \langle \text{out} | \mathcal{S} | \text{in} \rangle$$

Lattice QCD \longrightarrow QCD S-matrix

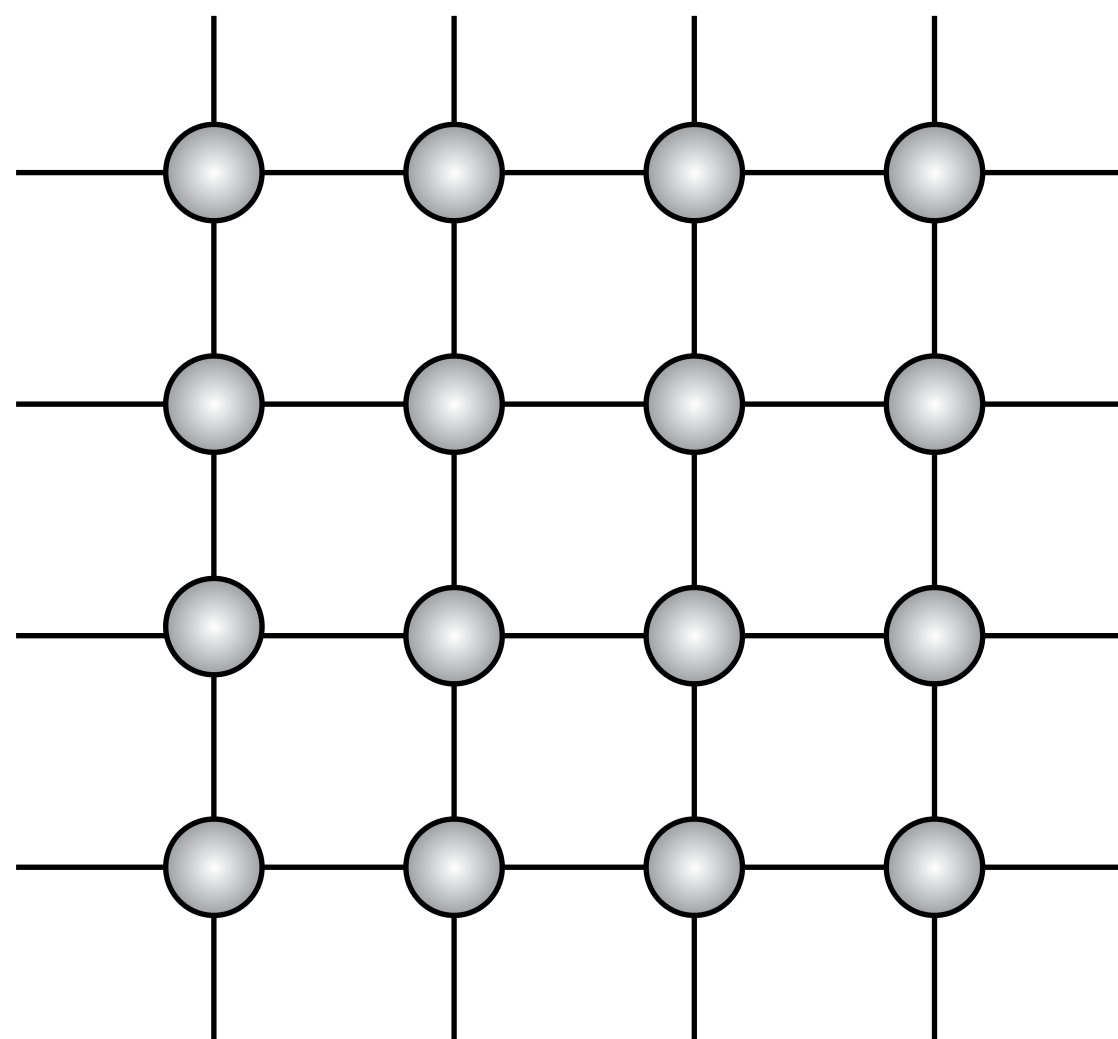
This Talk:

- How to extract S-matrix elements from Lattice QCD:
 - Two-particle scattering in finite volume
 - Lattice results for three particles



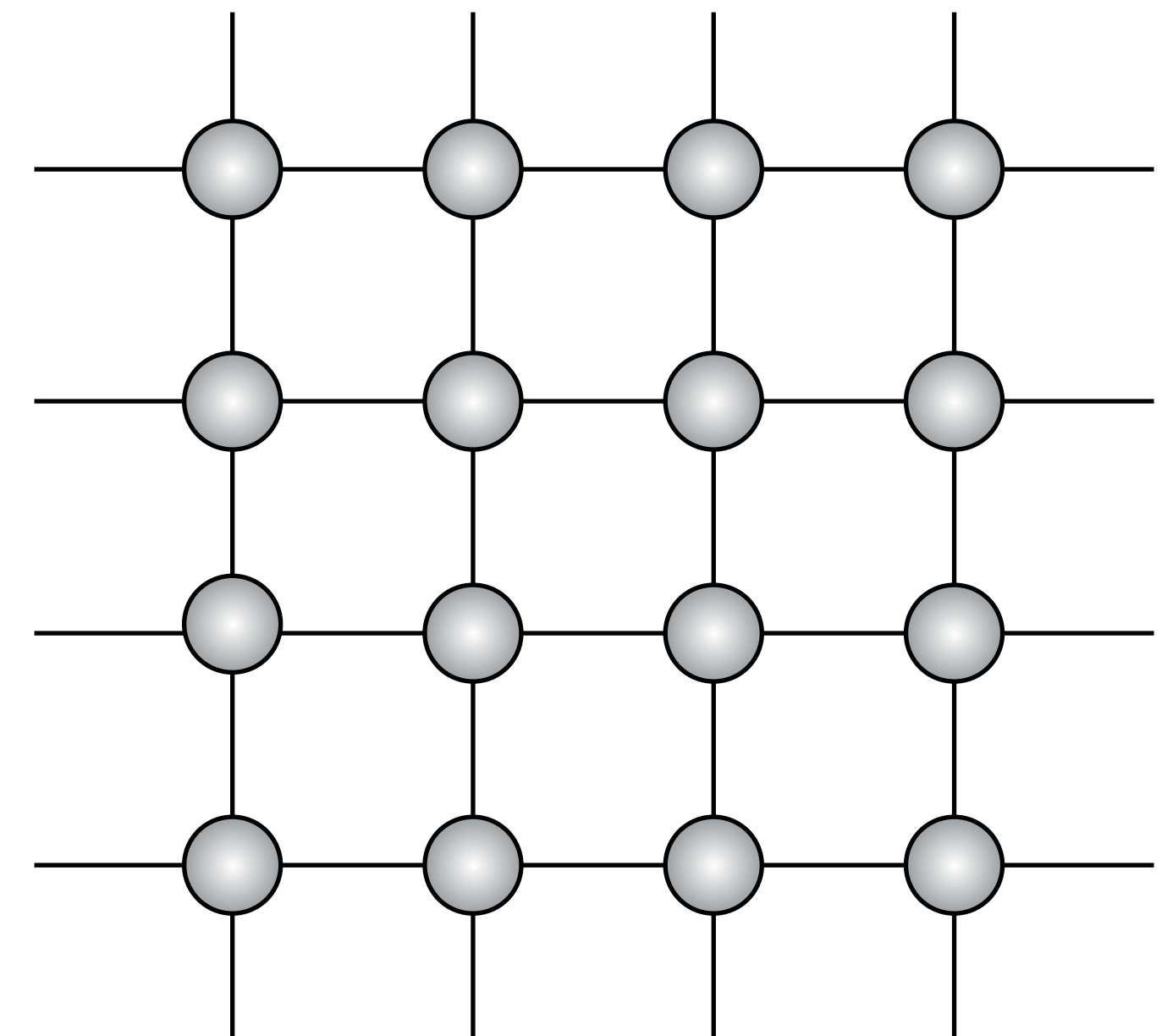
Lattice QCD

basics



Lattice QCD Basics (I)

- Lattice QCD is the state-of-the-art treatment of the strong interaction at hadronic energies

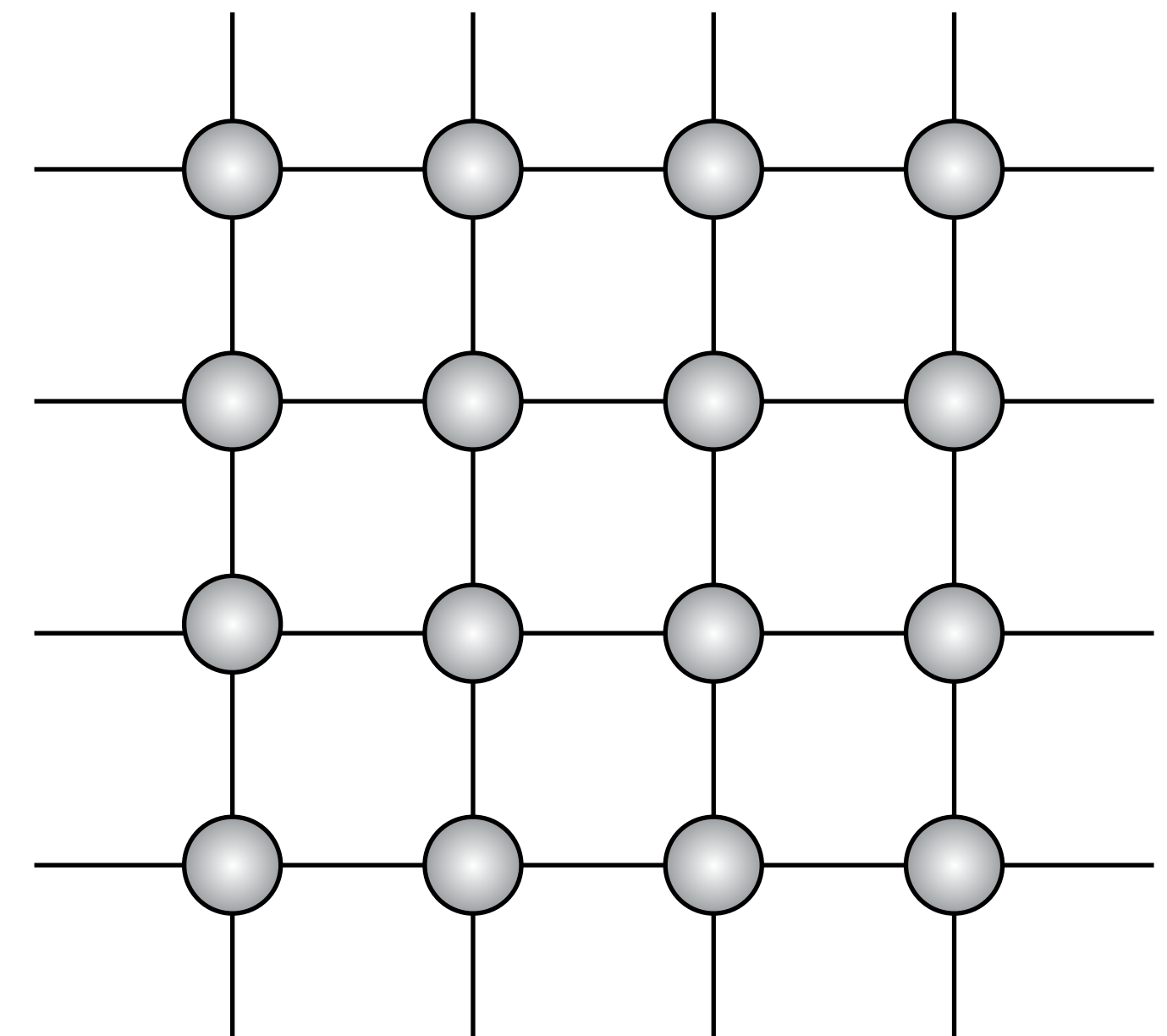


Lattice QCD Basics (I)

○ Lattice QCD is the state-of-the-art treatment of the strong interaction at hadronic energies

● Euclidean time: action has statistical meaning

$$\mathcal{Z} = \int D\psi D\bar{\psi} DA e^{-S_E(\psi, \bar{\psi}, A_\mu)}$$



Lattice QCD Basics (I)

○ Lattice QCD is the state-of-the-art treatment of the strong interaction at hadronic energies

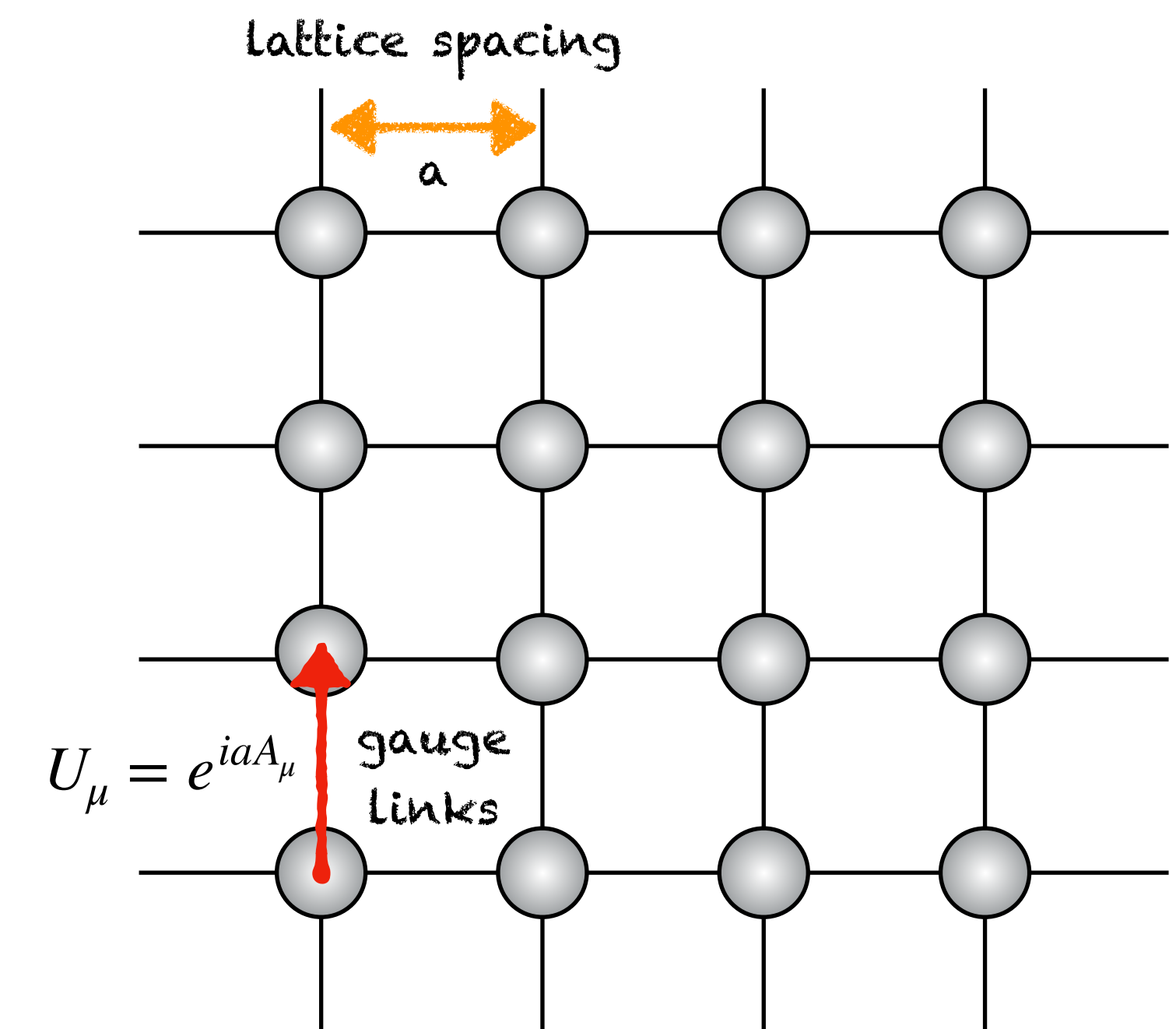
● Euclidean time: action has statistical meaning

$$\mathcal{Z} = \int D\psi D\bar{\psi} DA e^{-S_E(\psi, \bar{\psi}, A_\mu)}$$

● Discretize gauge fields and fermion fields:

→ Under control but technical

(e.g., discretization effects and continuum limit)



Lattice QCD Basics (I)

○ Lattice QCD is the state-of-the-art treatment of the strong interaction at hadronic energies

● Euclidean time: action has statistical meaning

$$\mathcal{Z} = \int D\psi D\bar{\psi} DA e^{-S_E(\psi, \bar{\psi}, A_\mu)}$$

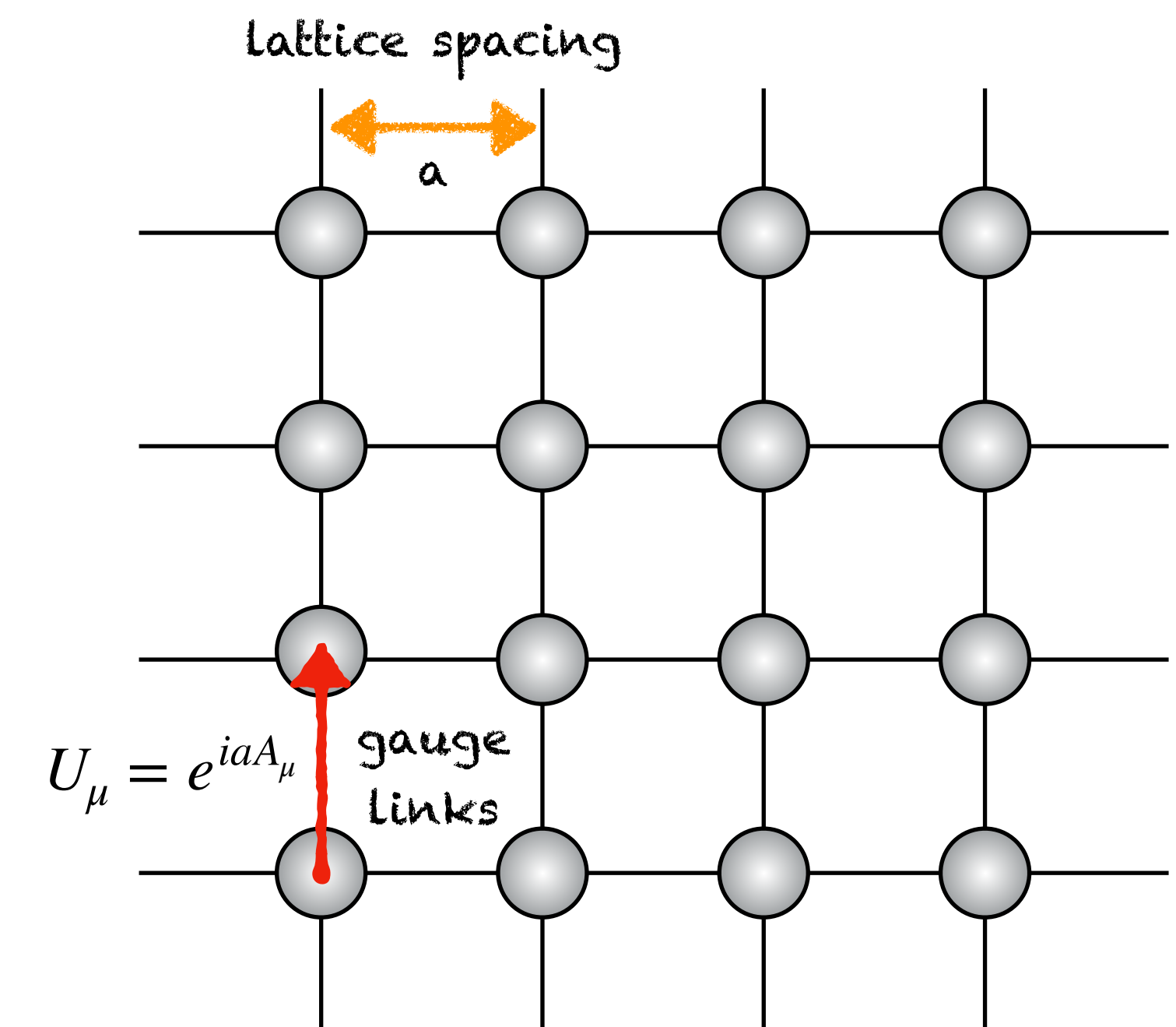
● Discretize gauge fields and fermion fields:

→ Under control but technical

(e.g., discretization effects and continuum limit)

● Compute correlation functions

$$\langle \mathcal{O}(t) \mathcal{O}(0) \rangle = \frac{1}{\mathcal{Z}} \int D\psi D\bar{\psi} DA \mathcal{O}(t) \mathcal{O}(0) e^{-S(\psi, \bar{\psi}, A_\mu)}$$



Lattice QCD Basics (II)

- In Lattice QCD, we measure **energy levels** and **matrix elements**: "Spectral decomposition"

$$\begin{aligned} C(t) = \langle \mathcal{O}(t)\mathcal{O}(0) \rangle &= \sum_n \langle 0 | \mathcal{O}(t) | n \rangle \langle n | \mathcal{O}(0) | 0 \rangle \\ &= \sum_n \left| \langle 0 | \mathcal{O}(0) | n \rangle \right|^2 e^{-E_n t} \end{aligned}$$

Lattice QCD Basics (II)

- In Lattice QCD, we measure **energy levels** and **matrix elements**: "Spectral decomposition"

$$\begin{aligned} C(t) = \langle \mathcal{O}(t)\mathcal{O}(0) \rangle &= \sum_n \langle 0 | \mathcal{O}(t) | n \rangle \langle n | \mathcal{O}(0) | 0 \rangle \\ &= \sum_n \left| \langle 0 | \mathcal{O}(0) | n \rangle \right|^2 e^{-E_n t} \xrightarrow{t \rightarrow \infty} A_0 e^{-E_0 t} \end{aligned}$$

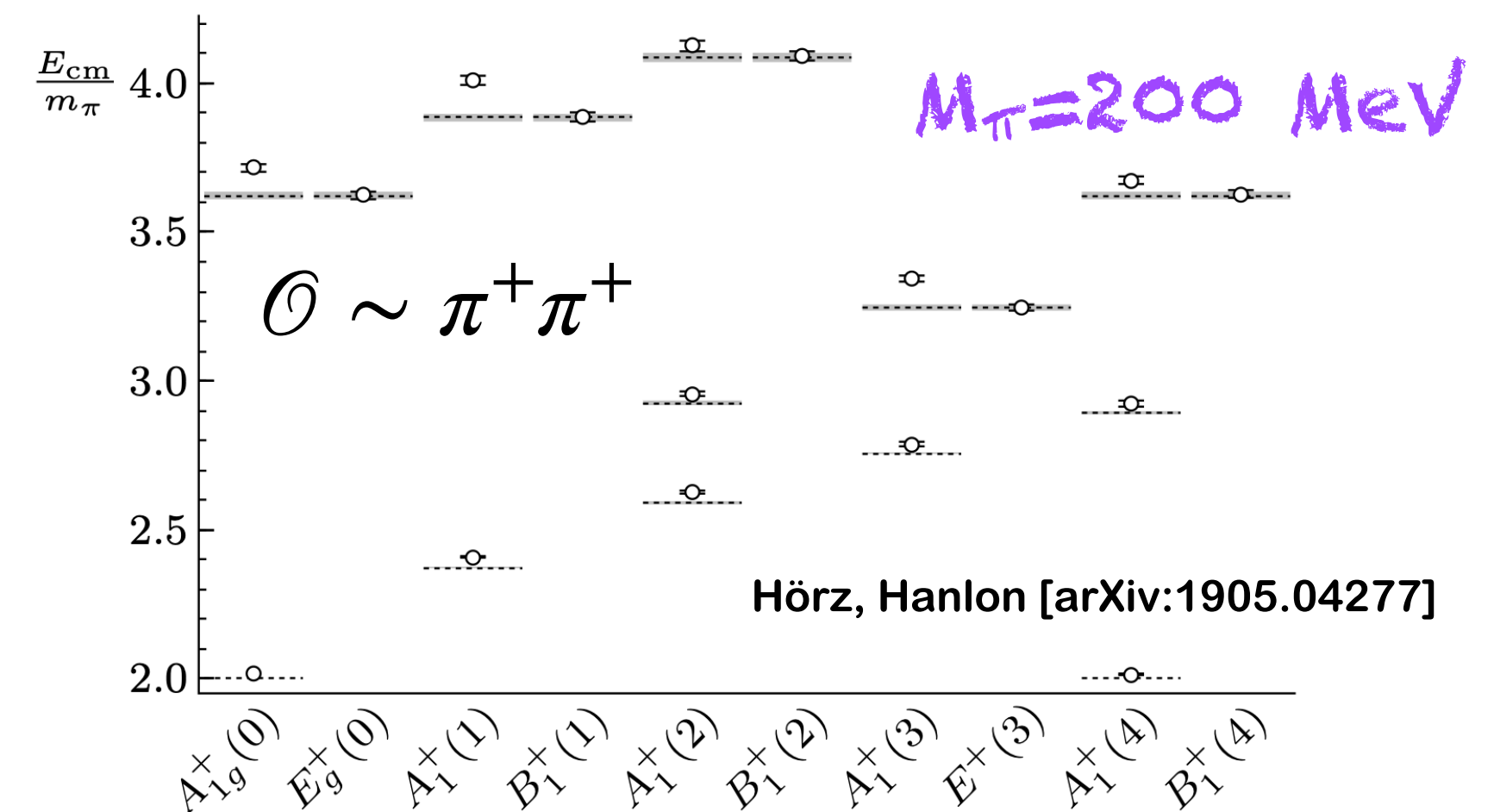
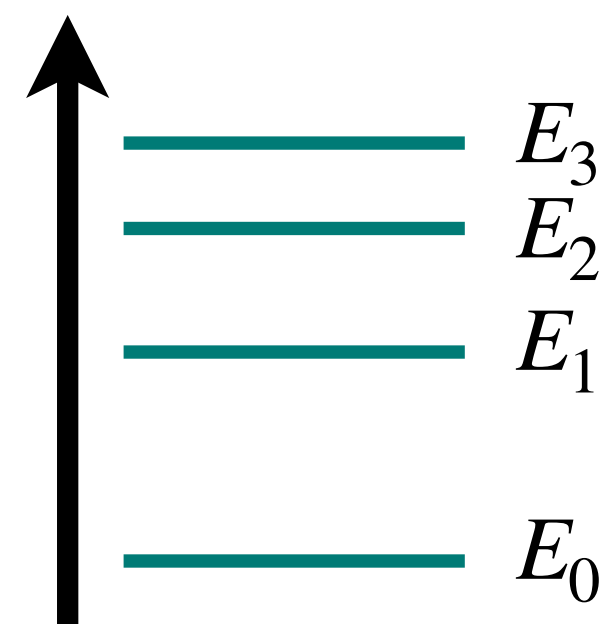
Lattice QCD Basics (II)

- In Lattice QCD, we measure **energy levels** and **matrix elements**: "Spectral decomposition"

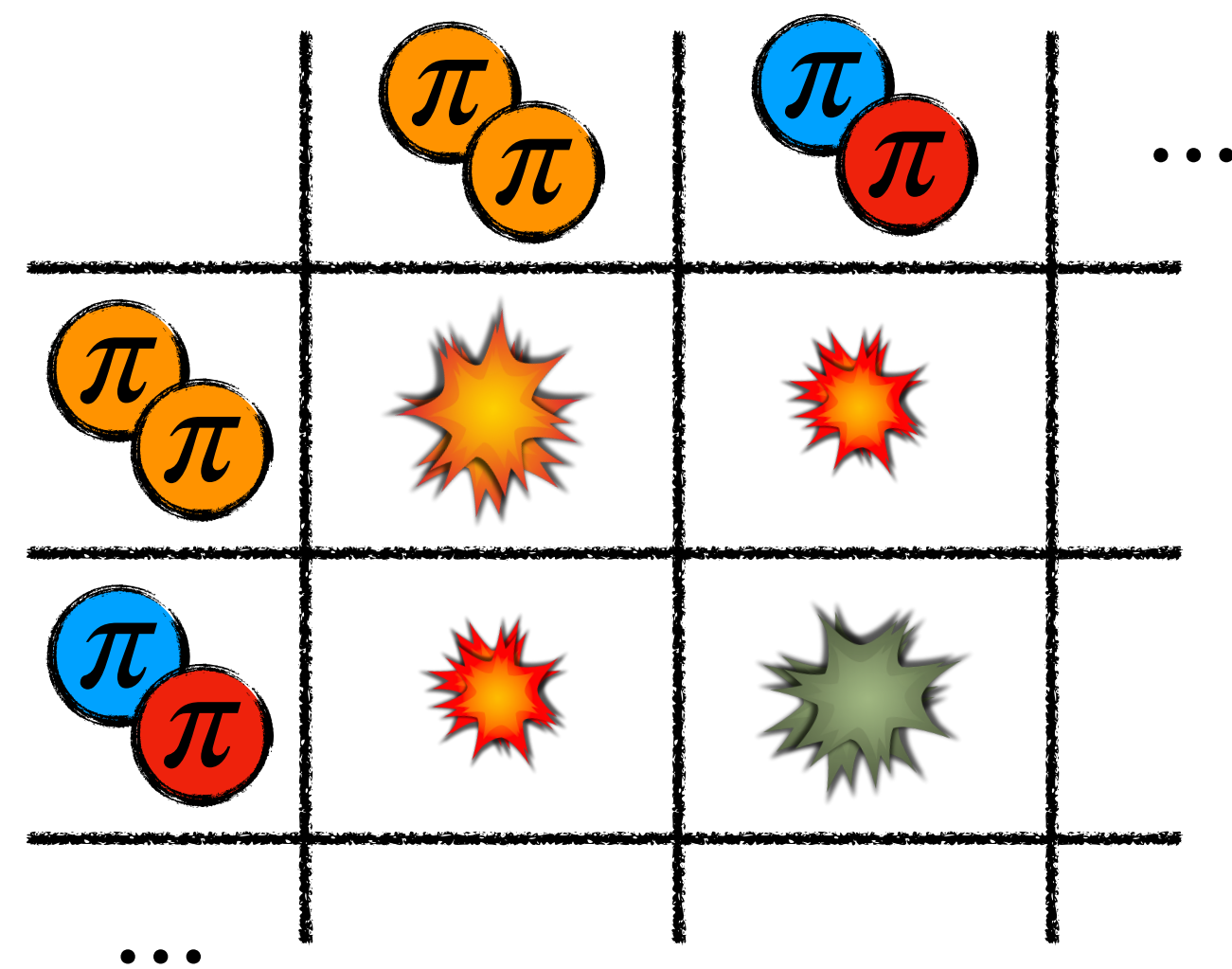
$$\begin{aligned}
 C(t) &= \langle \mathcal{O}(t)\mathcal{O}(0) \rangle = \sum_n \langle 0 | \mathcal{O}(t) | n \rangle \langle n | \mathcal{O}(0) | 0 \rangle \\
 &= \sum_n \left| \langle 0 | \mathcal{O}(0) | n \rangle \right|^2 e^{-E_n t} \xrightarrow{t \rightarrow \infty} A_0 e^{-E_0 t}
 \end{aligned}$$

- Multiple operators to obtain several energy levels

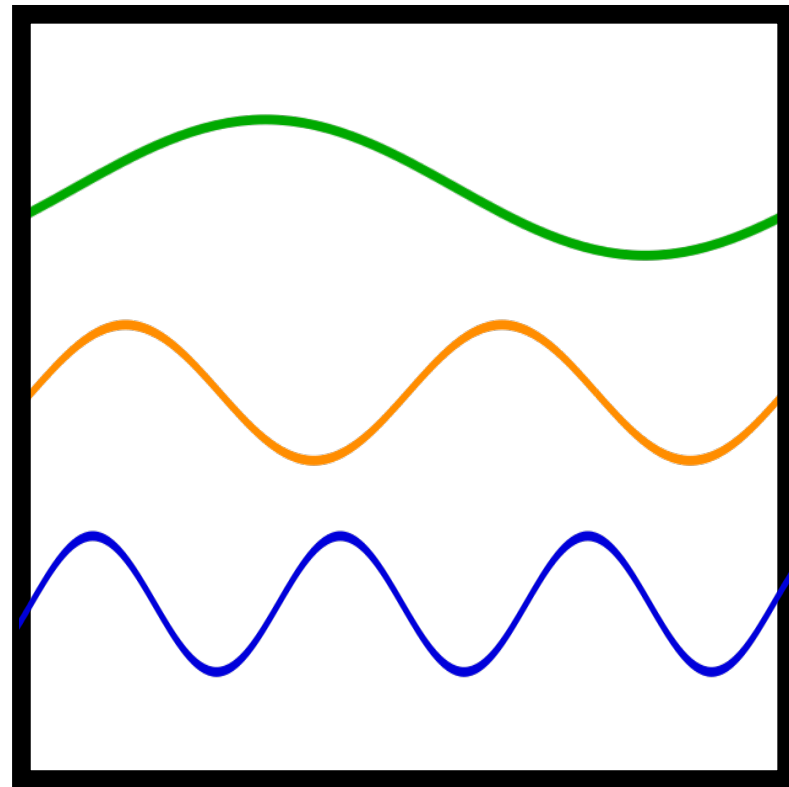
The Spectrum



Two particles in finite-volume

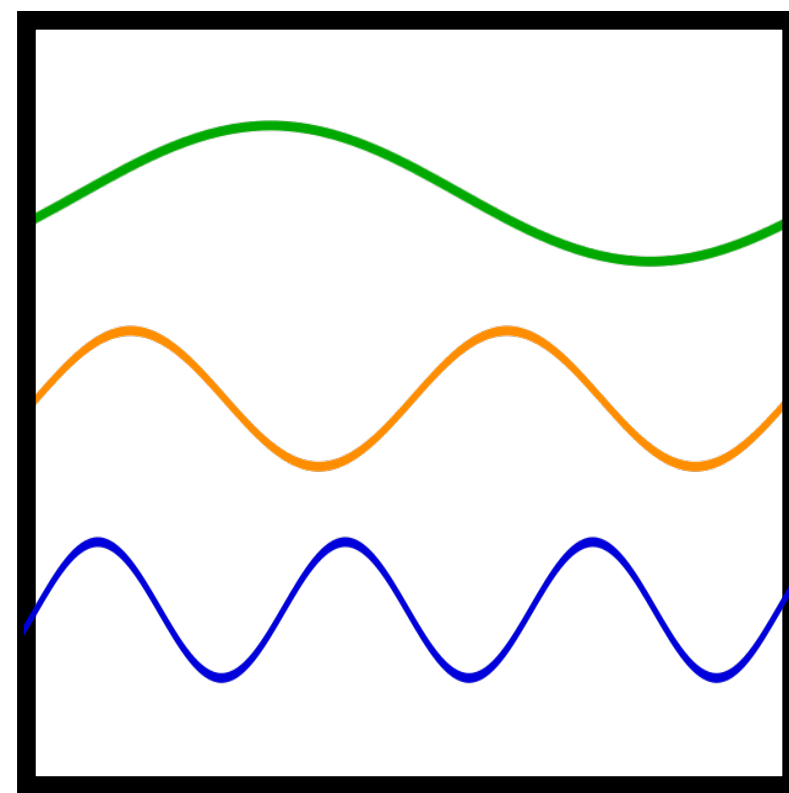


Finite-Volume Spectrum



Finite-Volume Spectrum

Free scalar particles in finite volume
with periodic BC

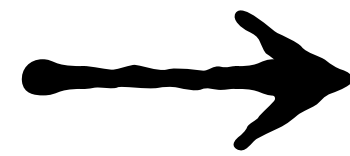


$$\vec{p} = \frac{2\pi}{L}(n_x, n_y, n_z)$$

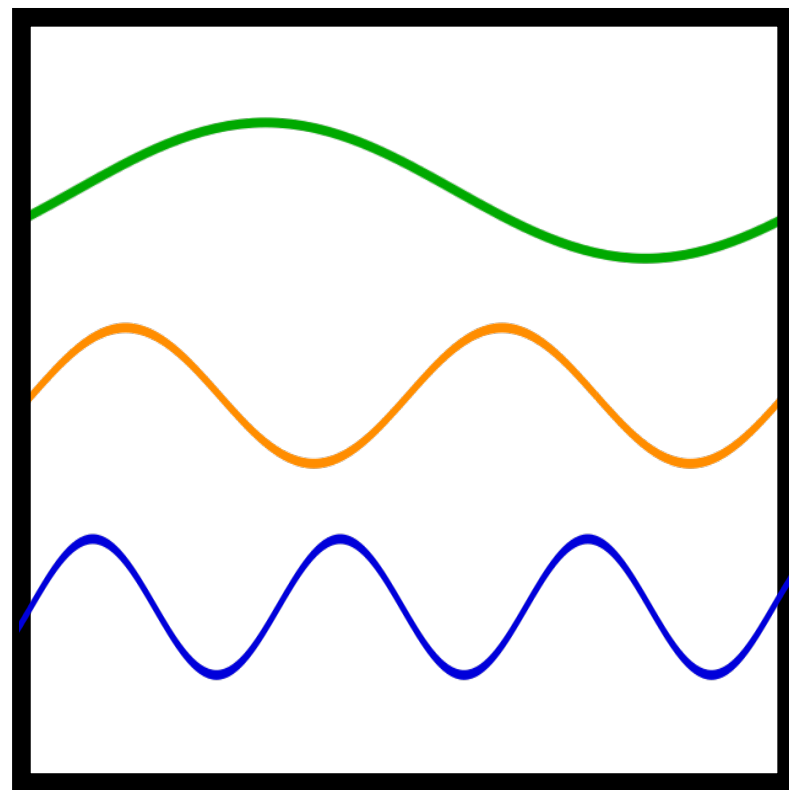
Two particles:
$$E = 2\sqrt{m^2 + \frac{4\pi^2}{L^2} \vec{n}^2}$$

Finite-Volume Spectrum

Free scalar particles in finite volume
with periodic BC



Interactions change the spectrum:
it can be treated as a perturbation

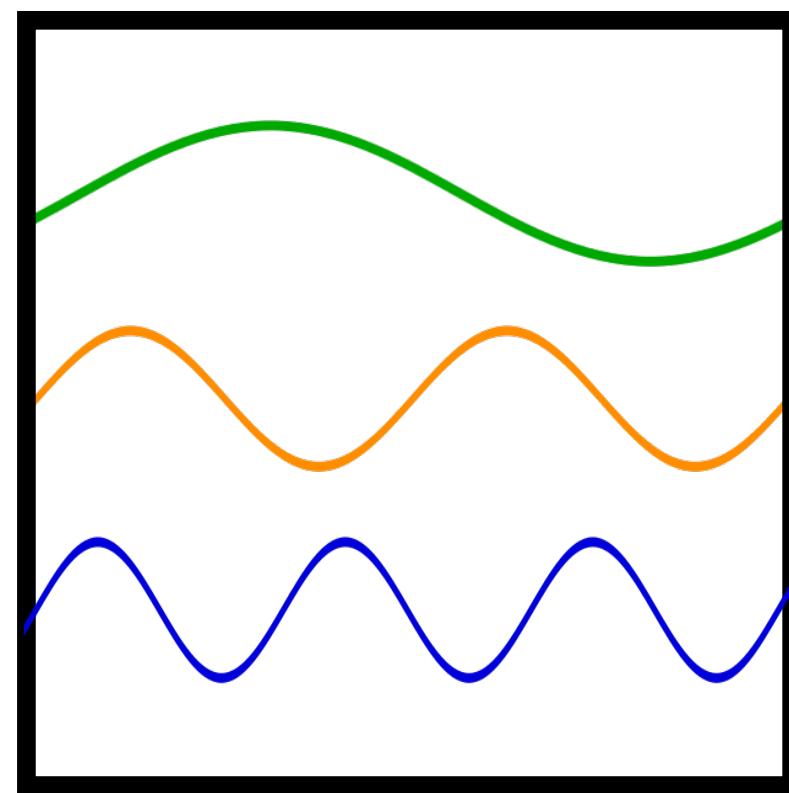


$$\vec{p} = \frac{2\pi}{L}(n_x, n_y, n_z)$$

Two particles:
$$E = 2\sqrt{m^2 + \frac{4\pi^2}{L^2} \vec{n}^2}$$

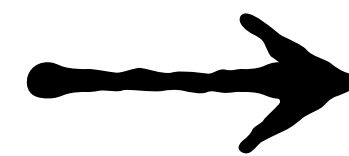
Finite-Volume Spectrum

Free scalar particles in finite volume
with periodic BC



$$\vec{p} = \frac{2\pi}{L}(n_x, n_y, n_z)$$

Two particles: $E = 2\sqrt{m^2 + \frac{4\pi^2}{L^2} \vec{n}^2}$



Interactions change the spectrum:
it can be treated as a perturbation

Ground state to leading order

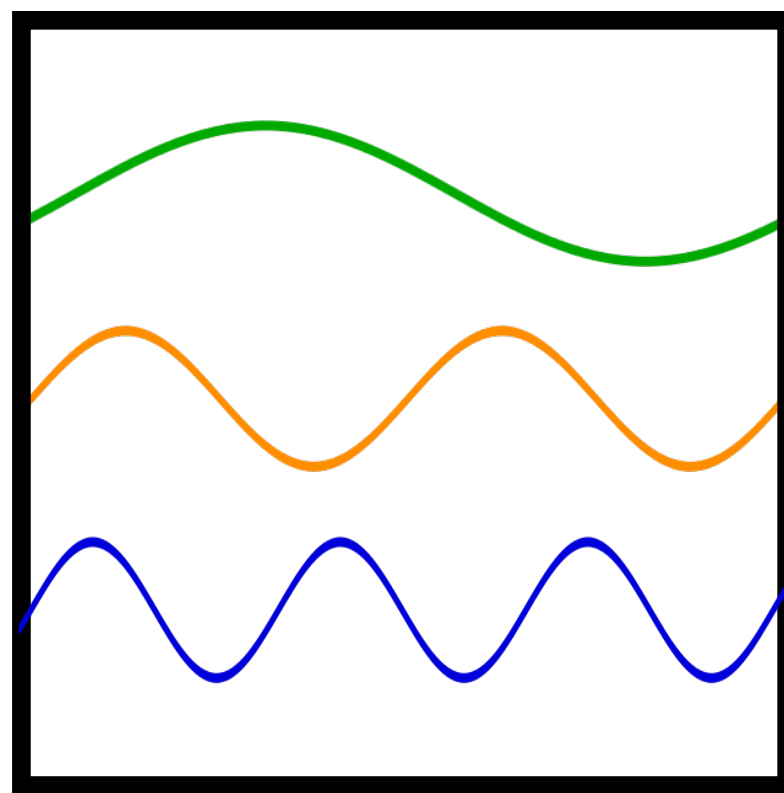
$$E_2 - 2m = \langle \phi(\vec{0})\phi(\vec{0}) | \mathbf{H}_I | \phi(\vec{0})\phi(\vec{0}) \rangle$$

$$\Delta E_2 = \frac{\mathcal{M}_2(E = 2m)}{8m^2L^3} + O(L^{-4})$$

[Huang, Yang, 1958]

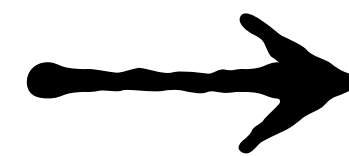
Finite-Volume Spectrum

Free scalar particles in finite volume
with periodic BC



$$\vec{p} = \frac{2\pi}{L}(n_x, n_y, n_z)$$

Two particles: $E = 2\sqrt{m^2 + \frac{4\pi^2}{L^2} \vec{n}^2}$



Interactions change the spectrum:
it can be treated as a perturbation

Ground state to leading order

$$E_2 - 2m = \langle \phi(\vec{0})\phi(\vec{0}) | \mathbf{H}_I | \phi(\vec{0})\phi(\vec{0}) \rangle$$

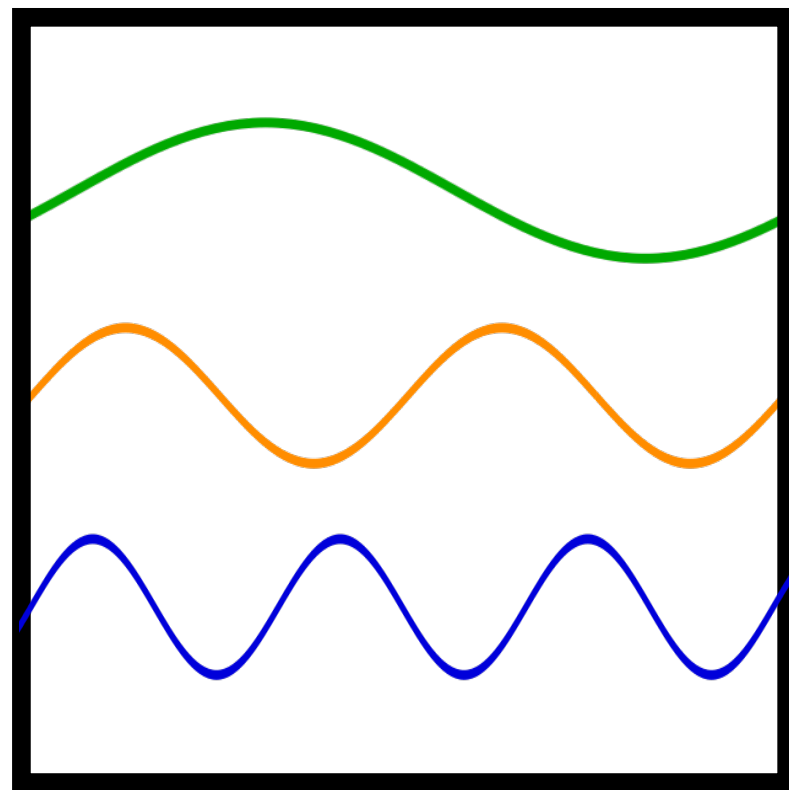
$$\Delta E_2 = \frac{\mathcal{M}_2(E = 2m)}{8m^2L^3} + O(L^{-4})$$

[Huang, Yang, 1958]

The **energy shift** of the two-particle ground state
is related to the $2 \rightarrow 2$ **scattering amplitude**

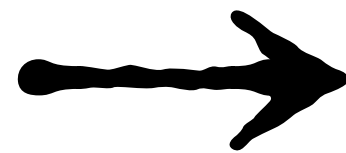
Finite-Volume Spectrum

Free scalar particles in finite volume with periodic BC



$$\vec{p} = \frac{2\pi}{L}(n_x, n_y, n_z)$$

Two particles: $E = 2\sqrt{m^2 + \frac{4\pi^2}{L^2} \vec{n}^2}$



Interactions change the spectrum: it can be treated as a perturbation

In general a problem of Quantum Field Theory in finite volume

ground state to leading order

$$E_2 = \langle \phi(\vec{0})\phi(\vec{0}) | \mathbf{H}_I | \phi(\vec{0})\phi(\vec{0}) \rangle$$

$$\Delta E_2 = \frac{\mathcal{M}_2(E = 2m)}{8m^2L^3} + O(L^{-4})$$

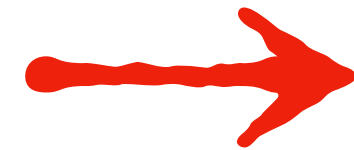
[Huang, Yang, 1958]

The **energy shift** of the two-particle ground state is related to the $2 \rightarrow 2$ **scattering amplitude**

The Lüscher Formalism

- A new field was opened by M. Lüscher in '86

finite-volume
spectrum of
two identical
scalars



s-wave
scattering
amplitude

Volume Dependence of the Energy Spectrum in Massive Quantum Field Theories

II. Scattering States

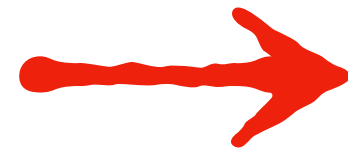
M. Lüscher

Theory Division, Deutsches Elektronen-Synchrotron DESY, D-2000 Hamburg 52, Federal Republic of Germany

The Lüscher Formalism

- A new field was opened by M. Lüscher in '86

finite-volume
spectrum of
two identical
scalars



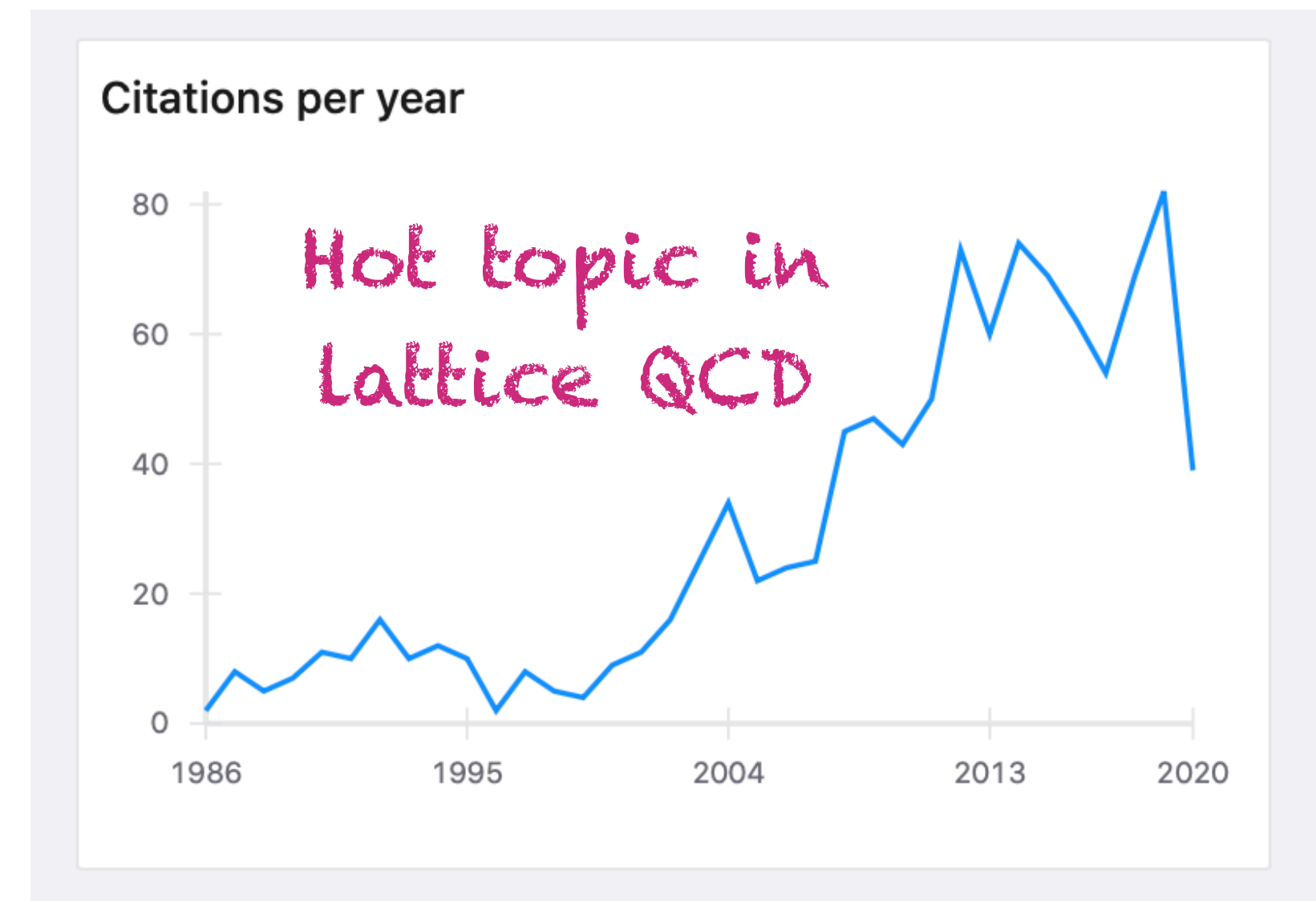
s-wave
scattering
amplitude

Volume Dependence of the Energy Spectrum in Massive Quantum Field Theories

II. Scattering States

M. Lüscher

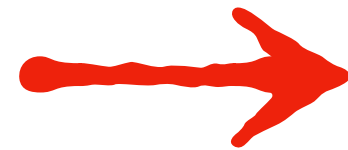
Theory Division, Deutsches Elektronen-Synchrotron DESY, D-2000 Hamburg 52, Federal Republic of Germany



The Lüscher Formalism

- A new field was opened by M. Lüscher in '86

finite-volume
spectrum of
two identical
scalars



s-wave
scattering
amplitude

- Fully general formalism exists up to date:

- Multichannel, non-identical $2 \rightarrow 2$ scattering for particles with spin in all partial waves. Including for weak decays, such as $K \rightarrow \pi\pi$ (Lellouch-Lüscher)

- Many people have contributed over the years:

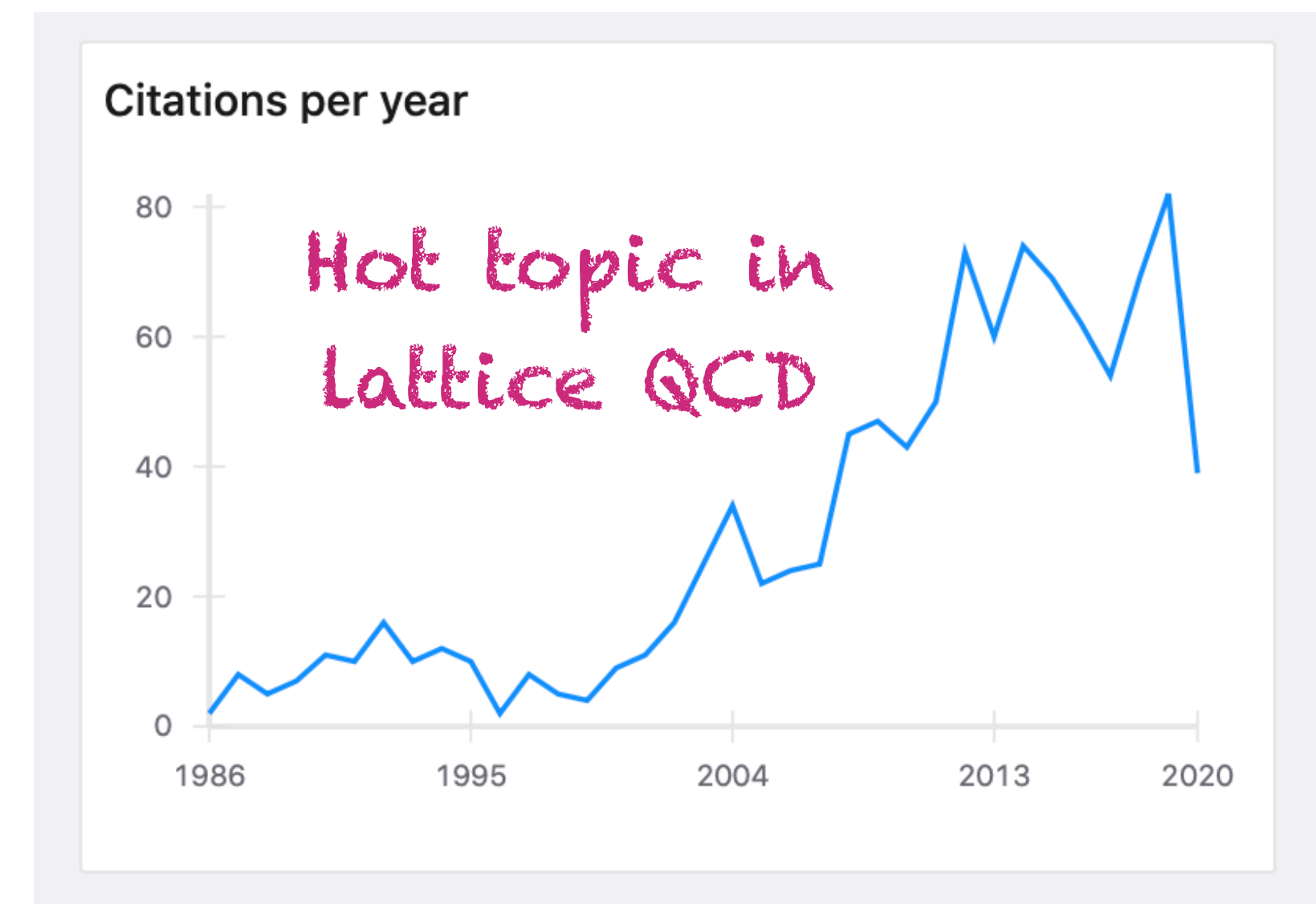
- ▶ Rummukainen and Gottlieb
- ▶ Kim, Sachrajda and Sharpe
- ▶ Göckeler, Horsley, Lage, Meißner, Rakow, Rusetsky, Schierholz, Zanotti
- ▶ Briceño

Volume Dependence of the Energy Spectrum in Massive Quantum Field Theories

II. Scattering States

M. Lüscher

Theory Division, Deutsches Elektronen-Synchrotron DESY, D-2000 Hamburg 52, Federal Republic of Germany



Quantization Condition(I)

- In order to derive the full relation, consider the finite-volume correlator:

$$C_L(E, \vec{P}) = \int e^{iPx} \langle \mathcal{O}(x) | \mathcal{O}(0) \rangle =$$

[à la Kim, Sachrajda, Sharpe]

Quantization Condition(I)

- In order to derive the full relation, consider the finite-volume correlator:

Skeleton expansion

$$C_L(E, \vec{P}) = \int e^{iPx} \langle \mathcal{O}(x) | \mathcal{O}(0) \rangle = \text{Diagram 1} + \text{Diagram 2} + \text{Diagram 3} + \dots$$

[à la Kim, Sachrajda, Sharpe]

Quantization Condition(I)

- In order to derive the full relation, consider the finite-volume correlator:

$$C_L(E, \vec{P}) = \int e^{iPx} \langle \mathcal{O}(x) | \mathcal{O}(0) \rangle = \text{Skeleton expansion} + \dots$$

[à la Kim, Sachrajda, Sharpe]

$\sum_{\vec{k}}$
Finite-volume sums

Quantization Condition(I)

○ In order to derive the full relation, consider the finite-volume correlator:

$$C_L(E, \vec{P}) = \int e^{iPx} \langle \mathcal{O}(x) | \mathcal{O}(0) \rangle = \text{Skeleton expansion} + \dots$$

[à la Kim, Sachrajda, Sharpe]

$B_2 = \text{Bethe-Salpeter Kernels} + \dots$

$\sum_{\vec{k}}$
 Finite-volume sums

Quantization Condition(I)

○ In order to derive the full relation, consider the finite-volume correlator:

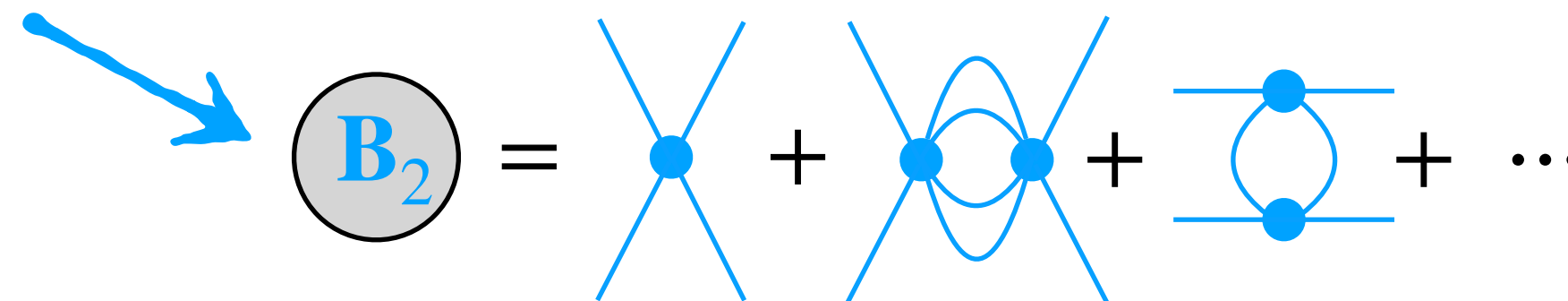
$$C_L(E, \vec{P}) = \int e^{iPx} \langle \mathcal{O}(x) | \mathcal{O}(0) \rangle = \text{Skeleton expansion} + \dots$$

[à la Kim, Sachrajda, Sharpe]

Only exponentially small effects in L

Bethe-Salpeter Kernels

$\sum_{\vec{k}}$
Finite-volume sums



Quantization Condition(I)

○ In order to derive the full relation, consider the finite-volume correlator:

$$C_L(E, \vec{P}) = \int e^{iPx} \langle \mathcal{O}(x) | \mathcal{O}(0) \rangle = \text{Skeleton expansion} + \dots$$

[à la Kim, Sachrajda, Sharpe]

Only exponentially small effects in L

Bethe-Salpeter Kernels

Finite-volume sums

$$\sum_{\vec{k}} \rightarrow \int d^3k + \left[\sum_{\vec{k}} - \int d^3k \right]$$

$$B_2 = \text{[diagrams]} + \dots$$

Quantization Condition(I)

○ In order to derive the full relation, consider the finite-volume correlator:

$$C_L(E, \vec{P}) = \int e^{iPx} \langle \mathcal{O}(x) | \mathcal{O}(0) \rangle = \text{Skeleton expansion} + \dots$$

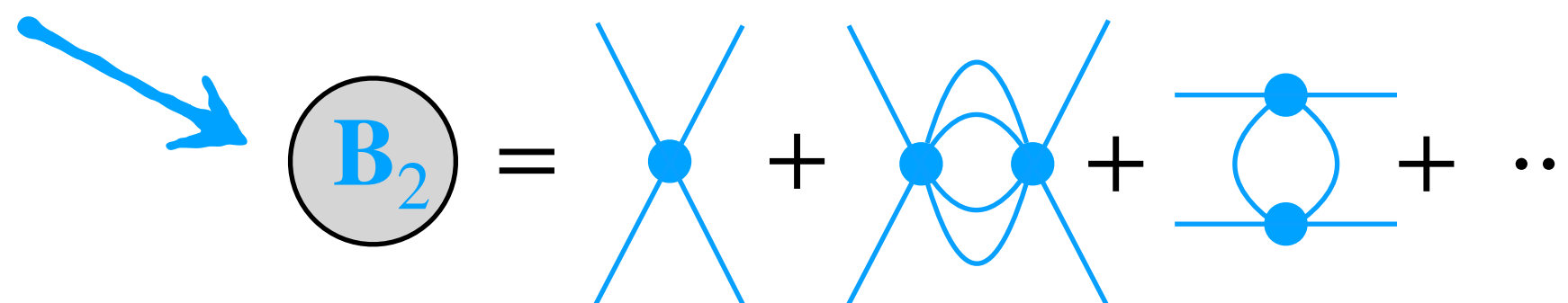
[à la Kim, Sachrajda, Sharpe]

Only exponentially small effects in L

Bethe-Salpeter Kernels

Finite-volume sums

$$\sum_{\vec{k}} \rightarrow \int d^3k + \left[\sum_{\vec{k}} - \int d^3k \right]$$



1. Separation of finite-volume effects
2. Resummation of diagrams

Quantization Condition(I)

In order to derive the full relation, consider the finite-volume correlator:

Skeleton expansion

$$C_L(E, \vec{P}) = \int e^{iPx} \langle \mathcal{O}(x) | \mathcal{O}(0) \rangle = \text{Diagram 1} + \text{Diagram 2} + \text{Diagram 3} + \dots$$

[à la Kim, Sachrajda, Sharpe]

Only exponentially small effects in L

Bethe-Salpeter Kernels

Finite-volume sums

$$\sum_{\vec{k}} \rightarrow \int d^3k + \left[\sum_{\vec{k}} - \int d^3k \right]$$

$$B_2 = \text{Diagram 1} + \text{Diagram 2} + \text{Diagram 3} + \dots$$

Known kinematic function

1. Separation of finite-volume effects
2. Resummation of diagrams

$$C_L(E, \vec{P}) = C_\infty(E, \vec{P}) + A^\dagger \frac{1}{\mathcal{K}_2 + F^{-1}} A + O(e^{-mL})$$

$$\mathcal{M}_2^{-1} = \mathcal{K}_2^{-1} - i\sqrt{s - 2m^2}$$

Quantization Condition (II)

$$C_L(E, \vec{P}) = \text{some algebra ...} = C_\infty(E, \vec{P}) + A^\dagger \frac{1}{\mathcal{K}_2 + F^{-1}} A + O(e^{-mL})$$

Quantization Condition (II)

$$C_L(E, \vec{P}) = \text{some algebra ...} = C_\infty(E, \vec{P}) + A^\dagger \frac{1}{\mathcal{K}_2 + F^{-1}} A + O(e^{-mL})$$

K-matrix parametrized
in terms of phase shift

$$\mathcal{K}_2^\ell = \frac{16\pi\sqrt{s}}{q^{2\ell+1} \cot \delta_\ell}$$

Quantization Condition (II)

$$C_L(E, \vec{P}) = \text{some algebra ...} = C_\infty(E, \vec{P}) + A^\dagger \frac{1}{\mathcal{K}_2 + F^{-1}} A + O(e^{-mL})$$

K-matrix parametrized
in terms of phase shift

$$\mathcal{K}_2^\ell = \frac{16\pi\sqrt{s}}{q^{2\ell+1} \cot \delta_\ell}$$

Finite-volume states appear
when the correlation function
has a pole

Quantization Condition (II)

$$C_L(E, \vec{P}) = \text{some algebra ...} = C_\infty(E, \vec{P}) + A^\dagger \frac{1}{\mathcal{K}_2 + F^{-1}} A + O(e^{-mL})$$

K-matrix parametrized
in terms of phase shift

$$\mathcal{K}_2^\ell = \frac{16\pi\sqrt{s}}{q^{2\ell+1} \cot \delta_\ell}$$

Finite-volume states appear
when the correlation function
has a pole

Two-particle Quantization Condition

$$\det \left[\mathcal{K}_2(E_n) + F^{-1}(E_n, \vec{P}, L) \right] = 0$$

Scattering
K-Matrix

Known kinematic
function

"QC2"

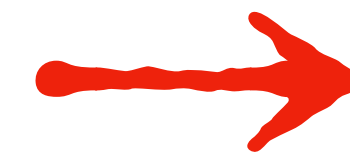
! It holds below $E_{cm} < 4m$

Isospin-2 $\pi\pi$ scattering

Two pions in s-wave

$$\mathcal{K}_2^{s\text{-wave}}(E_n) = \frac{-1}{F_{00}(E_n, \vec{P}, L)}$$

one
energy
level



a phase
shift
point

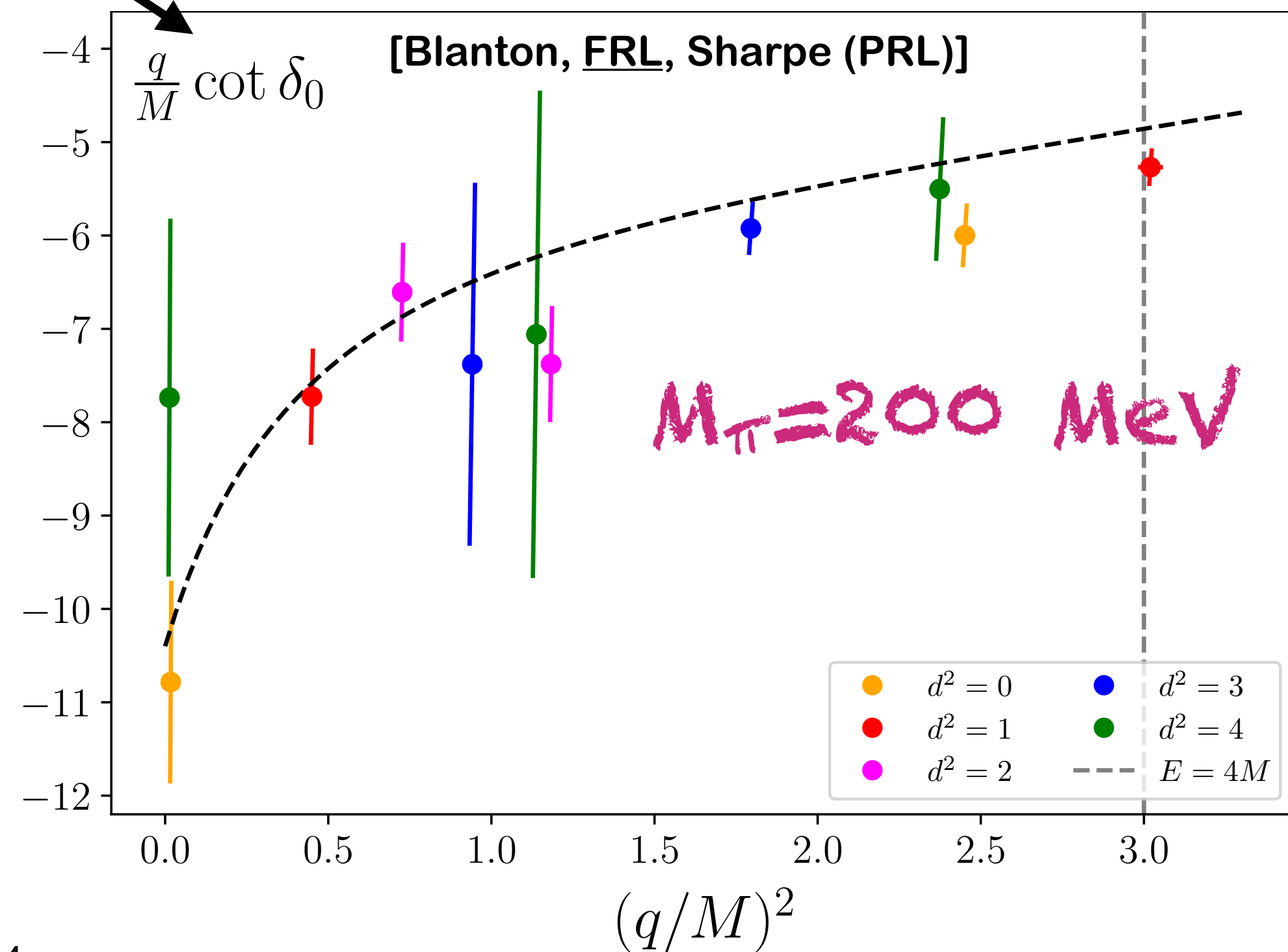
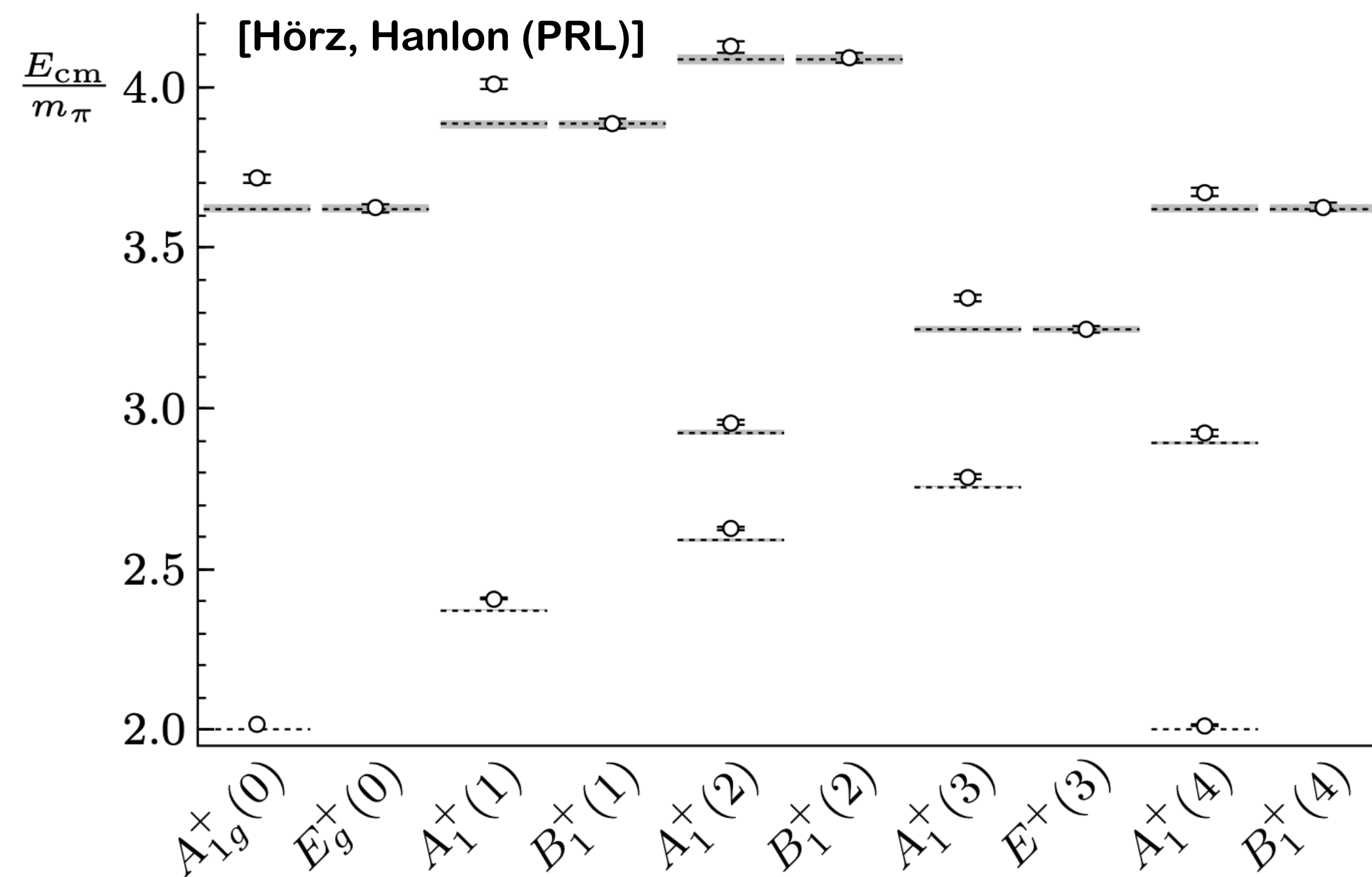
Isospin-2 $\pi\pi$ scattering

Two pions in s-wave

$$\mathcal{K}_2^{s\text{-wave}}(E_n) = \frac{-1}{F_{00}(E_n, \vec{P}, L)}$$

$$\mathcal{K}_2^{s\text{-wave}} \sim \frac{16\pi\sqrt{s}}{q \cot \delta_0}$$

one energy level \longrightarrow a phase shift point



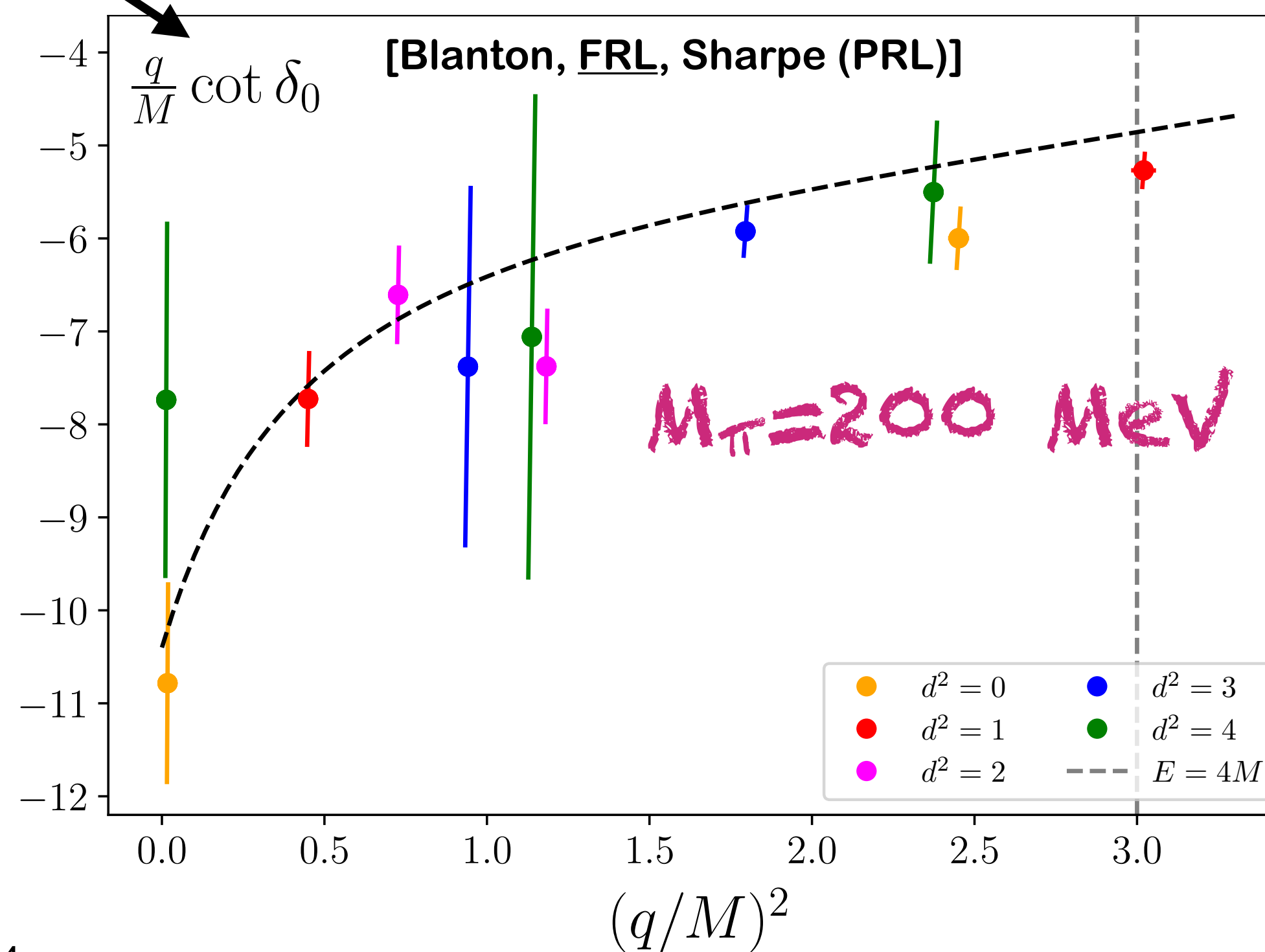
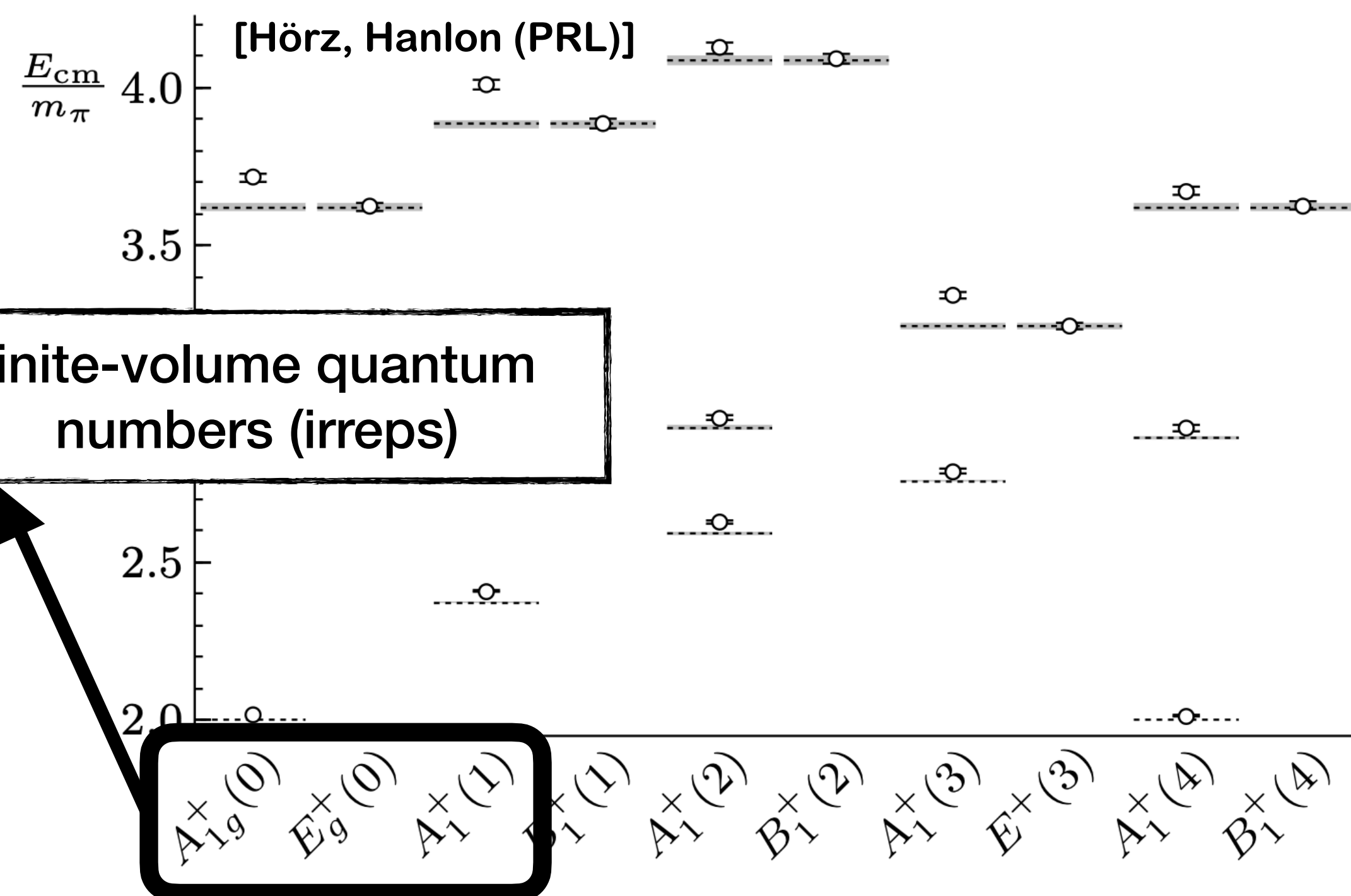
Isospin-2 $\pi\pi$ scattering

Two pions in s-wave

$$\mathcal{K}_2^{s\text{-wave}}(E_n) = \frac{-1}{F_{00}(E_n, \vec{P}, L)}$$

$$\mathcal{K}_2^{s\text{-wave}} \sim \frac{16\pi\sqrt{s}}{q \cot \delta_0}$$

one energy level \rightarrow a phase shift point



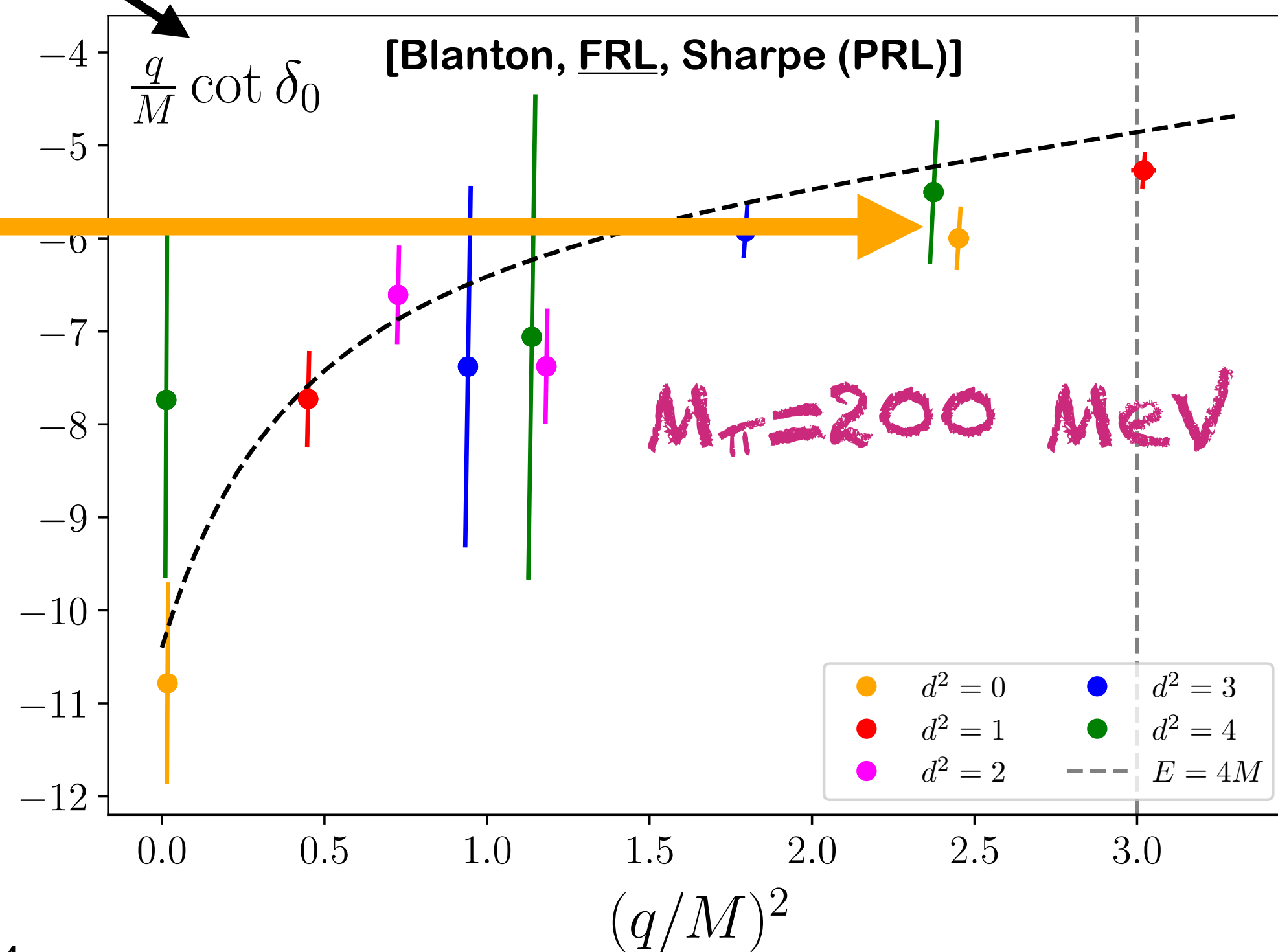
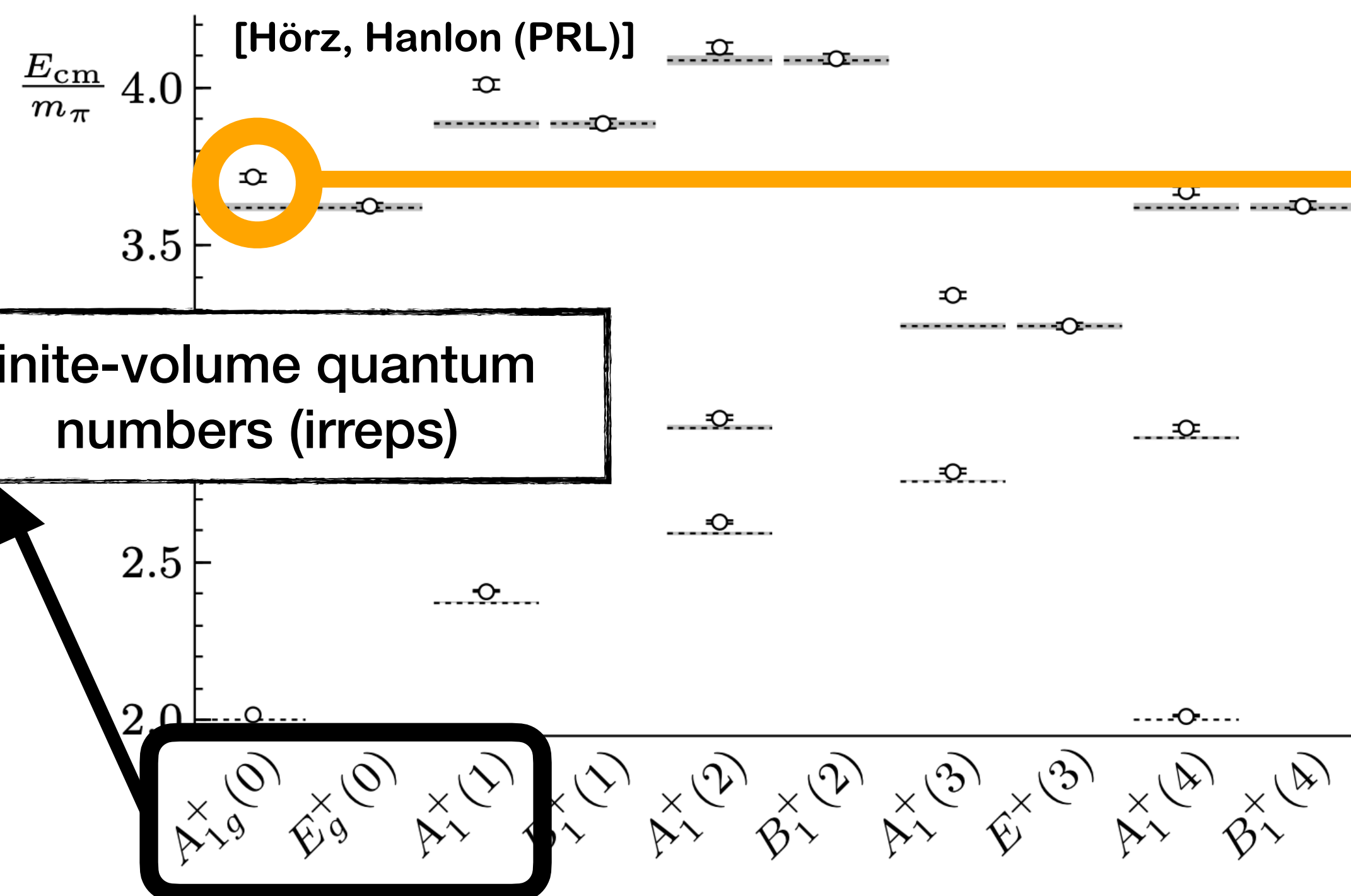
Isospin-2 $\pi\pi$ scattering

Two pions in s-wave

$$\mathcal{K}_2^{s\text{-wave}}(E_n) = \frac{-1}{F_{00}(E_n, \vec{P}, L)}$$

$$\mathcal{K}_2^{s\text{-wave}} \sim \frac{16\pi\sqrt{s}}{q \cot \delta_0}$$

one energy level \longrightarrow a phase shift point



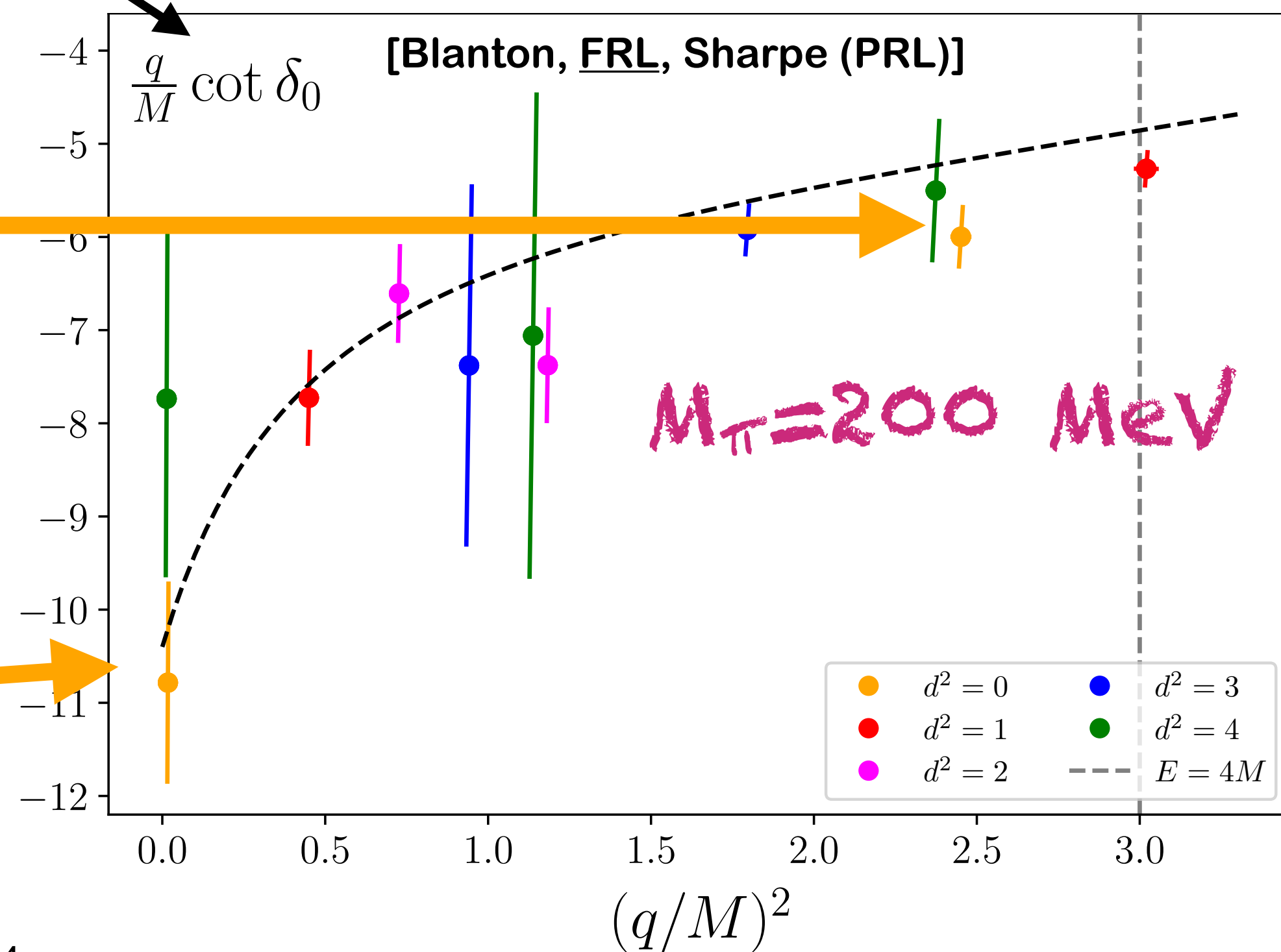
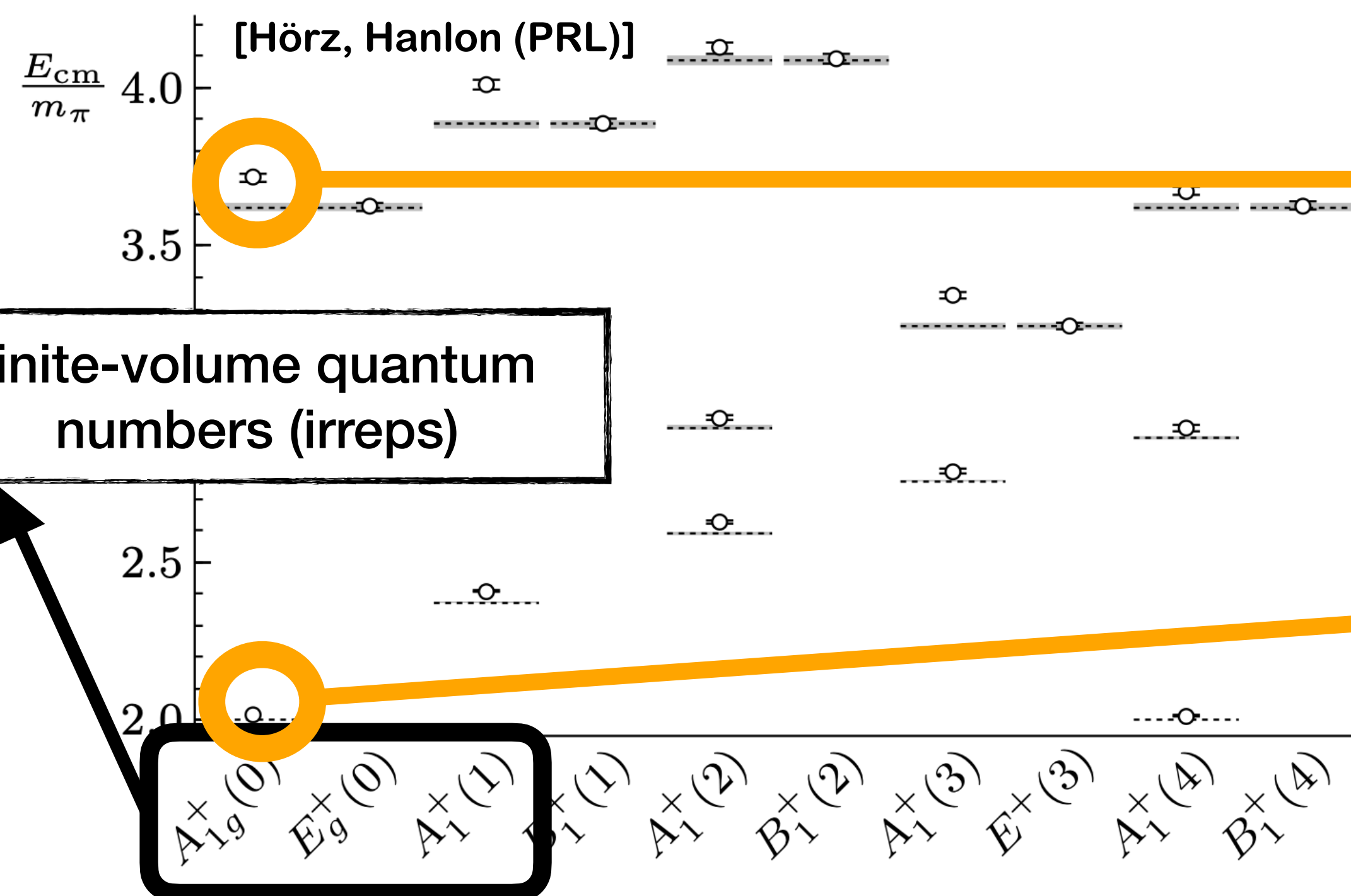
Isospin-2 $\pi\pi$ scattering

Two pions in s-wave

$$\mathcal{K}_2^{s\text{-wave}}(E_n) = \frac{-1}{F_{00}(E_n, \vec{P}, L)}$$

$$\mathcal{K}_2^{s\text{-wave}} \sim \frac{16\pi\sqrt{s}}{q \cot \delta_0}$$

one energy level \longrightarrow a phase shift point



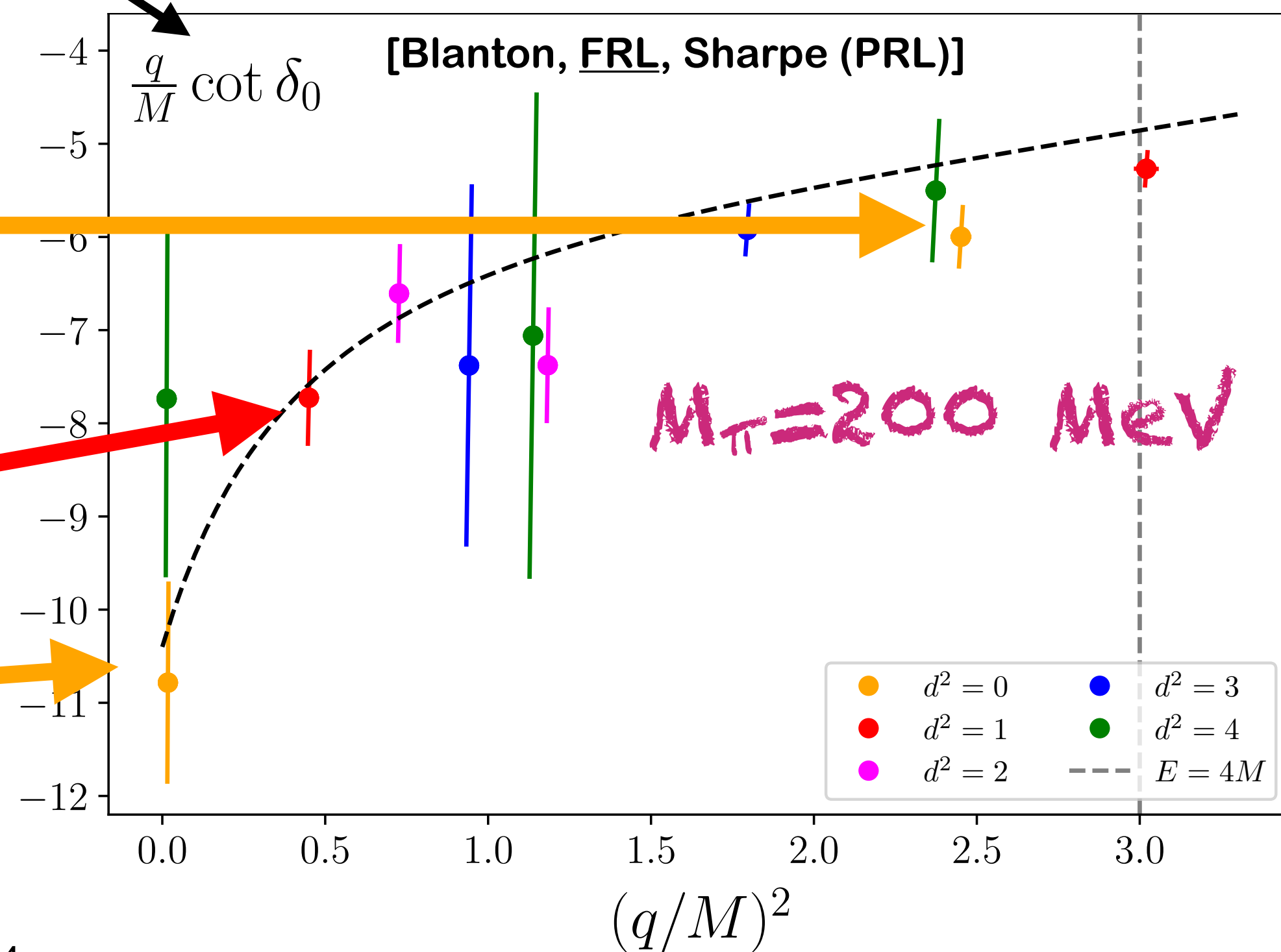
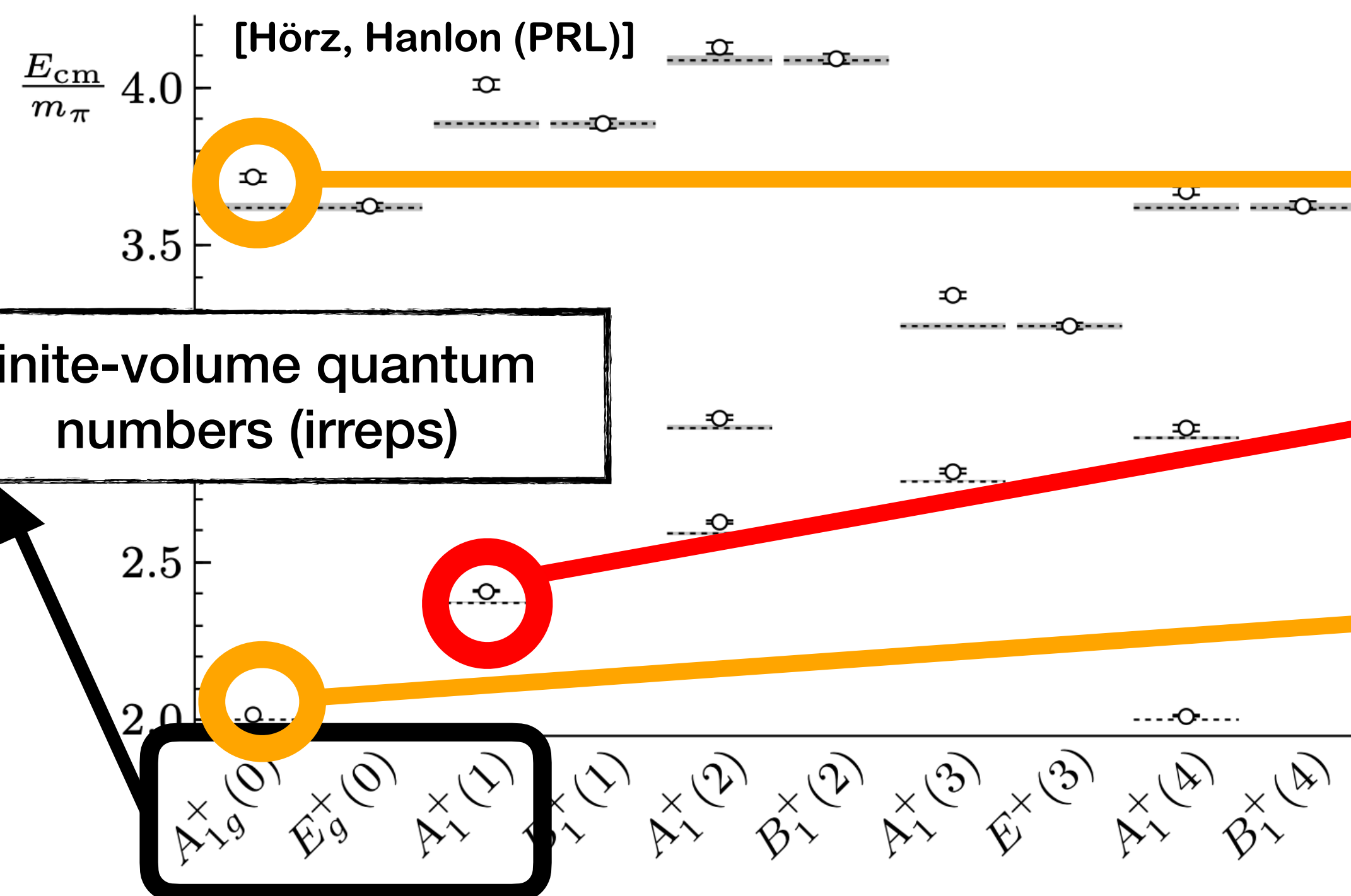
Isospin-2 $\pi\pi$ scattering

Two pions in s-wave

$$\mathcal{K}_2^{s\text{-wave}}(E_n) = \frac{-1}{F_{00}(E_n, \vec{P}, L)}$$

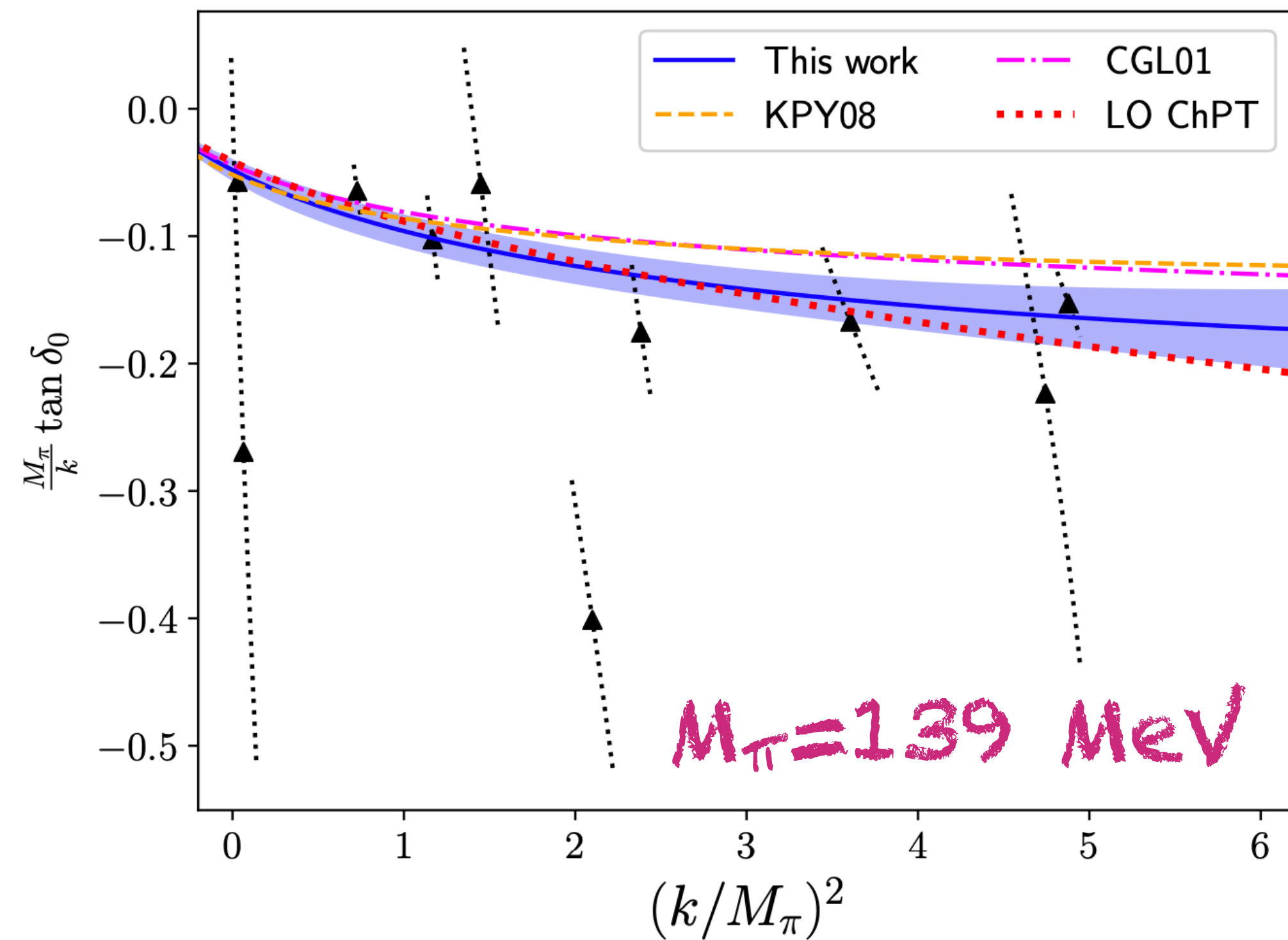
$$\mathcal{K}_2^{s\text{-wave}} \sim \frac{16\pi\sqrt{s}}{q \cot \delta_0}$$

one energy level \longrightarrow a phase shift point



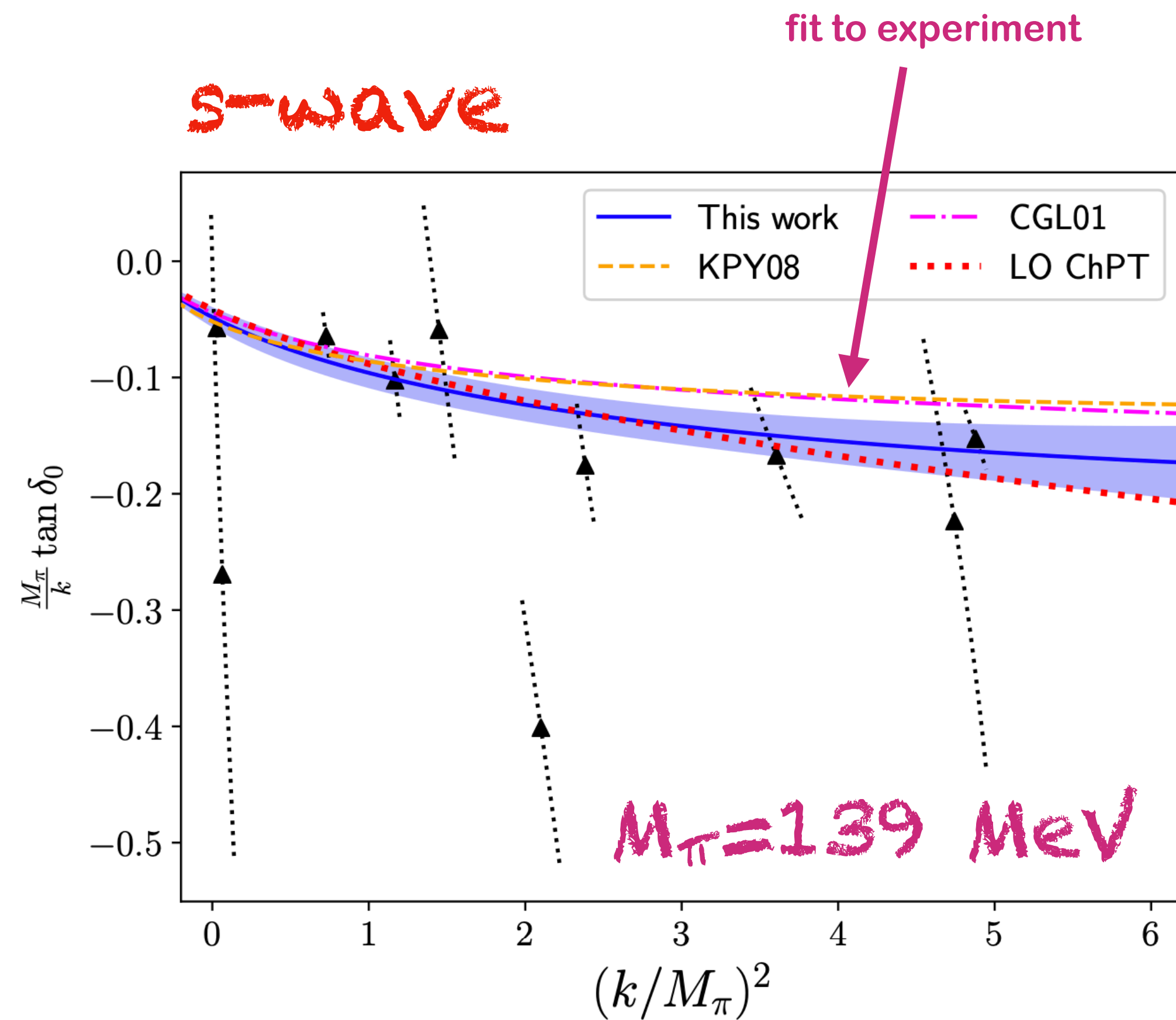
Isospin-2 $\pi\pi$ scattering

S-wave



[Fischer, Kostrzewa, Liu, [FRL](#), Ueding, Urbach (ETMC)]

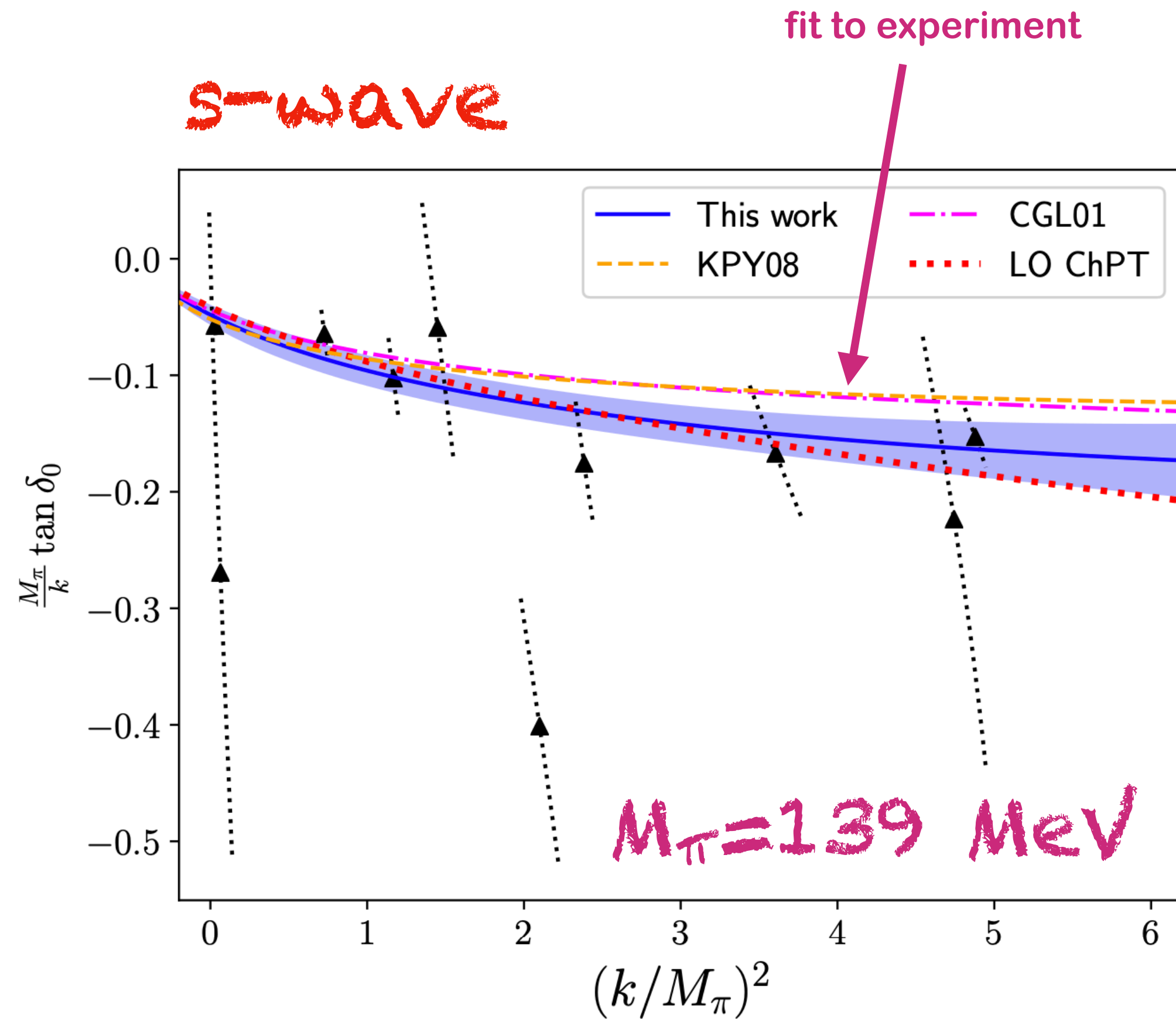
Isospin-2 $\pi\pi$ scattering



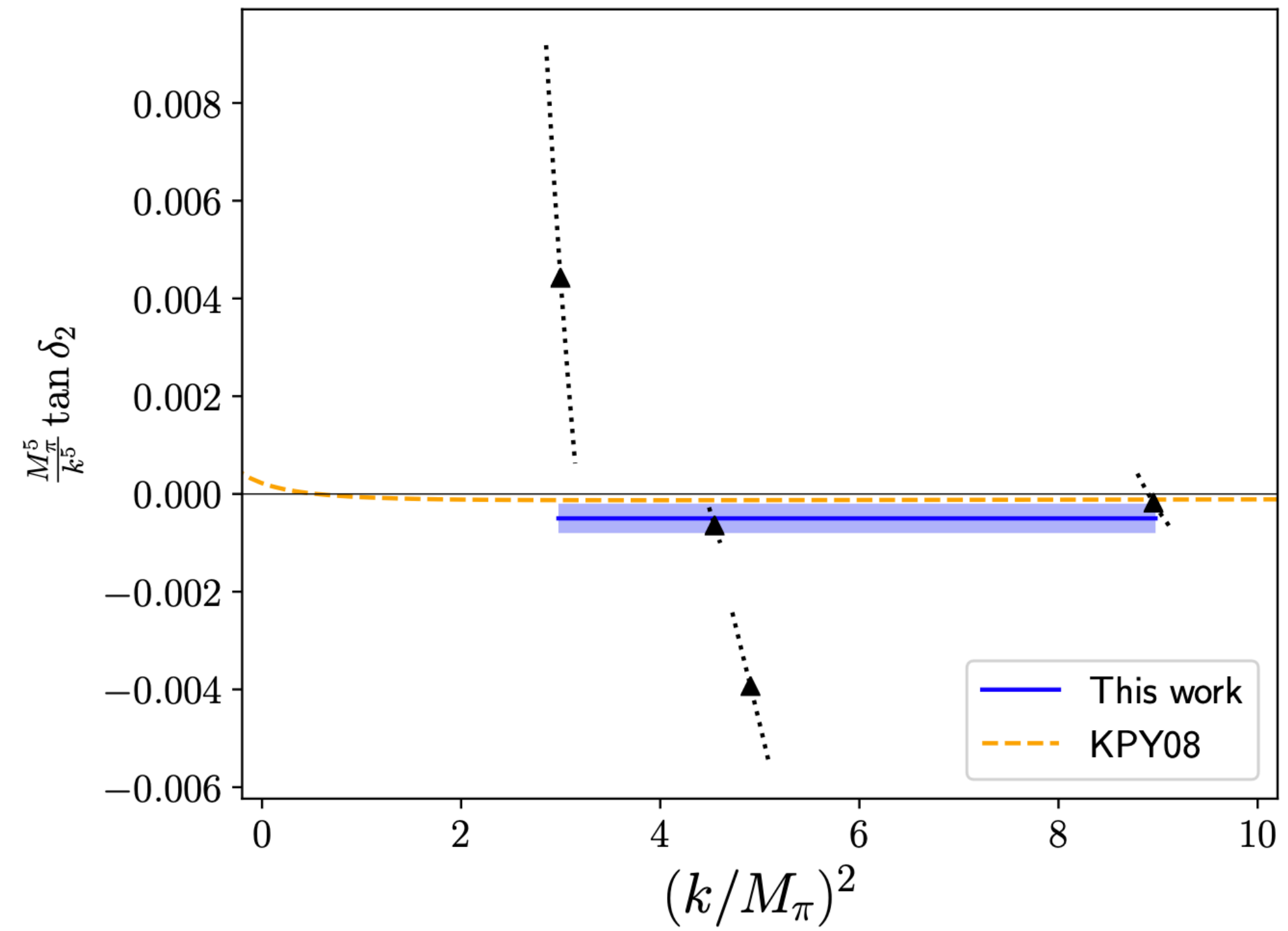
[Fischer, Kostrzewa, Liu, FRL, Ueding, Urbach (ETMC)]

Isospin-2 $\pi\pi$ scattering

s-wave

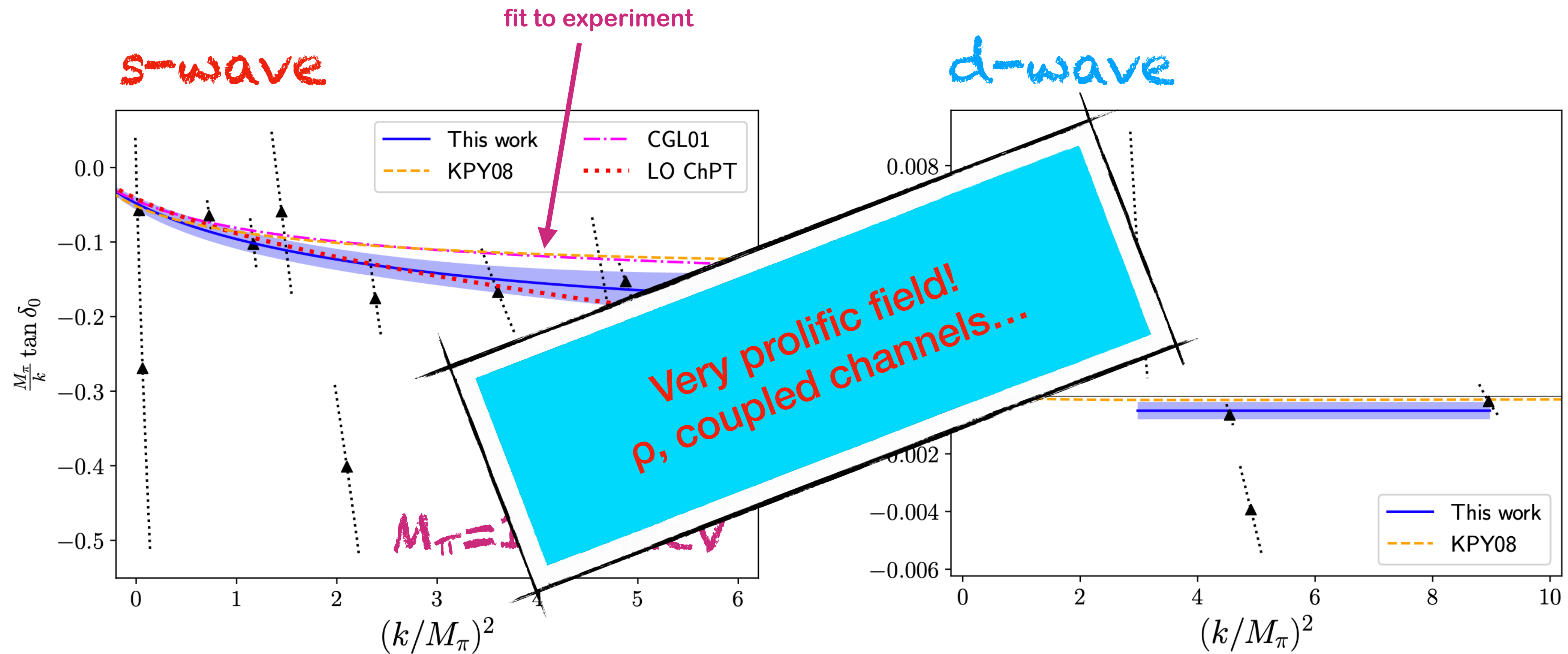


d-wave



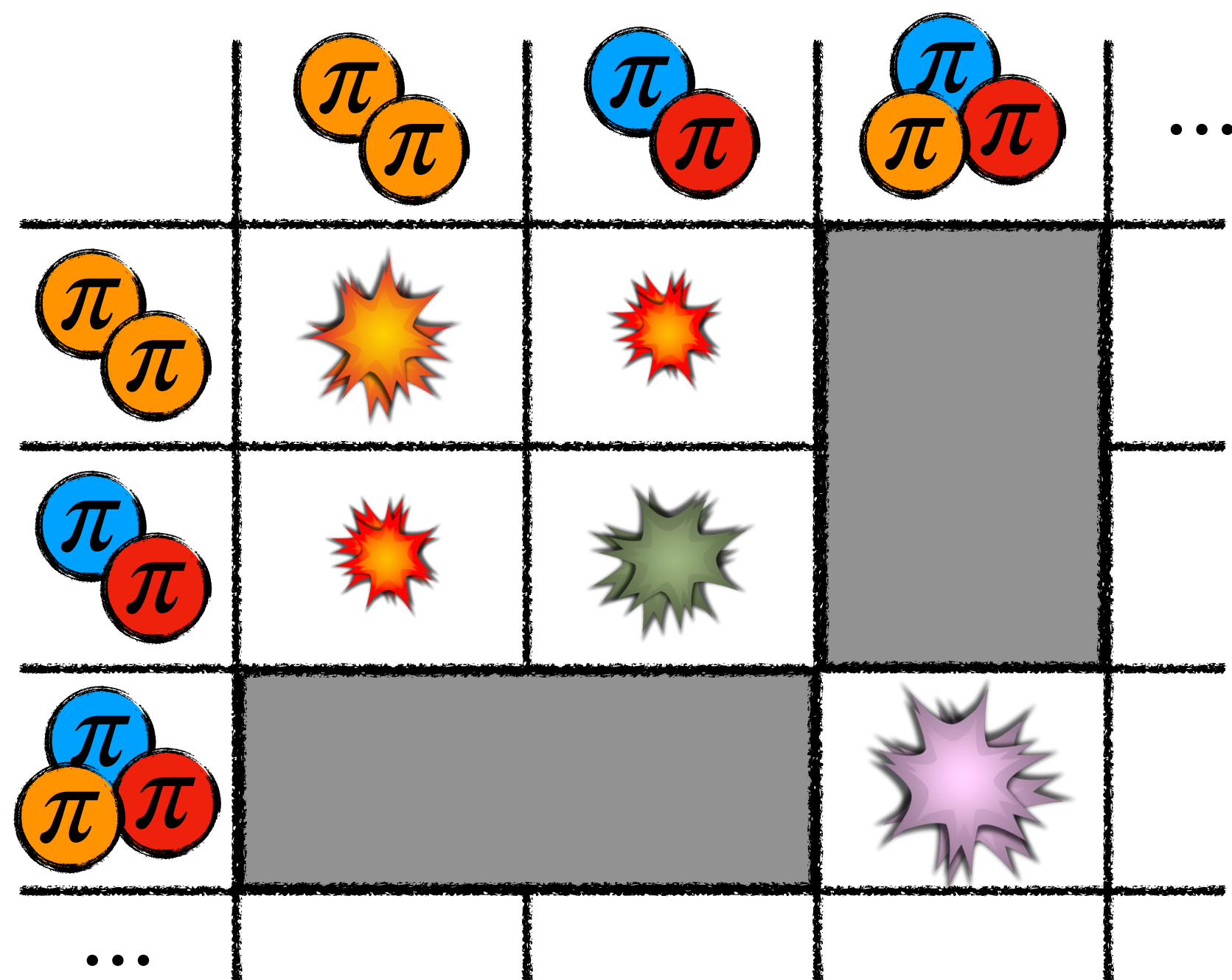
[Fischer, Kostrzewa, Liu, FRL, Ueding, Urbach (ETMC)]

Isospin-2 $\pi\pi$ scattering



[Fischer, Kostrzewa, Liu, FRL, Ueding, Urbach (ETMC)]

Three particles in finite volume



Quantization Condition (I)

Three-particle Quantization Condition
for identical scalars with G-parity

$$\det \left[\mathcal{K}_{df,3}(E) + F_3^{-1}(E, \vec{P}, L) \right] = 0$$

[Hansen, Sharpe]

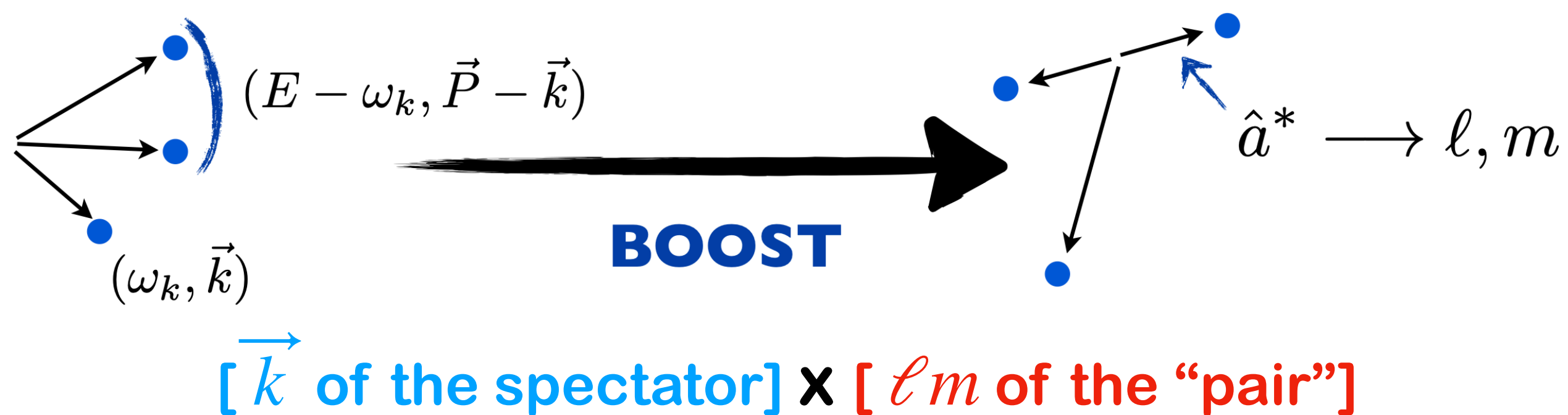
Quantization Condition (I)

Three-particle Quantization Condition
for identical scalars with G-parity

$$\det \left[\mathcal{K}_{df,3}(E) + F_3^{-1}(E, \vec{P}, L) \right] = 0$$

[Hansen, Sharpe]

Matrix indices are more complicated:



Quantization Condition (I)

Three-particle Quantization Condition
for identical scalars with G-parity

$$\det \left[\mathcal{K}_{df,3}(E) + F_3^{-1}(E, \vec{P}, L) \right] = 0$$

[Hansen, Sharpe]

Matrix indices are more complicated:



$[\vec{k}$ of the spectator] \times [ℓm of the “pair”]

Truncation:
neglect higher ℓ + cutoff function

Quantization Condition (I)

Three-particle Quantization Condition
for identical scalars with G-parity

$$\det \left[\mathcal{K}_{df,3}(E) + F_3^{-1}(E, \vec{P}, L) \right] = 0$$

[Hansen, Sharpe]

- $\mathcal{K}_{df,3}$ is real, divergence-free. It is an intermediate **cutoff-dependent** quantity with the symmetries of the physical amplitude

Matrix indices are more complicated:



$[\vec{k}$ of the spectator] \times [ℓm of the "pair"] \rightarrow

Truncation:
neglect higher ℓ + cutoff function

Quantization Condition (I)

Three-particle Quantization Condition
for identical scalars with G-parity

$$\det \left[\mathcal{K}_{df,3}(E) + F_3^{-1}(E, \vec{P}, L) \right] = 0$$

[Hansen, Sharpe]

- $\mathcal{K}_{df,3}$ is real, divergence-free. It is an intermediate **cutoff-dependent** quantity with the symmetries of the physical amplitude
- F_3 depends on **kinematical functions** and on the **two-to-two scattering** amplitude

Matrix indices are more complicated:



$[\vec{k}$ of the spectator] \times [ℓm of the "pair"] \rightarrow

Truncation:
neglect higher ℓ + cutoff function

Quantization Condition (I)

Three-particle Quantization Condition
for identical scalars with G-parity

$$\det \left[\mathcal{K}_{df,3}(E) + F_3^{-1}(E, \vec{P}, L) \right] = 0$$

[Hansen, Sharpe]

- $\mathcal{K}_{df,3}$ is real, divergence-free. It is an intermediate **cutoff-dependent** quantity with the symmetries of the physical amplitude
- F_3 depends on **kinematical functions** and on the **two-to-two scattering** amplitude

Matrix indices are more complicated:



$[\vec{k}$ of the spectator] \times [ℓm of the “pair”]



Recovering the physical
amplitude requires a further step

Truncation:
neglect higher ℓ + cutoff function

Quantization Condition (II)

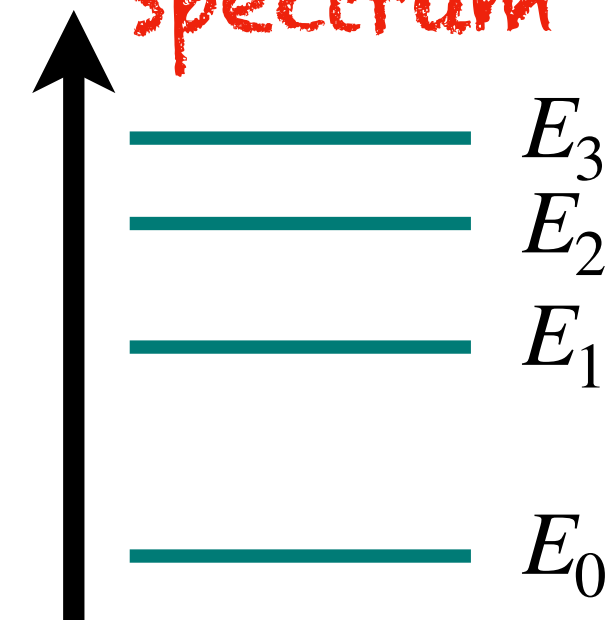
1. \mathcal{K}_2 and $\mathcal{K}_{df,3}$ parametrize interactions.
They can be obtained from the spectrum

[Hansen, Sharpe]

Alternative approaches:
[Mai, Döring], [Hammer, et al.]

Quantization Condition (II)

2π and 3π
Spectrum



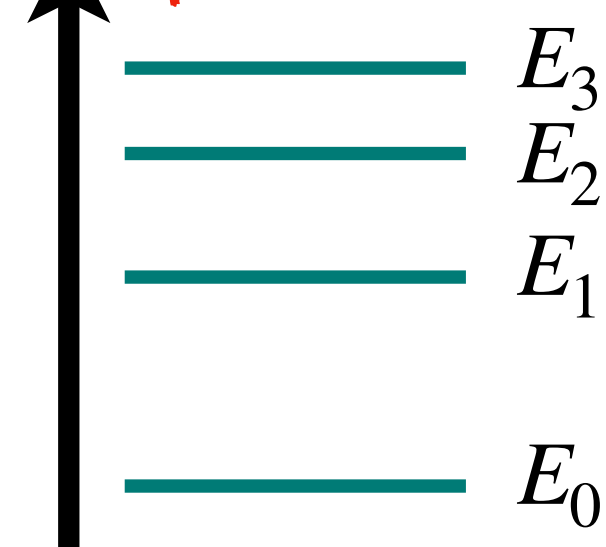
1. \mathcal{K}_2 and $\mathcal{K}_{df,3}$ parametrize interactions.
They can be obtained from the spectrum

[Hansen, Sharpe]

Alternative approaches:
[Mai, Döring], [Hammer, et al.]

Quantization Condition (II)

2π and 3π
Spectrum



$$\det [\mathcal{K}_2 + F_2^{-1}] = 0$$

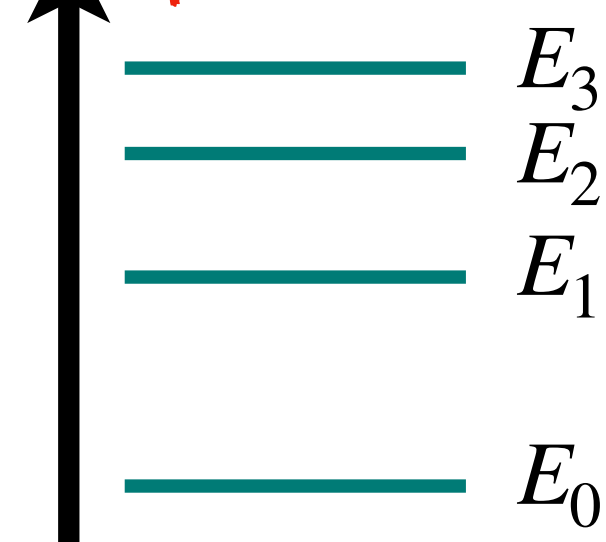
1. \mathcal{K}_2 and $\mathcal{K}_{df,3}$ parametrize interactions. They can be obtained from the spectrum

[Hansen, Sharpe]

Alternative approaches:
[Mai, Döring], [Hammer, et al.]

Quantization Condition (II)

2π and 3π
Spectrum



$$\det [\mathcal{K}_2 + F_2^{-1}] = 0$$

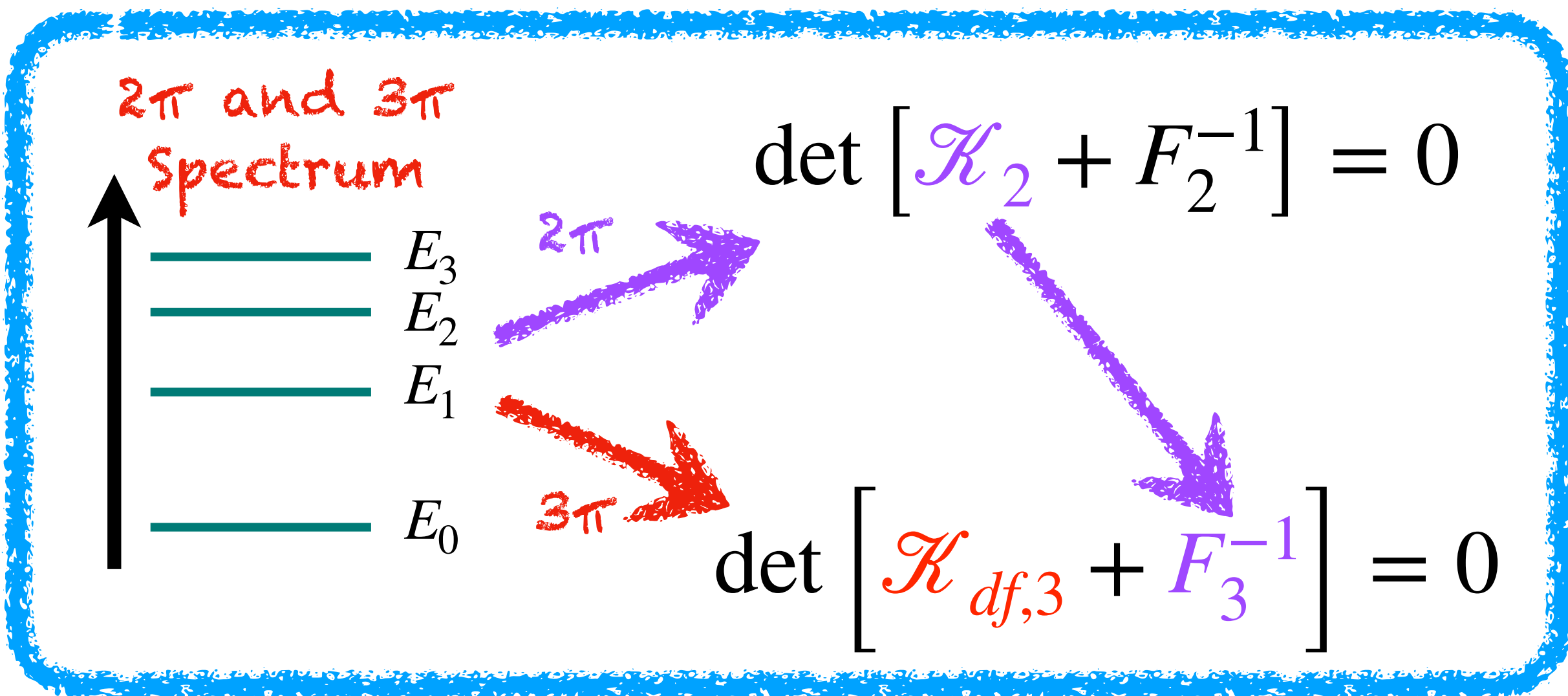
$$\det [\mathcal{K}_{df,3} + F_3^{-1}] = 0$$

1. \mathcal{K}_2 and $\mathcal{K}_{df,3}$ parametrize interactions. They can be obtained from the spectrum

[Hansen, Sharpe]

Alternative approaches:
[Mai, Döring], [Hammer, et al.]

Quantization Condition (II)



1. \mathcal{K}_2 and $\mathcal{K}_{df,3}$ parametrize interactions. They can be obtained from the spectrum

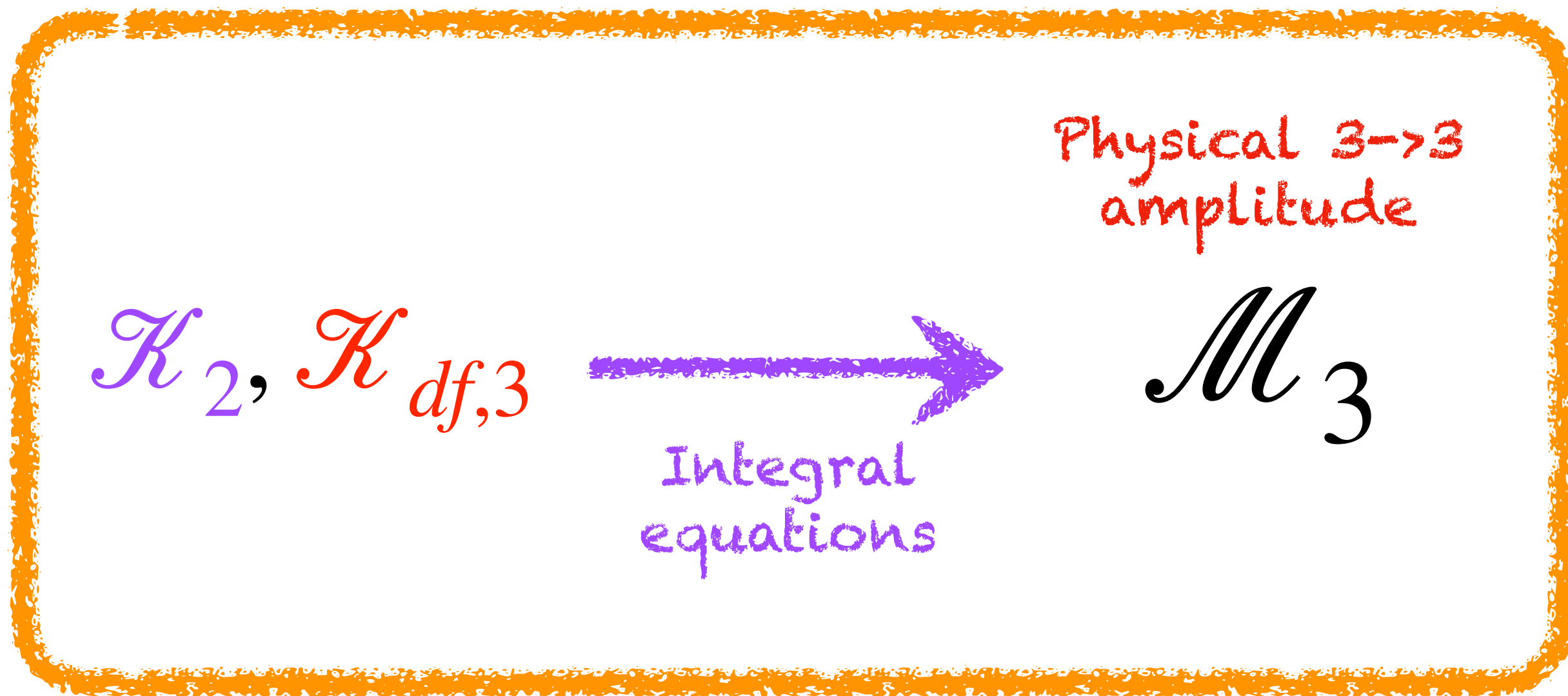
[Hansen, Sharpe]

Alternative approaches:
[Mai, Döring], [Hammer, et al.]

2. Solve integral equations to obtain the physical three-to-three amplitude

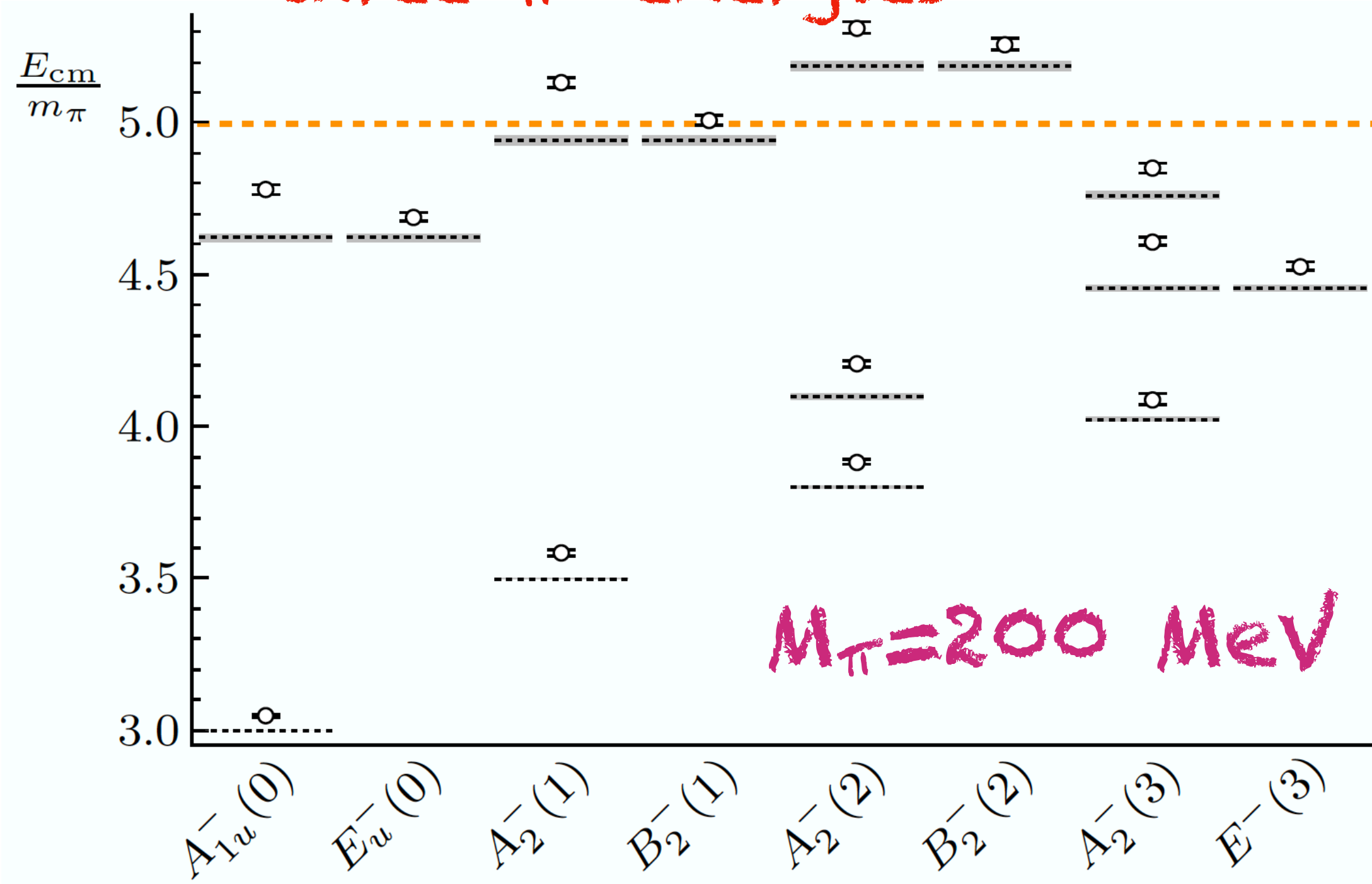
Derived by [Hansen, Sharpe]

Solved in [Briceño et al], [Hansen et al.], [Jackura et al.]



two & three-pion spectrum

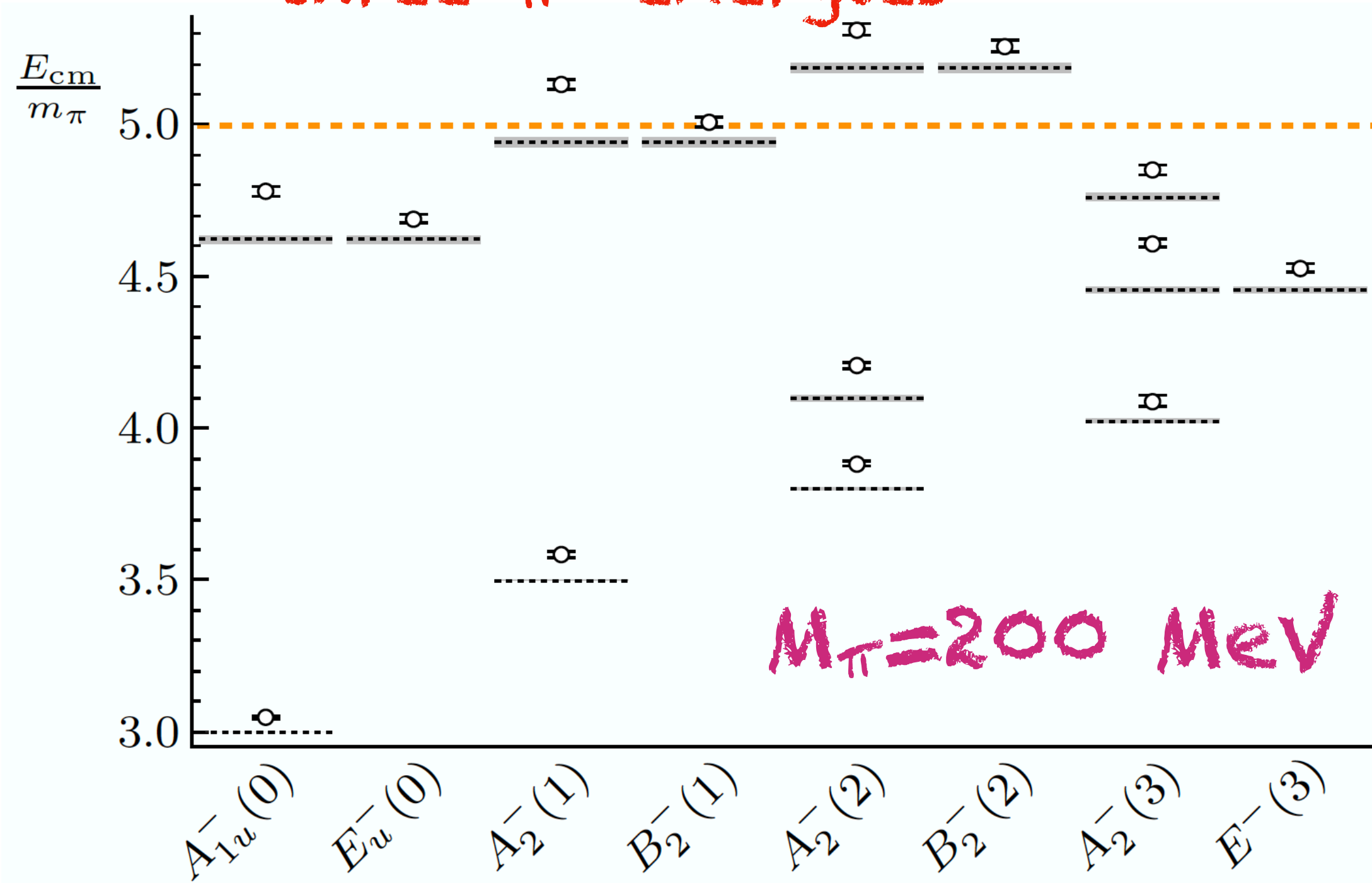
three- π^+ energies



[Hörz, Hanlon (PRL)]

two & three-pion spectrum

three- π^+ energies



$M_\pi = 200$ MeV

[Hörz, Hanlon (PRL)]

$I = 3$ three-pion scattering amplitude from lattice QCD

Tyler D. Blanton,^{1,*} Fernando Romero-López,^{2,†} and Stephen R. Sharpe^{1,‡}

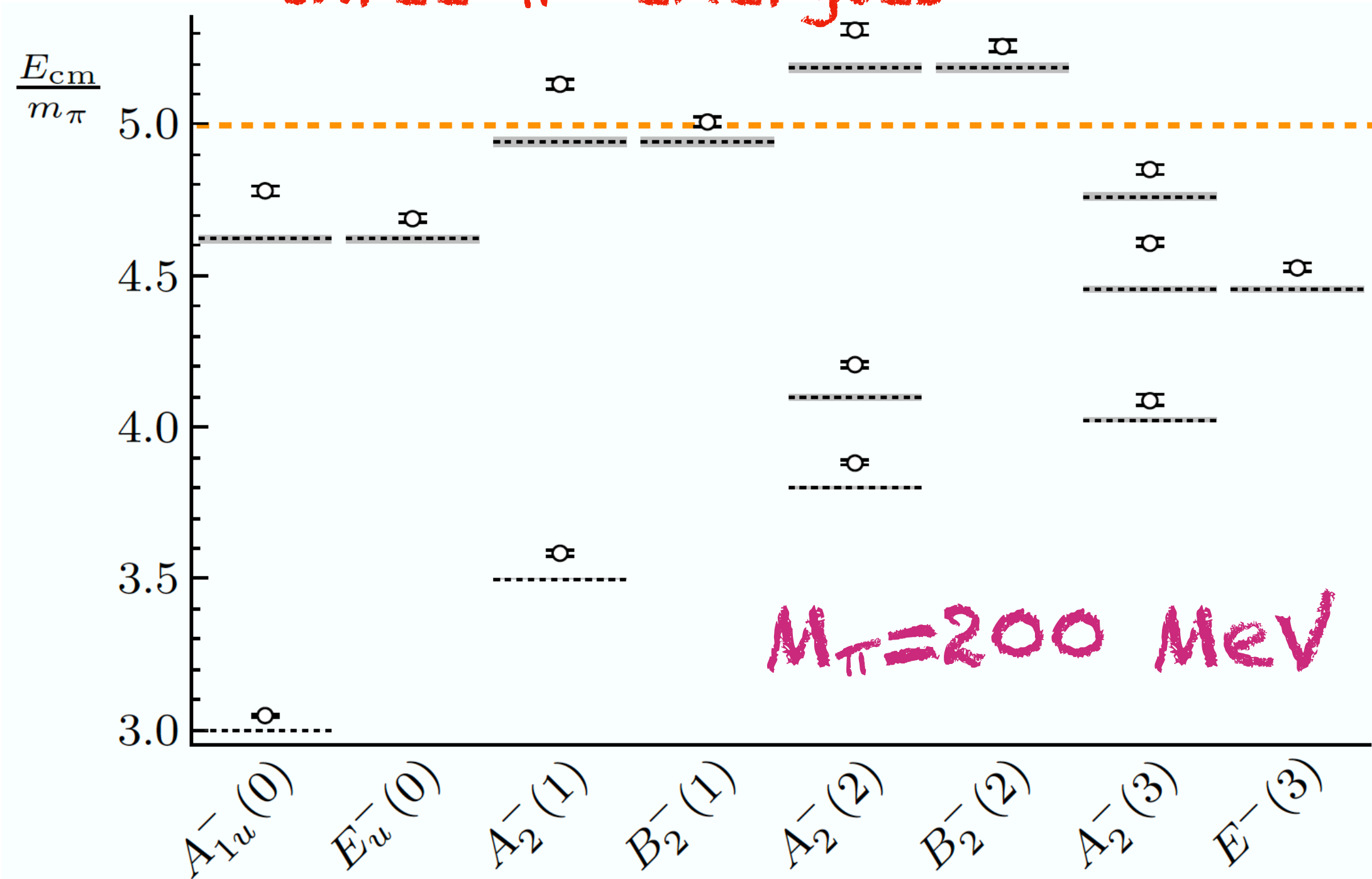
¹Physics Department, University of Washington, Seattle, WA 98195-1560, USA

²Instituto de Física Corpuscular, Universitat de València and CSIC, 46980 Paterna, Spain

(Dated: February 4, 2020)

Two & three-pion spectrum

three- π^+ energies



[Hörz, Hanlon (PRL)]

$I = 3$ three-pion scattering amplitude from lattice QCD

Tyler D. Blanton,^{1,*} Fernando Romero-López,^{2,†} and Stephen R. Sharpe^{1,‡}

¹Physics Department, University of Washington, Seattle, WA 98195-1560, USA

²Instituto de Física Corpuscular, Universitat de València and CSIC, 46980 Paterna, Spain

(Dated: February 4, 2020)

First full analysis of the finite-volume spectrum of $2\pi^+$ and $3\pi^+$!

Fit results

○ Parametrize $\mathcal{K}_{df,3}$ including only s-wave interactions:

$$\frac{q}{M} \cot \delta_0 = \frac{\sqrt{s}M}{s - z_2^2} (B_0 + B_1 q^2 + \dots)$$

$$\mathcal{K}_{df,3} = \mathcal{K}_{df,3}^{iso,0} + \mathcal{K}_{df,3}^{iso,1} \left(\frac{s - 9M^2}{9M^2} \right)$$

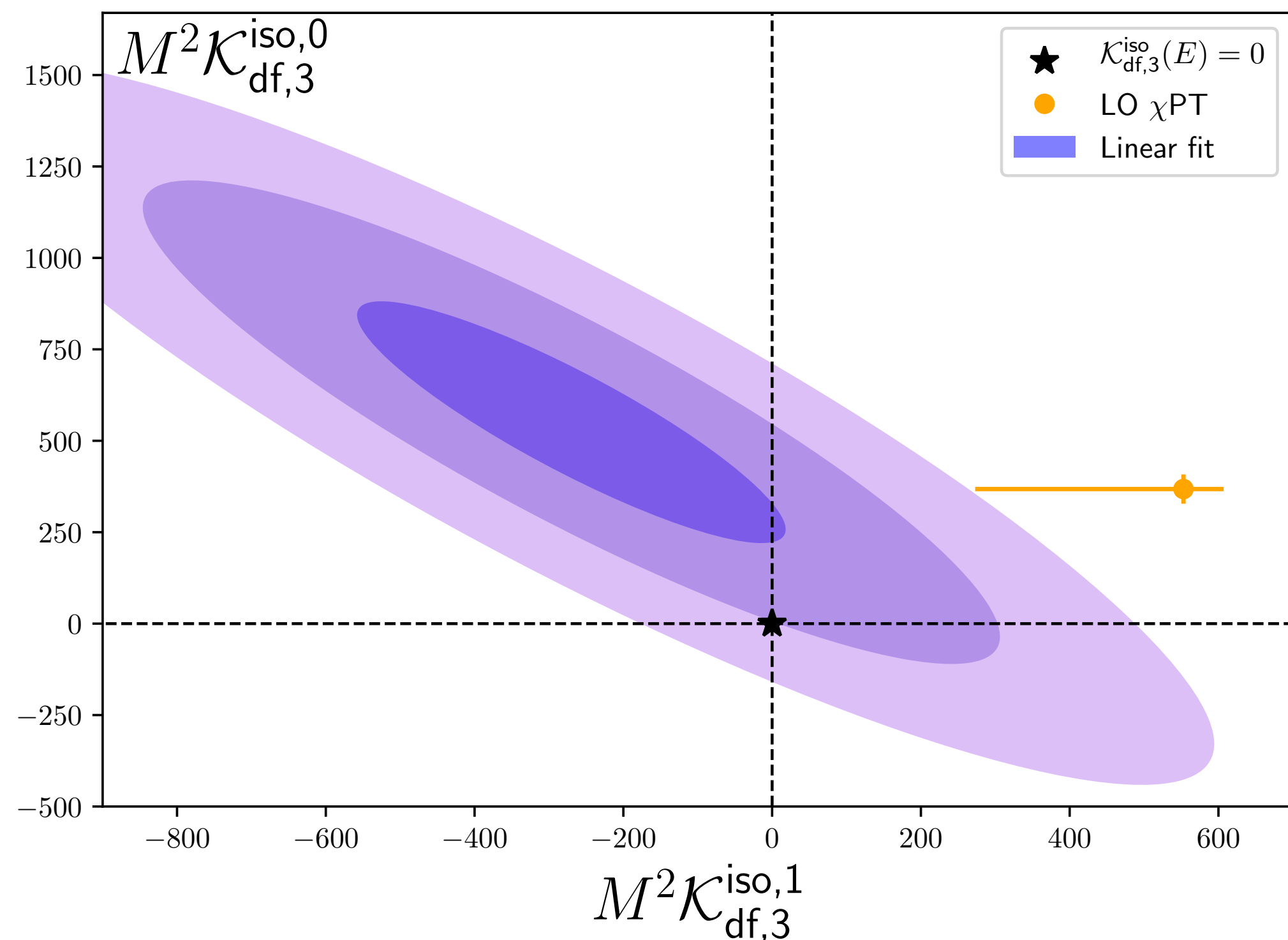
Fit results

Parametrize $\mathcal{K}_{df,3}$ including only s-wave interactions:

$$\frac{q}{M} \cot \delta_0 = \frac{\sqrt{s}M}{s - z_2^2} (B_0 + B_1 q^2 + \dots)$$

Fit	B_0	B_1	z_2^2/M^2	$M^2\mathcal{K}_{df,3}^{iso,0}$	$M^2\mathcal{K}_{df,3}^{iso,1}$	χ^2/dof	Ma_0	M^2ra_0
5	-11.1(7)	-2.4(3)	1 (fixed)	550(330)	-280(290)	26.04/(22-4)	0.090(5)	2.57(8)

$$\mathcal{K}_{df,3} = \mathcal{K}_{df,3}^{iso,0} + \mathcal{K}_{df,3}^{iso,1} \left(\frac{s - 9M^2}{9M^2} \right)$$



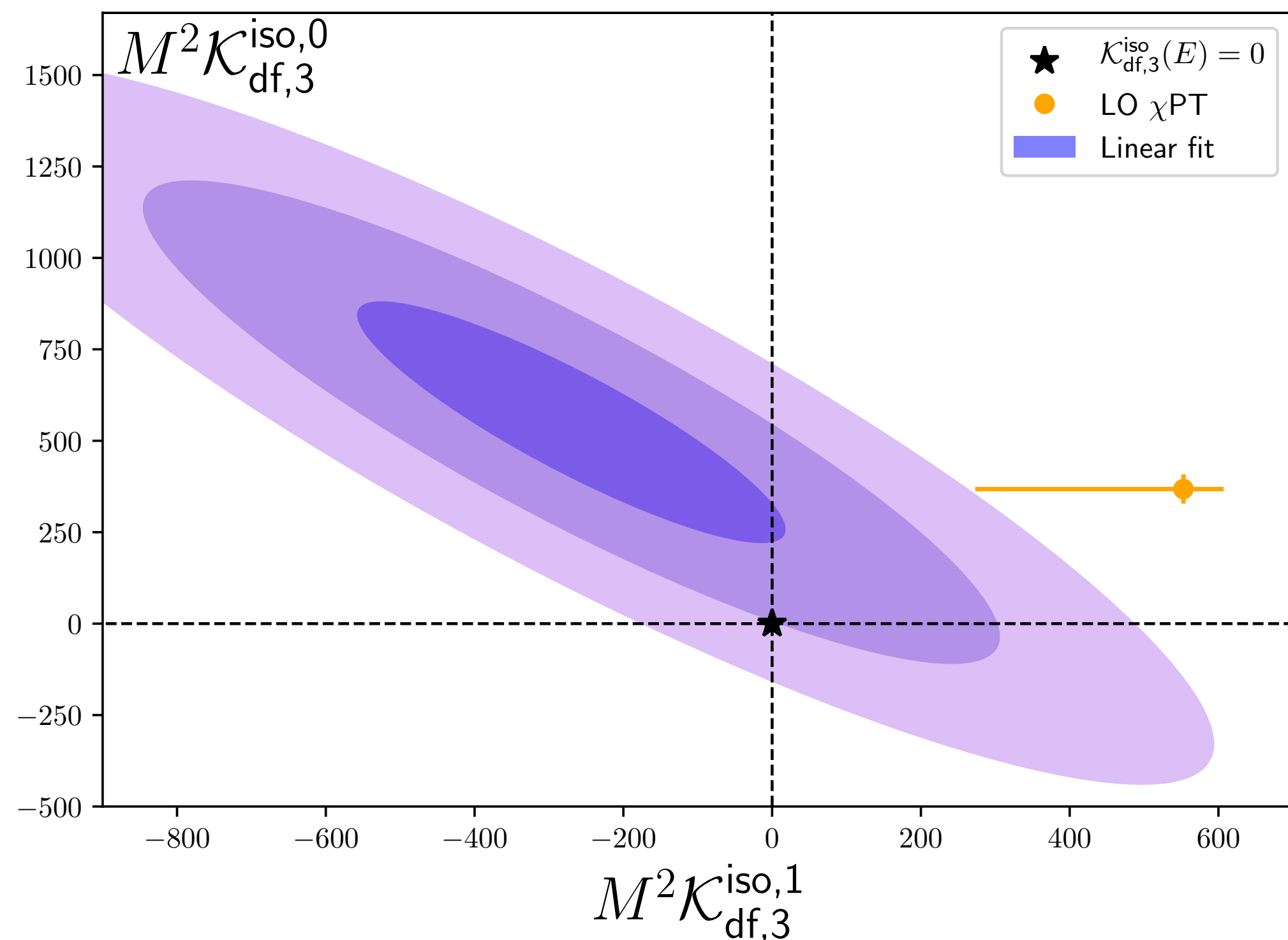
Fit results

Parametrize $\mathcal{K}_{df,3}$ including only s-wave interactions:

$$\frac{q}{M} \cot \delta_0 = \frac{\sqrt{s}M}{s - z_2^2} (B_0 + B_1 q^2 + \dots)$$

Fit	B_0	B_1	z_2^2/M^2	$M^2\mathcal{K}_{df,3}^{iso,0}$	$M^2\mathcal{K}_{df,3}^{iso,1}$	χ^2/dof	Ma_0	M^2ra_0
5	-11.1(7)	-2.4(3)	1 (fixed)	550(330)	-280(290)	26.04/(22-4)	0.090(5)	2.57(8)

$$\mathcal{K}_{df,3} = \mathcal{K}_{df,3}^{iso,0} + \mathcal{K}_{df,3}^{iso,1} \left(\frac{s - 9M^2}{9M^2} \right)$$



1. 2σ evidence for $\mathcal{K}_{df,3} \neq 0$.

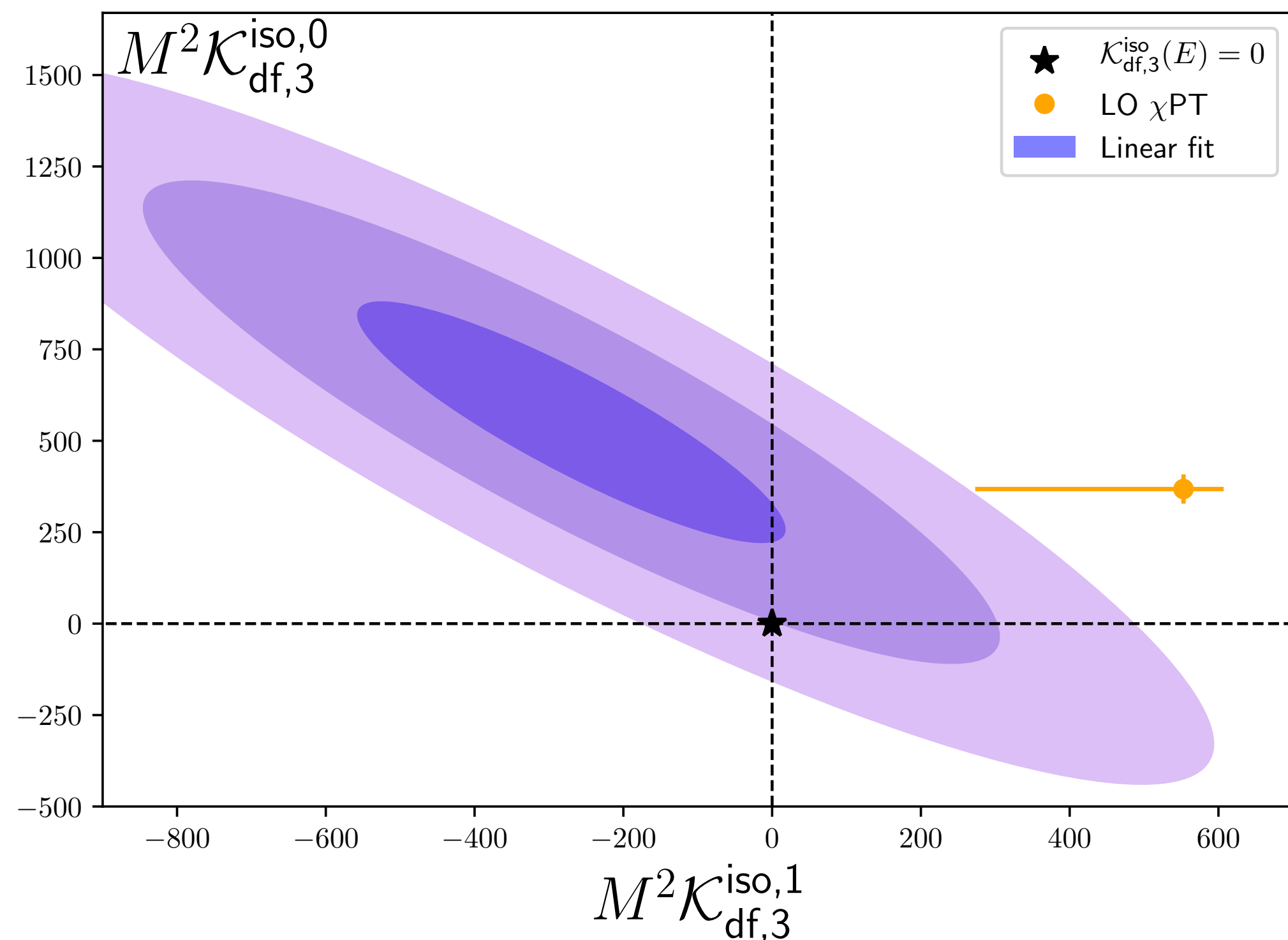
Fit results

Parametrize $\mathcal{K}_{df,3}$ including only s-wave interactions:

$$\frac{q}{M} \cot \delta_0 = \frac{\sqrt{s}M}{s - z_2^2} (B_0 + B_1 q^2 + \dots)$$

Fit	B_0	B_1	z_2^2/M^2	$M^2\mathcal{K}_{df,3}^{iso,0}$	$M^2\mathcal{K}_{df,3}^{iso,1}$	χ^2/dof	Ma_0	M^2ra_0
5	-11.1(7)	-2.4(3)	1 (fixed)	550(330)	-280(290)	26.04/(22-4)	0.090(5)	2.57(8)

$$\mathcal{K}_{df,3} = \mathcal{K}_{df,3}^{iso,0} + \mathcal{K}_{df,3}^{iso,1} \left(\frac{s - 9M^2}{9M^2} \right)$$



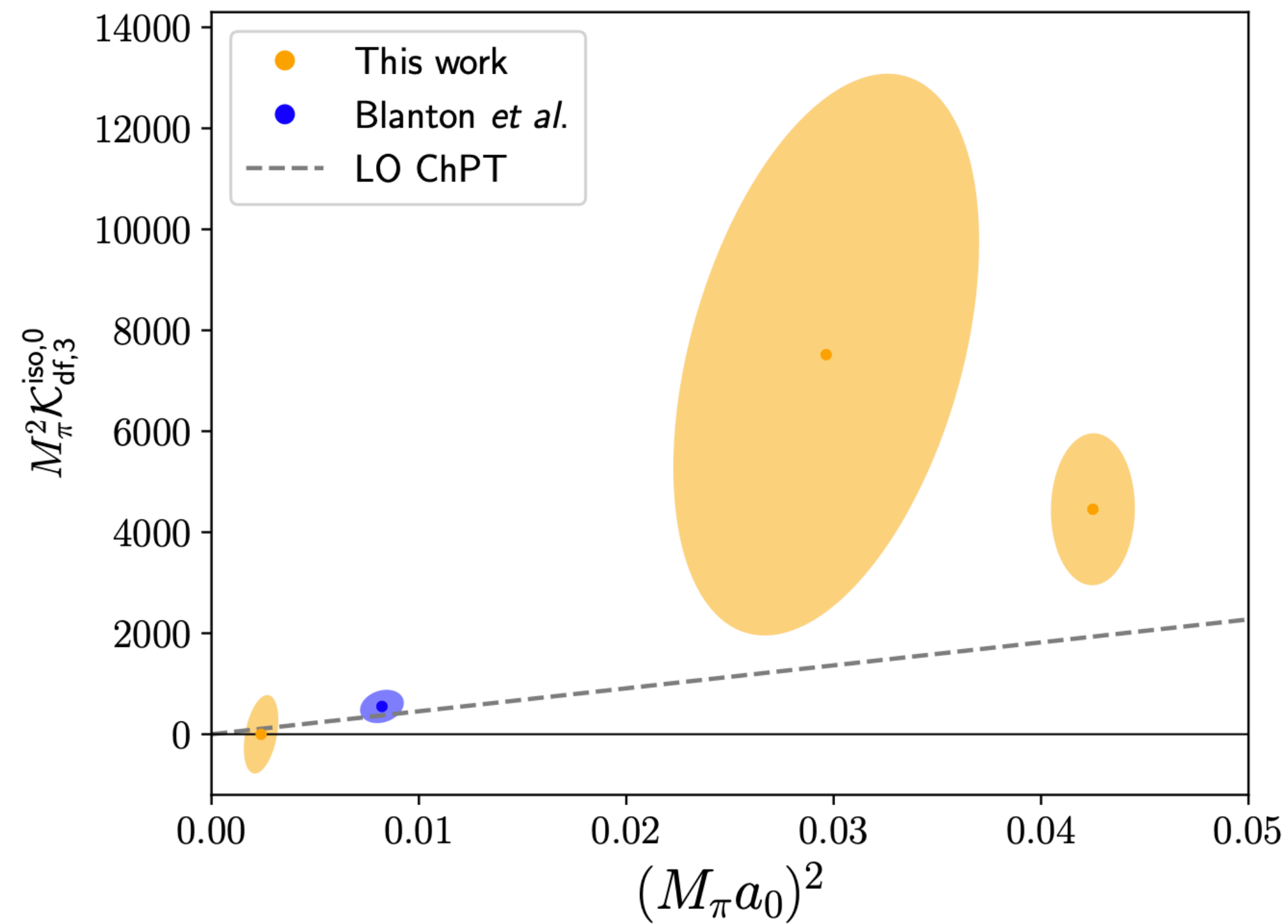
1. 2σ evidence for $\mathcal{K}_{df,3} \neq 0$.
2. Some tension with ChPT.

Chiral dependence

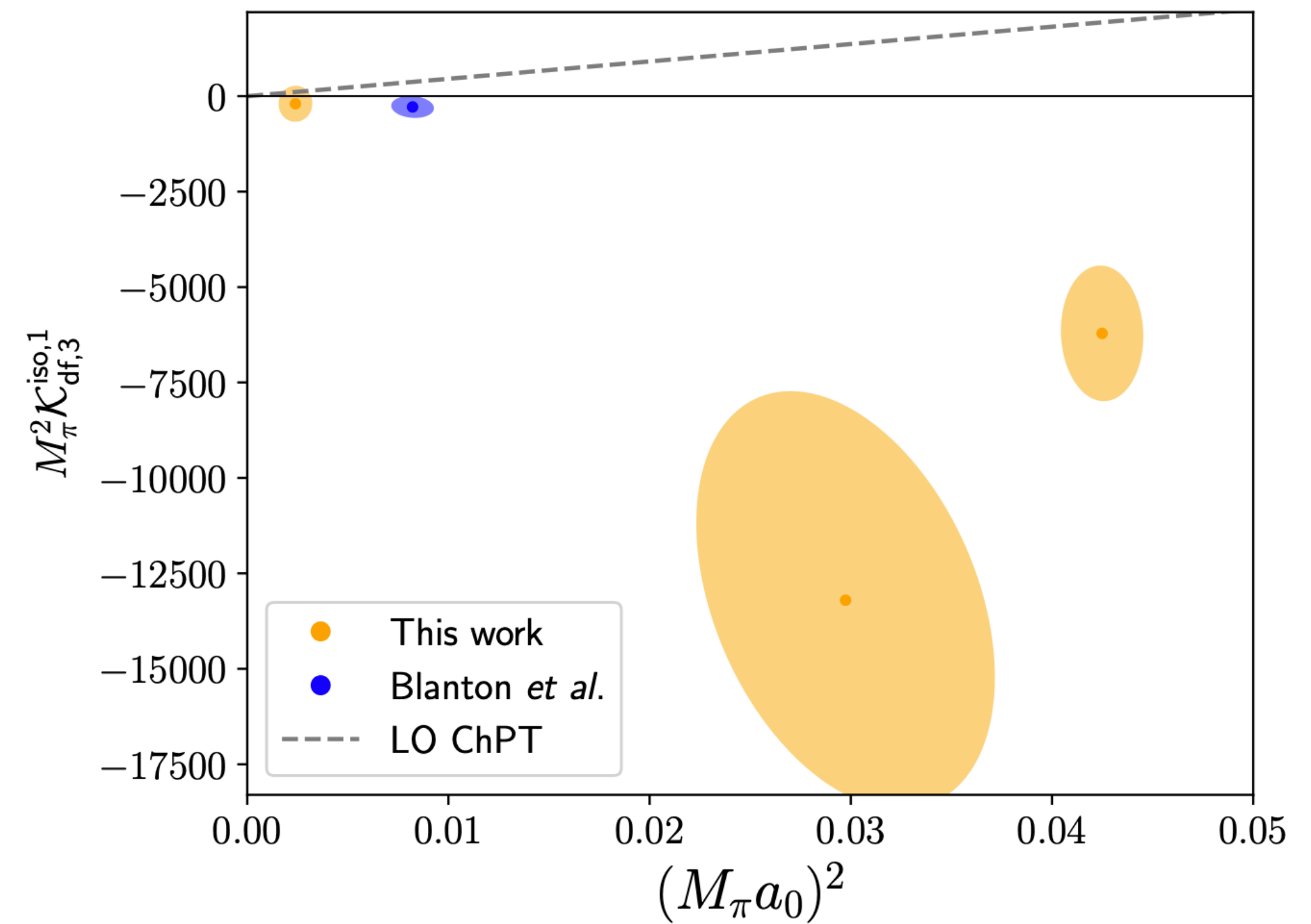
On a later article, the chiral dependence of $\mathcal{K}_{df,3}$ has been studied, including physical pions.

[Fischer, Kostrzewa, Liu, [FRL](#), Ueding, Urbach (ETMC)]

$$\mathcal{K}_{df,3} = \mathcal{K}_{df,3}^{iso,0} + \mathcal{K}_{df,3}^{iso,1} \left(\frac{s - 9M^2}{9M^2} \right)$$



(a) $\mathcal{K}_{df,3}^{iso,0}$



(b) $\mathcal{K}_{df,3}^{iso,1}$

see also other studies:
[Mai et al., Culver et al.]

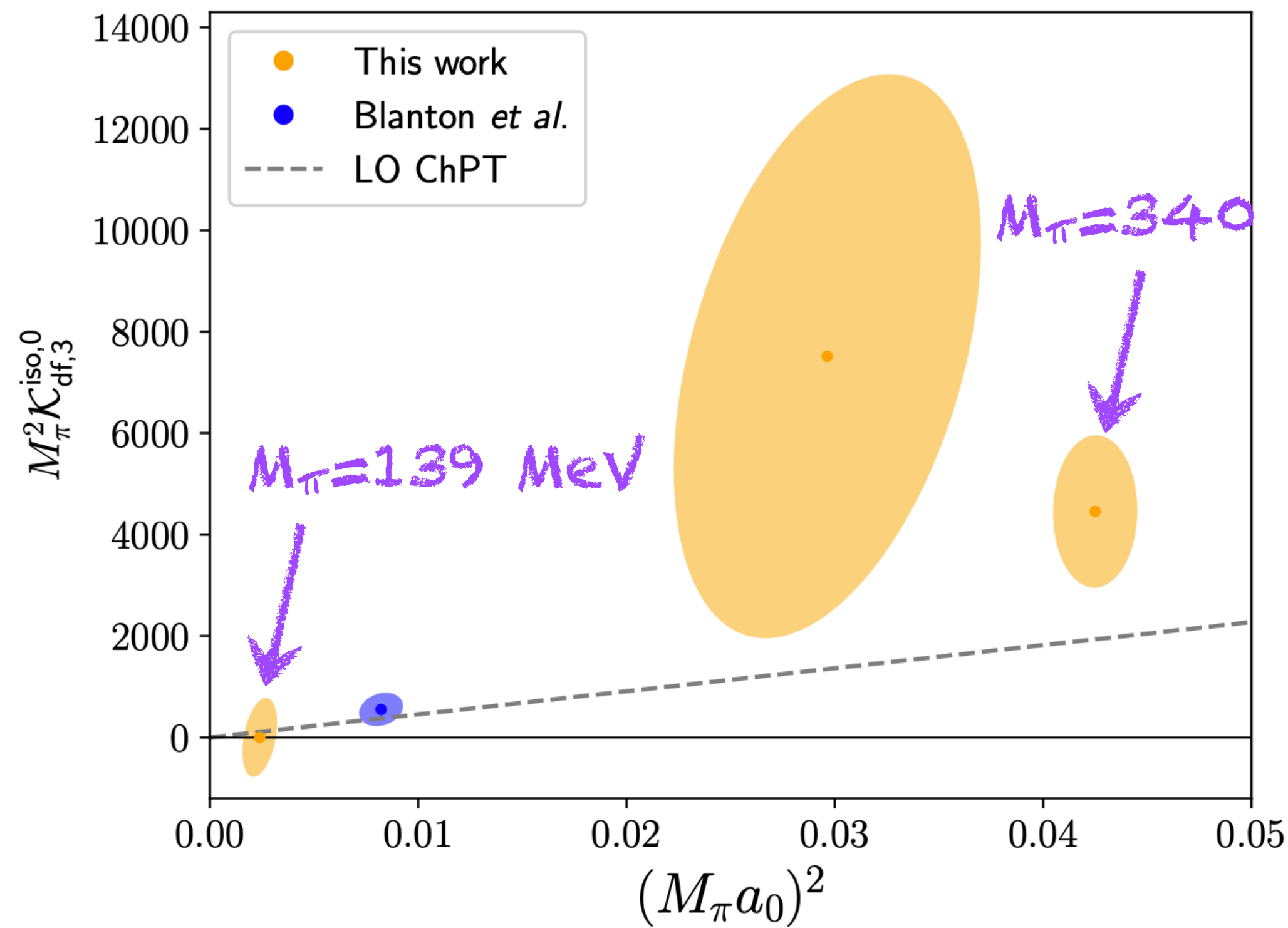
Chiral dependence

On a later article, the chiral dependence of $\mathcal{K}_{df,3}$ has been studied, including physical pions.

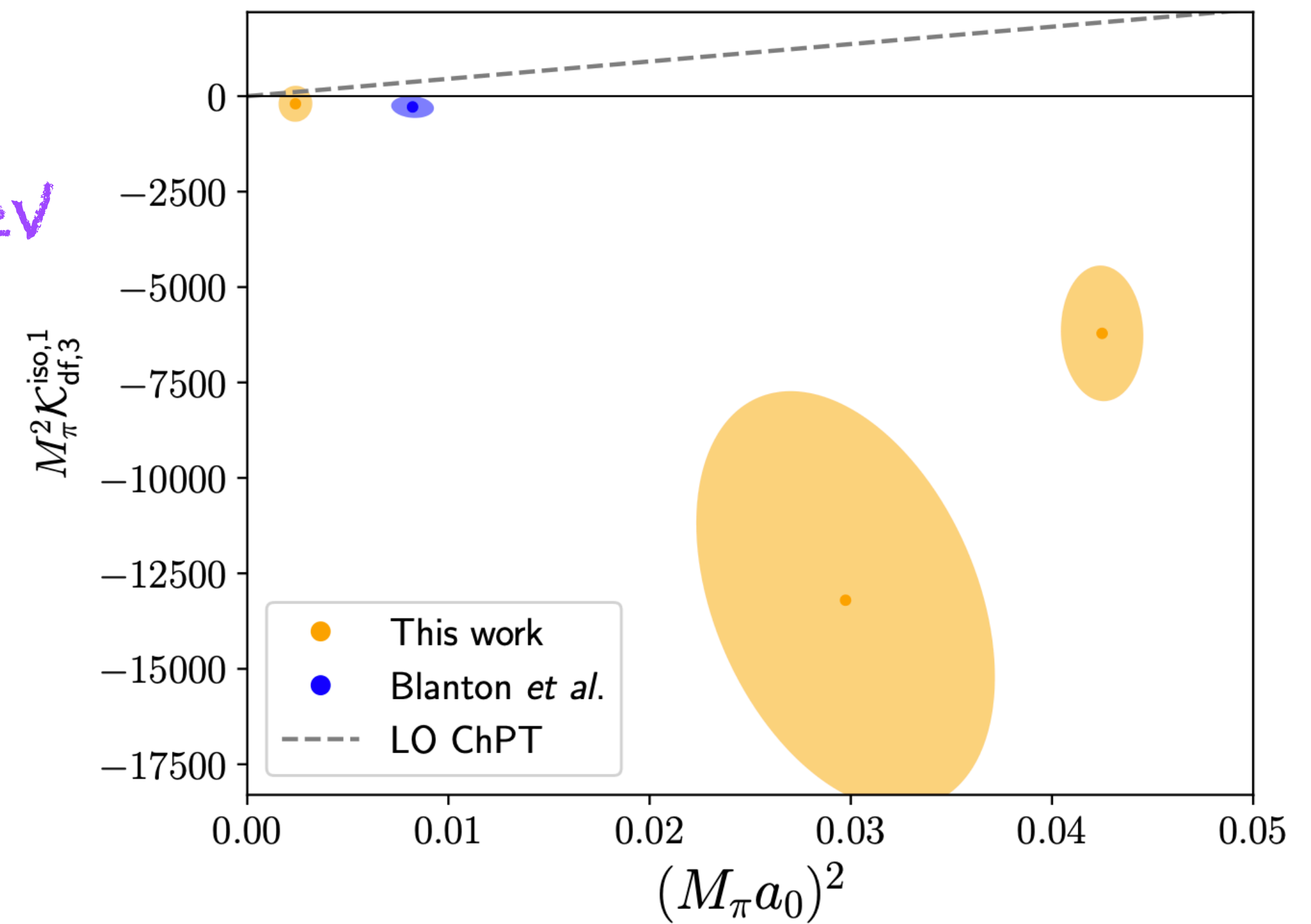
[Fischer, Kostrzewa, Liu, [FRL](#), Ueding, Urbach (ETMC)]

$$\mathcal{K}_{df,3} = \mathcal{K}_{df,3}^{iso,0} + \mathcal{K}_{df,3}^{iso,1} \left(\frac{s - 9M^2}{9M^2} \right)$$

Constant term seems well-behaved



(a) $\mathcal{K}_{df,3}^{iso,0}$



(b) $\mathcal{K}_{df,3}^{iso,1}$

see also other studies:
[Mai et al., Culver et al.]

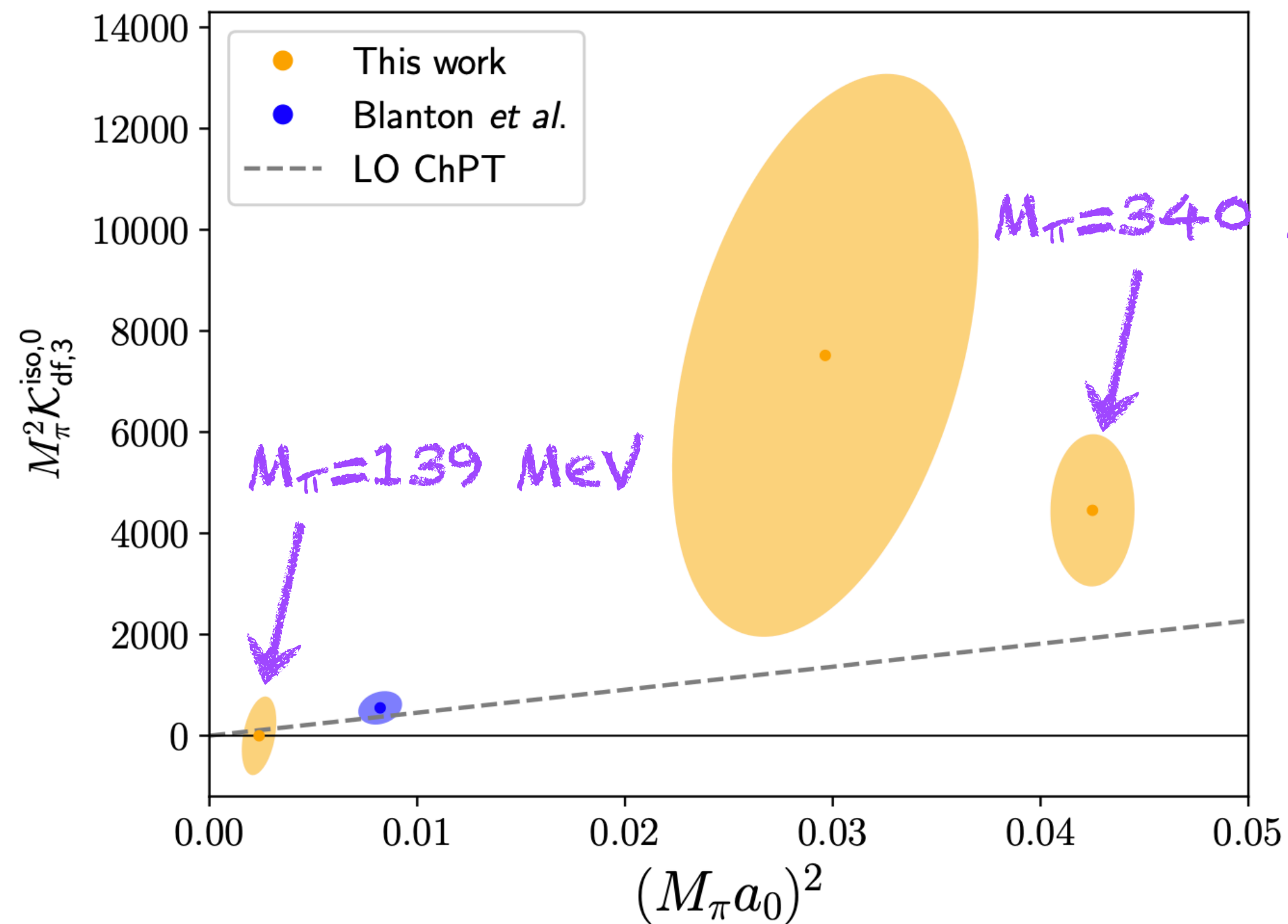
Chiral dependence

On a later article, the chiral dependence of $\mathcal{K}_{df,3}$ has been studied, including physical pions.

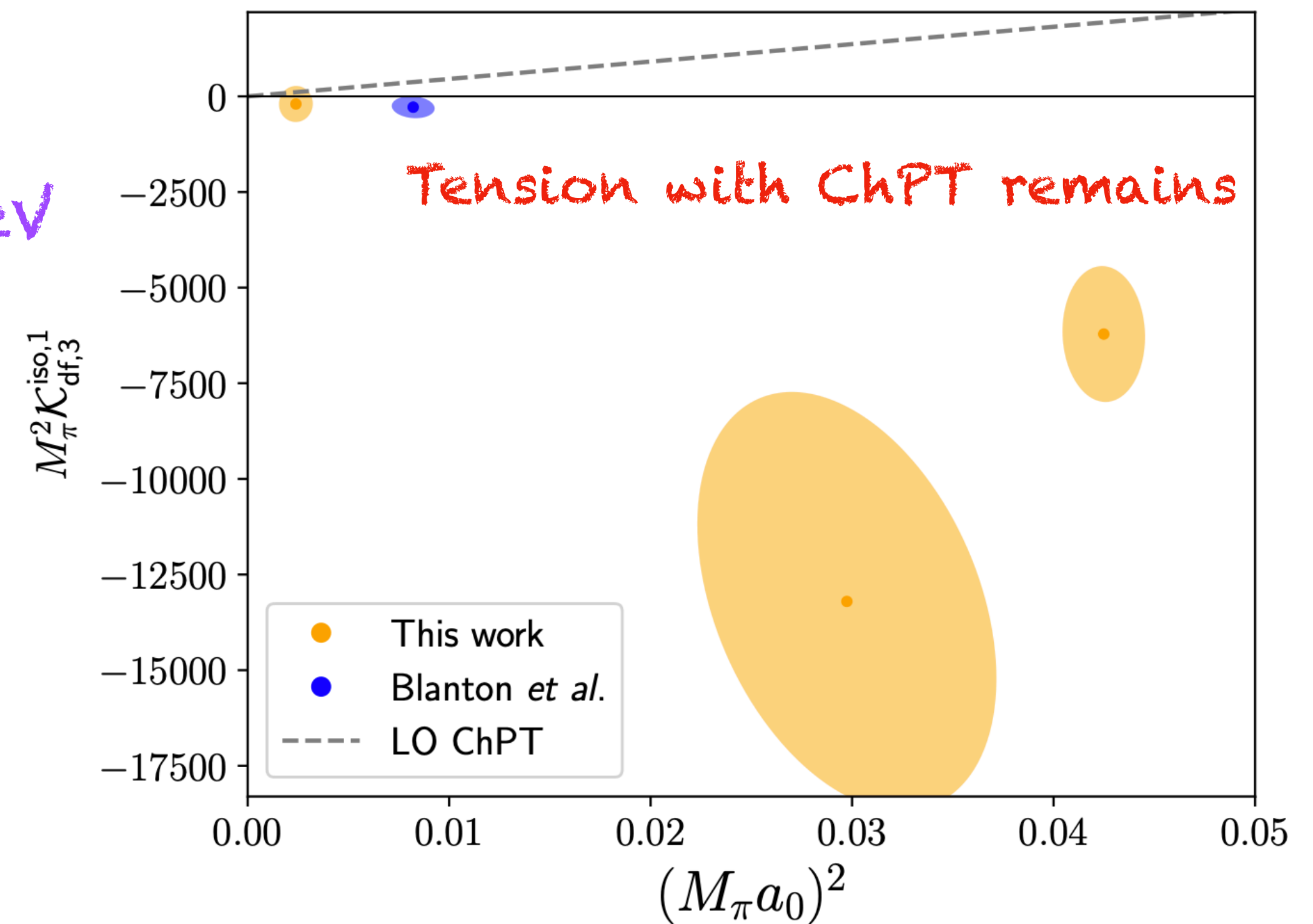
[Fischer, Kostrzewa, Liu, [FRL](#), Ueding, Urbach (ETMC)]

$$\mathcal{K}_{df,3} = \mathcal{K}_{df,3}^{iso,0} + \mathcal{K}_{df,3}^{iso,1} \left(\frac{s - 9M^2}{9M^2} \right)$$

Constant term seems well-behaved



(a) $\mathcal{K}_{df,3}^{iso,0}$



(b) $\mathcal{K}_{df,3}^{iso,1}$

see also other studies:
[Mai et al., Culver et al.]

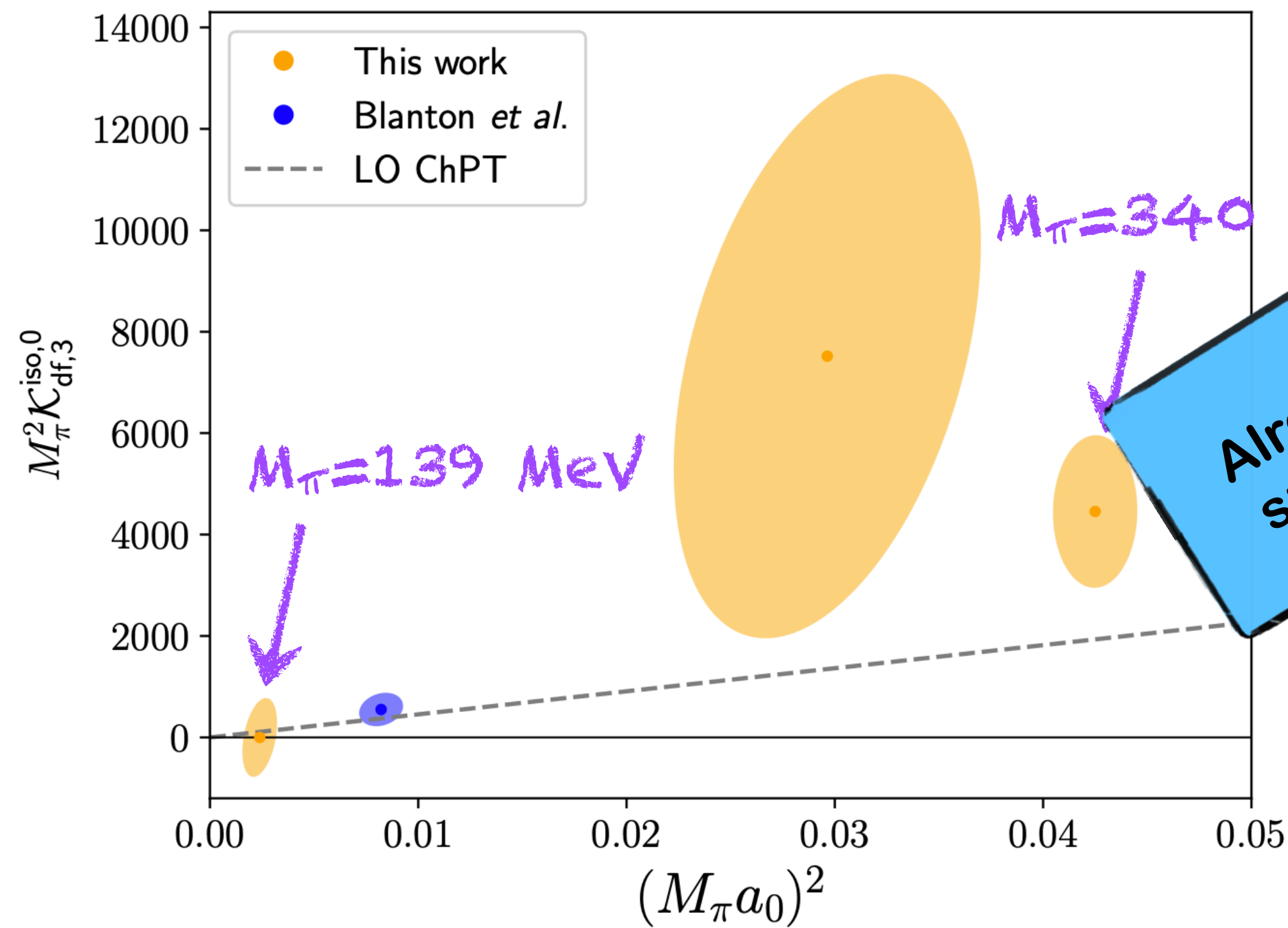
Chiral dependence

On a later article, the chiral dependence of $\mathcal{K}_{df,3}$ has been studied, including physical pions.

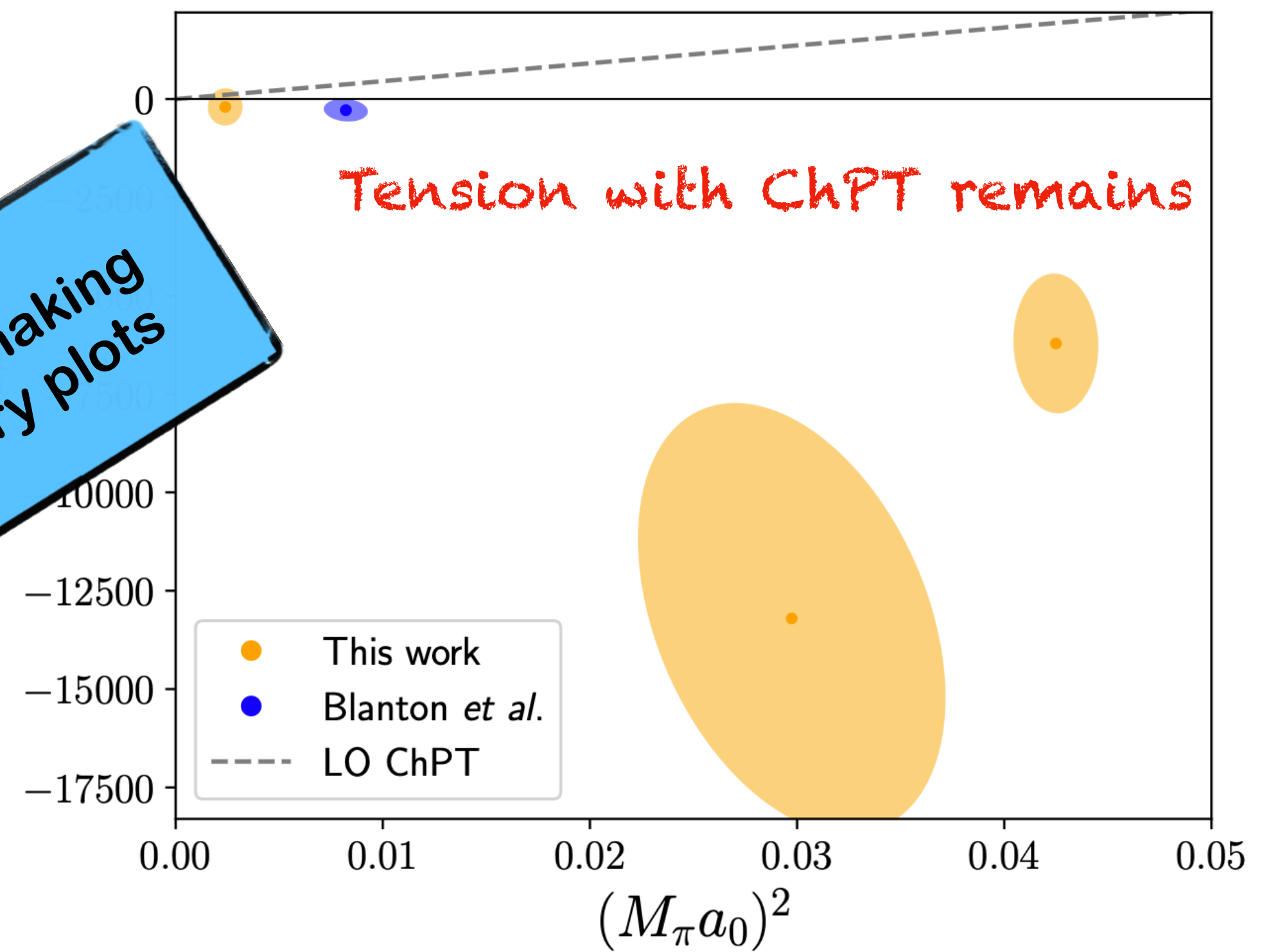
[Fischer, Kostrzewa, Liu, FRL, Ueding, Urbach (ETMC)]

$$\mathcal{K}_{df,3} = \mathcal{K}_{df,3}^{iso,0} + \mathcal{K}_{df,3}^{iso,1} \left(\frac{s - 9M^2}{9M^2} \right)$$

Constant term seems well-behaved



(a) $\mathcal{K}_{df,3}^{iso,0}$



(b) $\mathcal{K}_{df,3}^{iso,1}$

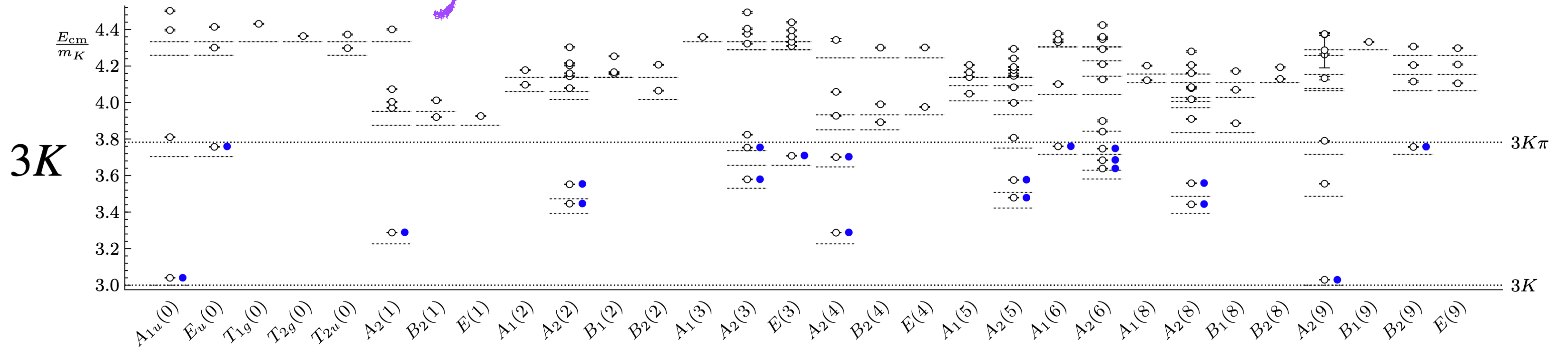
Already making summary plots

see also other studies:
[Mai et al., Culver et al.]

Two- and three-kaons

- Other simple systems can also be studied: $2K^+$ & $3K^+$
- Many energy levels that allow for **s- and d-wave** interactions to be extracted!

Preliminary!

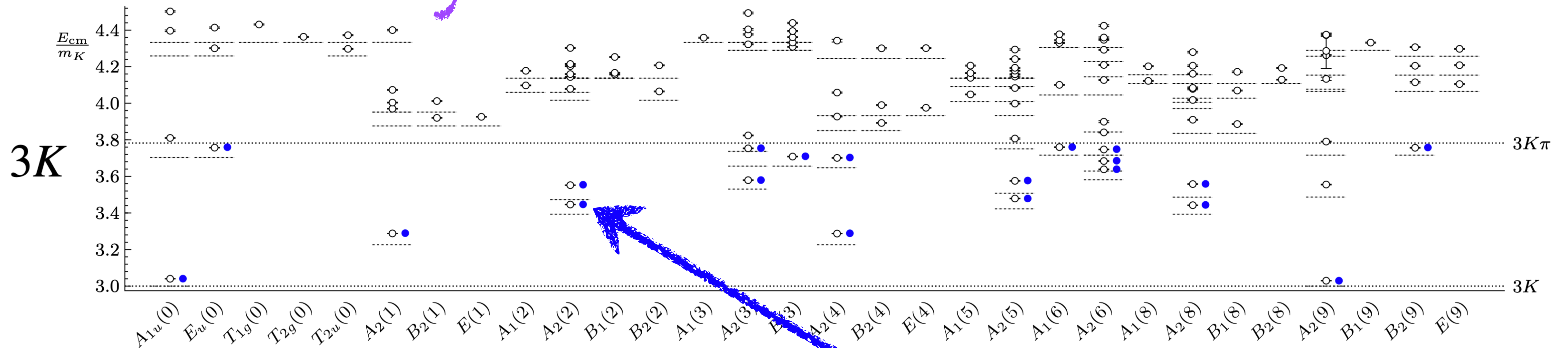


[Blanton, Hanlon, Hörz, Morningstar, FRL, Sharpe (in preparation)]

Two- and three-kaons

- Other simple systems can also be studied: $2K^+$ & $3K^+$
- Many energy levels that allow for **s- and d-wave** interactions to be extracted!

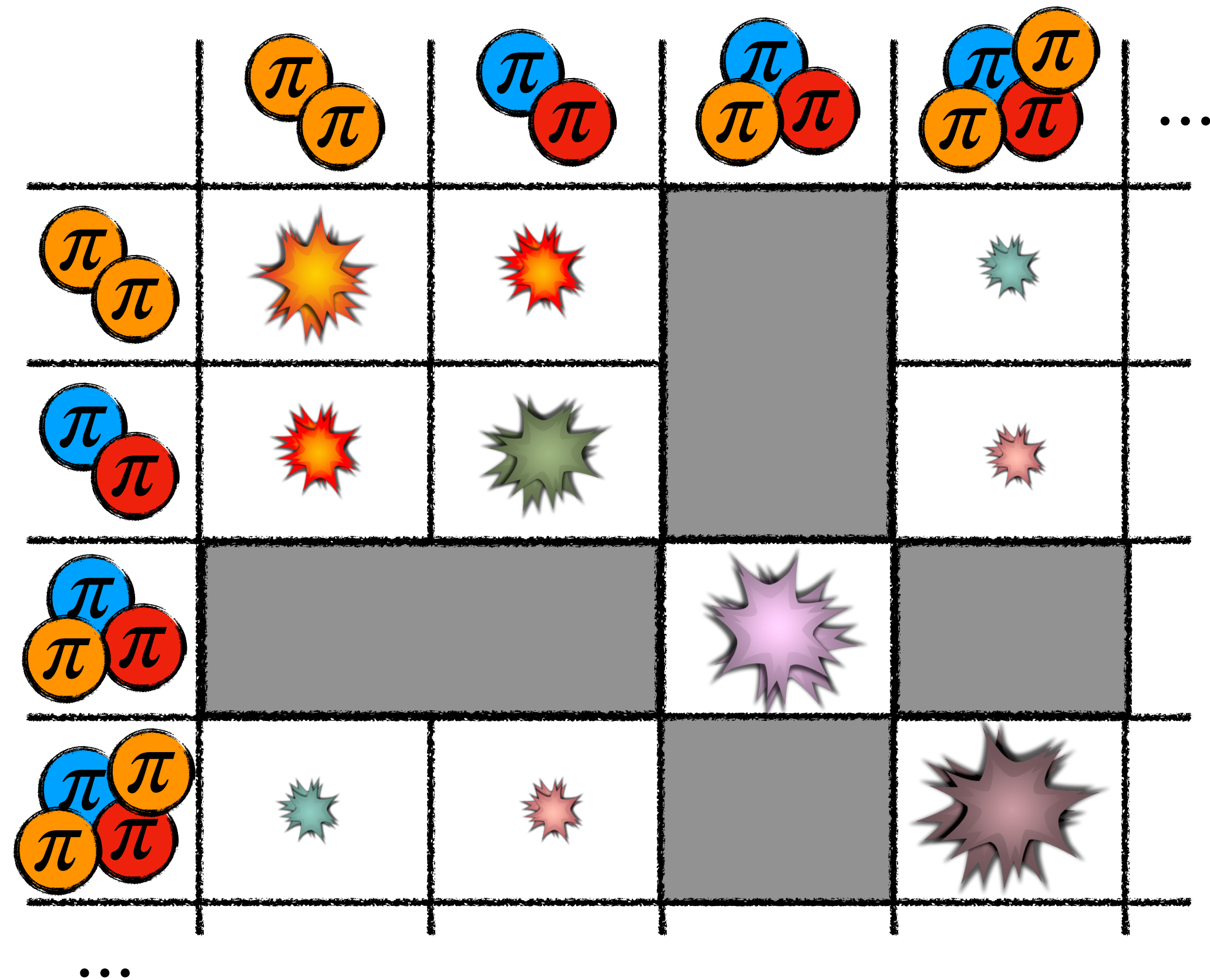
Preliminary!



[Blanton, Hanlon, Hörz, Morningstar, FRL, Sharpe (in preparation)]

prediction from QC3

Outlook



Next steps

Next steps

1. **Lattice studies of 3π resonances (ω , h_1)**
[Hansen, FRL, Sharpe, arXiv:2003.10974]

Next steps

1. **Lattice studies of 3π resonances (ω , h_1)**
[Hansen, FRL, Sharpe, arXiv:2003.10974]
2. **Generalizing the formalism for generic two- and three- particle systems, (e.g. nucleons, Roper resonance)**

Next steps

1. **Lattice studies of 3π resonances (ω , h_1)**
[Hansen, FRL, Sharpe, arXiv:2003.10974]
2. **Generalizing the formalism for generic two- and three- particle systems, (e.g. nucleons, Roper resonance)**
3. **Formalism for three-particle decays, such as $K \rightarrow 3\pi$, $\gamma \rightarrow 3\pi$**

Next steps

1. **Lattice studies of 3π resonances (ω , h_1)**
[Hansen, FRL, Sharpe, arXiv:2003.10974]
2. **Generalizing the formalism for generic two- and three- particle systems, (e.g. nucleons, Roper resonance)**
3. **Formalism for three-particle decays, such as $K \rightarrow 3\pi$, $\gamma \rightarrow 3\pi$**

Next steps

1. **Lattice studies of 3π resonances (ω , h_1)**
[Hansen, FRL, Sharpe, arXiv:2003.10974]
2. **Generalizing the formalism for generic two- and three- particle systems, (e.g. nucleons, Roper resonance)**
3. **Formalism for three-particle decays, such as $K \rightarrow 3\pi$, $\gamma \rightarrow 3\pi$**
4. **Beyond three particles!**

Next steps

1. Lattice studies of 3π resonances (ω , h_1)
[Hansen, FRL, Sharpe, arXiv:2003.10974]
2. Generalizing the formalism for generic two- and three- particle systems,
(e.g. nucleons, Roper resonance)
3. Formalism for three-particle decays, such as $K \rightarrow 3\pi$, $\gamma \rightarrow 3\pi$
4. Beyond three particles!

Thanks!