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**Dispersive study of  $\pi K$  and  $\pi\pi \rightarrow KK$  scattering:  
threshold parameters  
and  
 $\kappa/K_0^*(700)$  resonance determination**

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arXiv:2001.08153. Phys.Rev.Lett. 124 (2020) 17, 172001  
arXiv:2010.1122. To appear in Physics Reports

Theoretical Aspects of Hadron Spectroscopy and Phenomenology  
Valencia Spain, 15-17/12/2020.

Supported by:



## Motivation

- $\pi, K$  appear as final products of almost all hadronic strange processes:  
B, D, decays, CP violation studies...
- $\pi, K$  are Goldstone Bosons of QCD:  
Threshold parameters test Chiral Symmetry Breaking
- Main or relevant source for PDG parameters of:  
 $\kappa/K_0^*(700), K_0^*(1430), K_1^*(892), K_1^*(1410), K_2^*(1410), K_3^*(1780)$

## Problems

- Data: extracted from  $KN \rightarrow \pi KN$ , assuming one pion exchange.  
Large systematic uncertainties and inconsistencies.
- Large model-dependences:  
naïve models often used for parameterizations and resonance poles

## Dispersion Relations (This talk)

**Model independent** constraints,  
precise threshold parameters and pole determinations.  
Enhanced precision

# Data on $\pi K$ scattering: S-channel

Most reliable sets:

Estabrooks et al. 78 (SLAC)

Aston et al. 88 (SLAC-LASS)

$l=1/2$  and  $3/2$  combination

MANY DATA IN CONFLICT

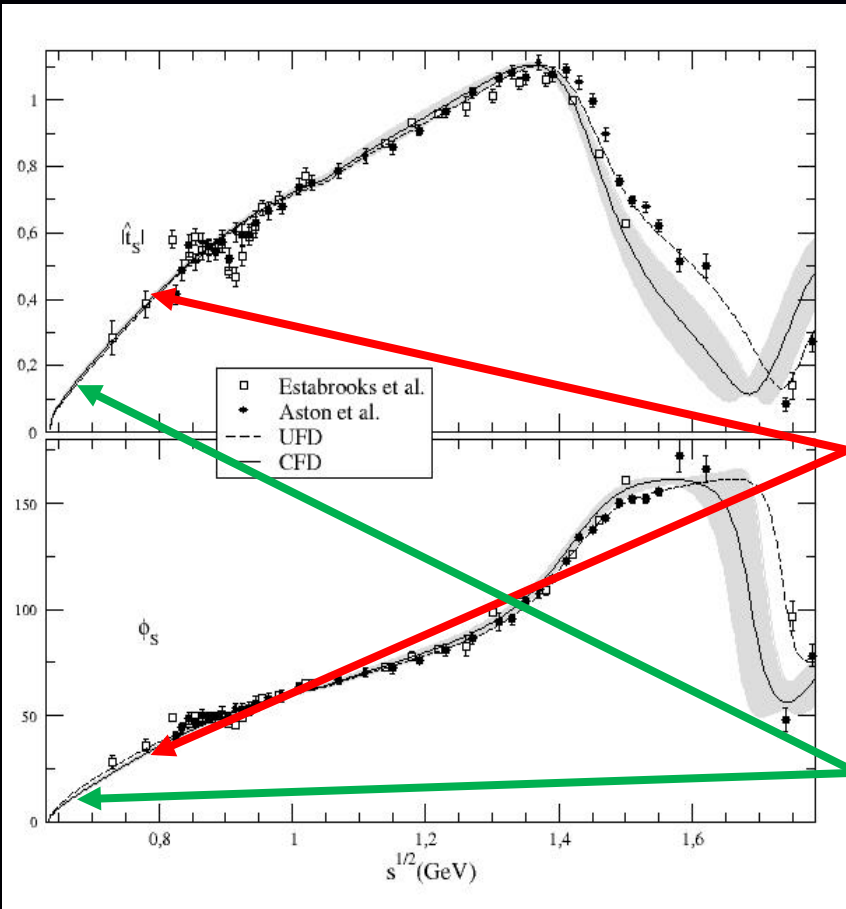
No clear “peak” or phase movement of  $\kappa/K_0^*(800)$  resonance

Definitely NO BREIT-WIGNER shape

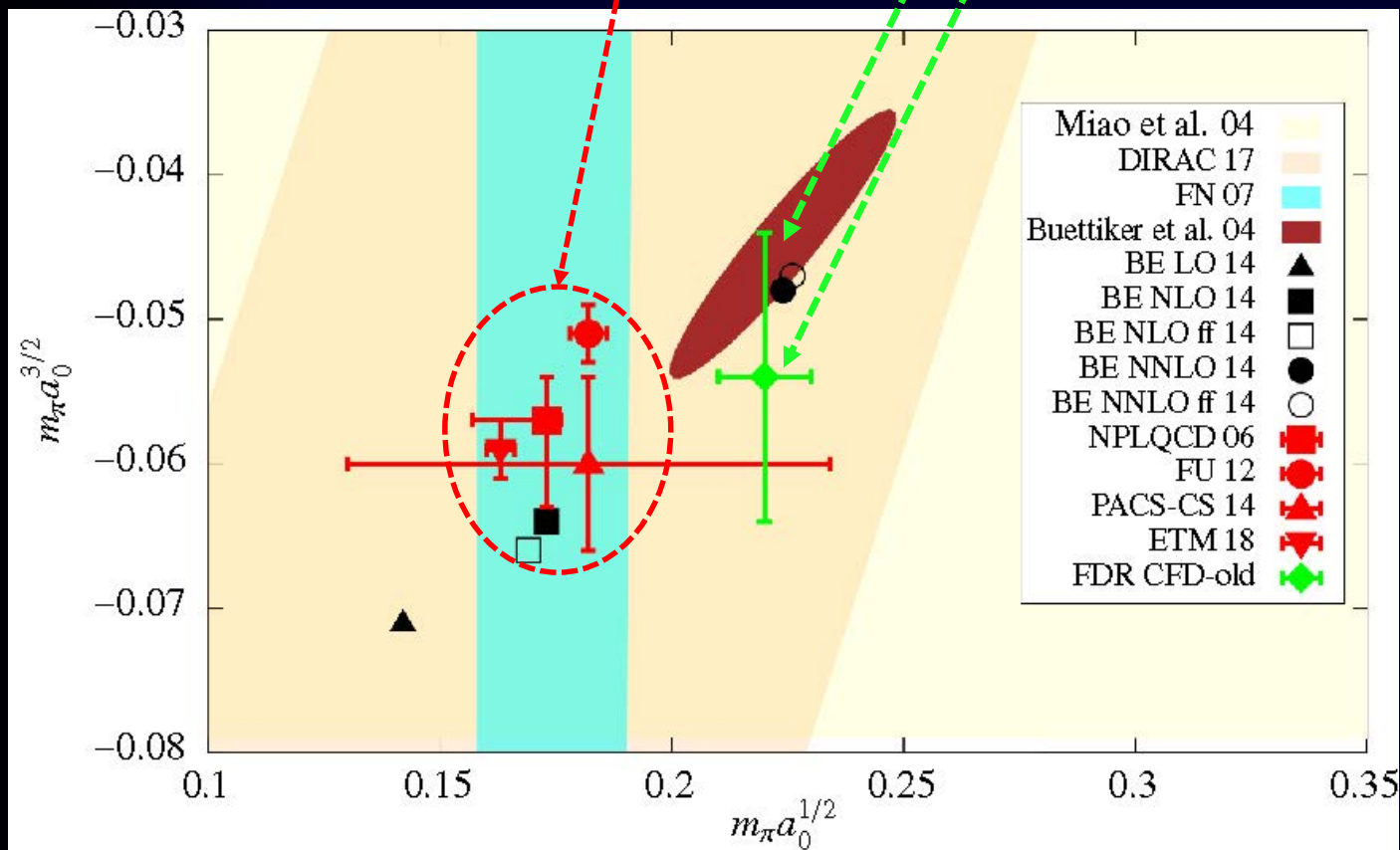
No data near threshold.

Models need dangerous extrapolations.

Dispersion relations  $\rightarrow$  **sum-rules**



- Threshold parameters relevant to test ChPT (NNLO at present).
- Present tension between **lattice** and dispersive results



- Dalitz 1965: “Quite apart from the model discussed here,...such  $K^*$  states are expected to exist simply on the basis of  $SU(3)$ ” Procs. Oxford Int. Conf. on Elementary Particles 1965
- Many claims at different masses, narrow, wide... claims of absence. Confusion

1967  
attitude

## REVIEWS OF

# MODERN PHYSICS

VOLUME 39, NUMBER 1

### Data on Particles and Resonant States\*

ARTHUR H. ROSENFELD, ANGELA BARBARO-GALTIERI, WILLIAM J. PODOLSKY, I.  
PAUL SODING, CHARLES G. WOHL  
*Lawrence Radiation Laboratory, University of California, Berkeley, California*  
MATTS ROOS  
*CERN, Geneva, Switzerland*  
WILLIAM J. WILLIS  
*Dept. of Physics, Yale University, New Haven, Connecticut*

Data on the properties of leptons, mesons, and baryons are listed, referenced, averaged, and summarized in tables and wallet cards. This is an updating of the *Reviews of Modern Physics* article of October 1965.

1. The  $\kappa(725)$  (Lynch, Rittenberg, Rosenfeld, Söding, Dec. 1966)

We are beginning to think that  $\kappa$  should be classified along with flying saucers, the Loch Ness Monster, and the Abominable Snowman. We have heard of several experiments which were supposed to confirm it, and each one has either failed completely or failed to find it in the sought-for channel, but found instead a small  $K\pi$  peak near 725 MeV in some other channel.

- Removed from Review of Particle Physics in 1976 (with the  $\sigma$ )
- Back to RPP in 2004 as “controversial”  $K_0^*(800)$ . Omitted from summary tables

Strong support for  $\kappa/K_0^*(800)$  from chiral theories and experimental decays of heavier mesons, but rigorous model-independent extractions absent. Often inadequate Breit-Wigner formalisms

- Omitted from the 2018PDG summary table since, “needs confirmation”

Since the 70's 90's, all descriptions of data respecting unitarity and chiral symmetry find a pole at  $M=650-770$  MeV and  $\Gamma\sim 550$  MeV or larger.

Best determination came from a SOLUTION of a Roy-Steiner dispersive formalism, consistent with UChPT

Decotes Genon et al 2006

2017PDG  $K_0^*(800)$  dominated by such a SOLUTION

$$M-i\Gamma/2=(682\pm 29)-i(273\pm 12) \text{ MeV}$$

PDG2018:

**(630-730)-i(260-340) MeV**  
name changed to  $K_0^*(700)$

PDG2020:

$K_0^*(700)$  Makes it to the summary tables.  
Still “Needs Confirmation”

We were encouraged by PDG to confirm it with a dispersive DATA analysis (this talk)

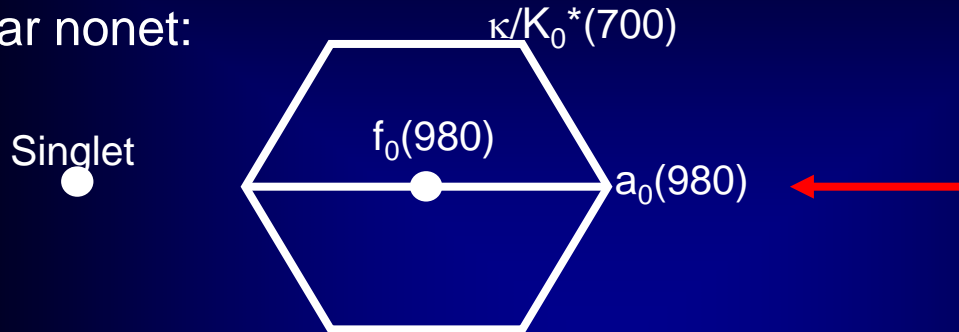
# MOTIVATION: The light scalar controversy.

## ● Scalar SU(3) multiplets identification controversial

- Too many or too few resonances for decades  
But there is an emerging picture



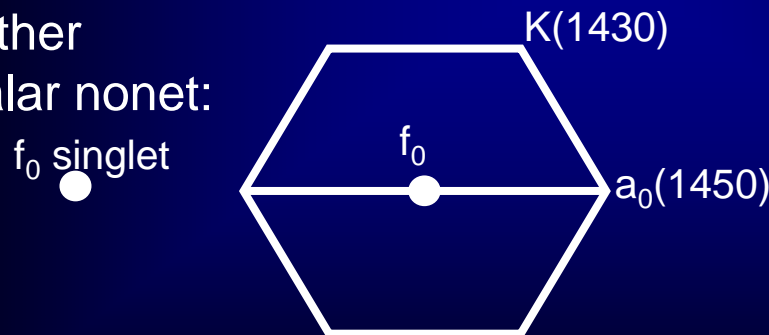
A Light scalar nonet:



Non-strange heavier!!  
Inverted hierarchy problem  
For quark-antiquark

$f_0(500)$  and  $f_0(980)$  are really octet/singlet mixtures

+ Another heavier scalar nonet:



+ glueball ?



Enough  $f_0$  states have been observed:  $f_0(1370)$ ,  $f_0(1500)$ ,  $f_0(1700)$ .

The whole picture is complicated by mixture between them (lots of works here)

**Only the light  $\kappa(700)$  or  $K_0^*(700)$  "Needs Confirmation" @ PDG2020**



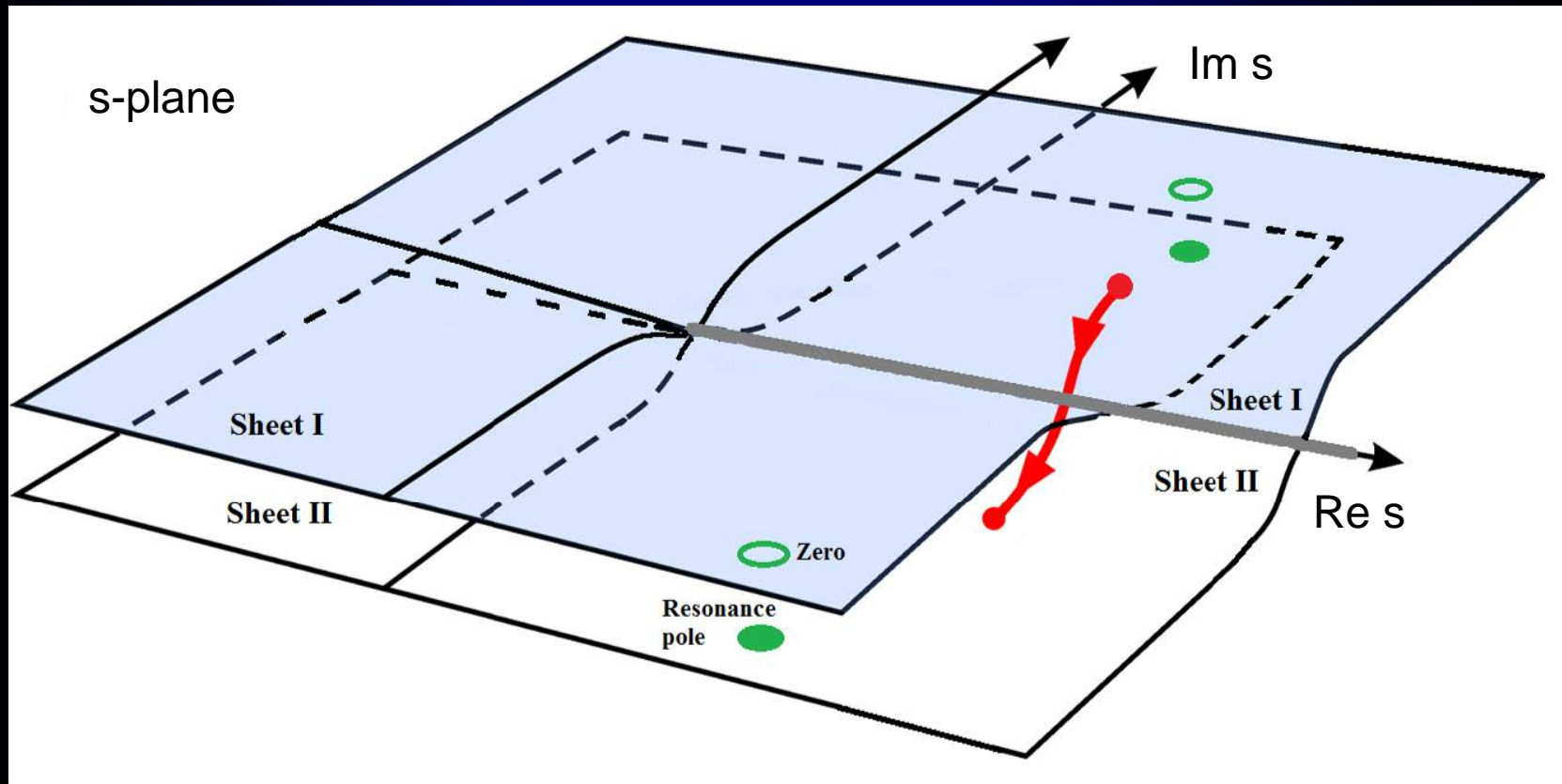
## Resonances as poles

The Breit-Wigner shape is just an approximation for narrow and isolated resonances

The universal features of resonances are their pole positions and residues \*

$$\sqrt{s_{pole}} \approx M - i \Gamma/2$$

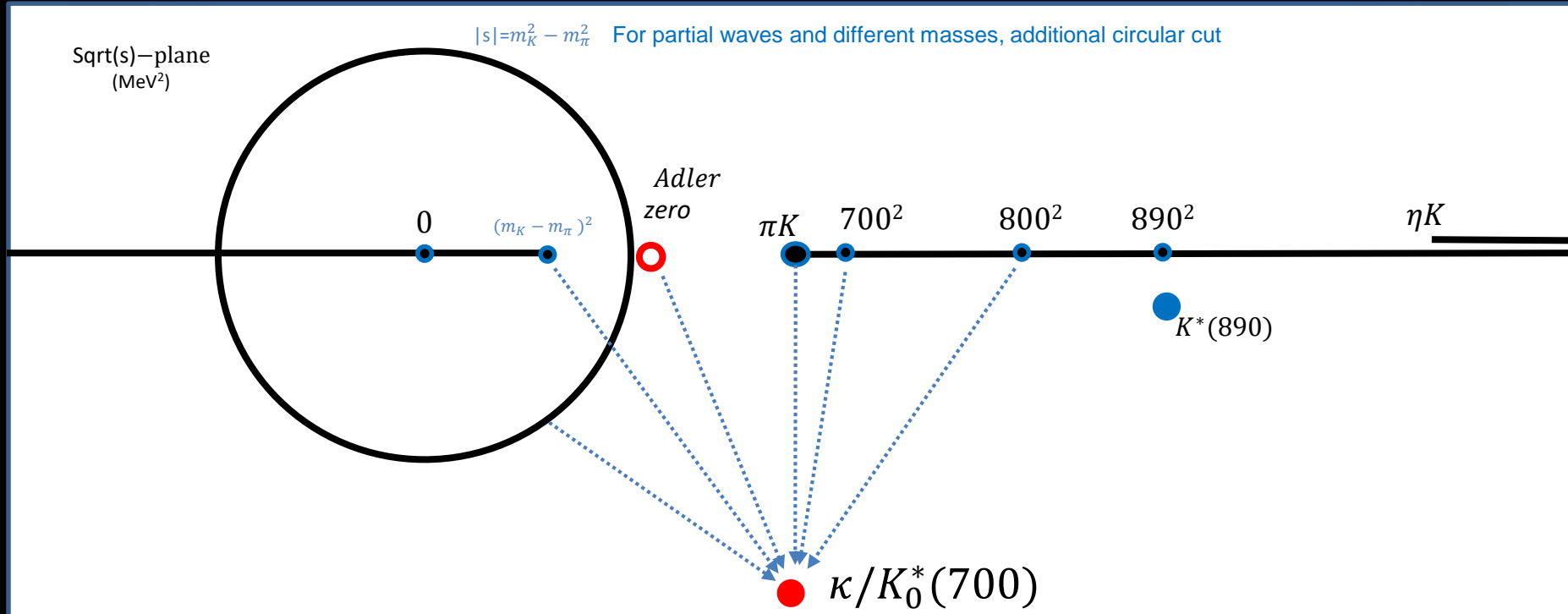
\*in the Riemann sheet obtained from an analytic continuation through the physical cut





# Why so much worries about low energy and CORRECT ANALYTIC STRUCTURE?

Analyticity is expressed in the  $s$ -variable, not in  $\text{Sqrt}(s)$



Important for  
the  $\kappa/K_0^*(700)$   
and threshold  
parameters

- Threshold behavior (chiral symmetry)
- Subthreshold behavior (chiral symmetry  $\rightarrow$  Adler zeros)
- Other cuts (Left & circular)
- Avoid spurious singularities

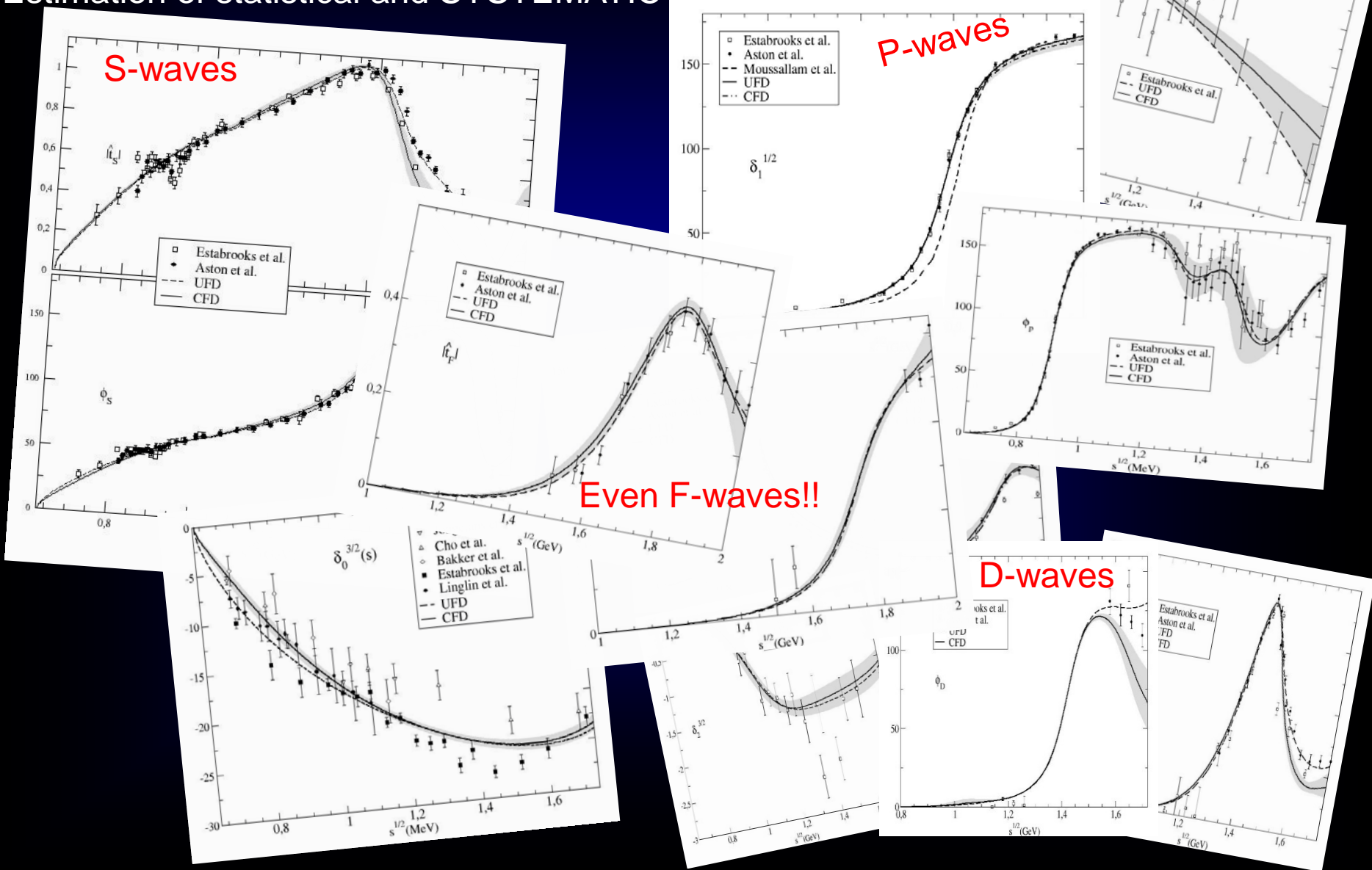
Less important for other resonances...

# Our Dispersive/Analytic Approach for $\pi K$ and strange resonances

## FIRST STEP:

**Simple Unconstrained Fits (UFD)** to  $\pi K$  and  $\pi\pi \rightarrow KK$  partial-

Estimation of statistical and SYSTEMATIC errors



**Simple Unconstrained Fits** to  $\pi K$  partial-wave Data (UFD).

Estimation of statistical and SYSTEMATIC errors

**Forward Dispersion Relations:**

Left cut easy to rewrite

Relate amplitudes, not partial waves

Not direct access to pole

## Forward dispersion relations for $K \pi$ scattering.

Since interested in the resonance region, we use minimal number of subtractions

Defining the  $s \leftrightarrow u$  symmetric and anti-symmetric amplitudes at  $t=0$

$$T^+(s) = \frac{T^{1/2}(s) + 2T^{3/2}(s)}{3} = \frac{T^{I_t=0}(s)}{\sqrt{6}},$$
$$T^-(s) = \frac{T^{1/2}(s) - T^{3/2}(s)}{3} = \frac{T^{I_t=1}(s)}{2}.$$

We need one subtraction for the symmetric amplitude

$$\text{Re}T^+(s) = T^+(s_{\text{th}}) + \frac{(s - s_{\text{th}})}{\pi} P \int_{s_{\text{th}}}^{\infty} ds' \left[ \frac{\text{Im}T^+(s')}{(s' - s)(s' - s_{\text{th}})} - \frac{\text{Im}T^+(s')}{(s' + s - 2\Sigma_{\pi K})(s' + s_{\text{th}} - 2\Sigma_{\pi K})} \right],$$

And none for the antisymmetric

$$\text{Re}T^-(s) = \frac{(2s - 2\Sigma_{\pi K})}{\pi} P \int_{s_{\text{th}}}^{\infty} ds' \frac{\text{Im}T^-(s')}{(s' - s)(s' + s - 2\Sigma_{\pi K})}.$$

where  $\Sigma_{\pi K} = m_{\pi}^2 + m_K^2$

Simple Unconstrained Fits to  $\pi K$  partial-wave Data (UFD).

Estimation of statistical and SYSTEMATIC errors

## Forward Dispersion Relations:

Left cut easy to rewrite

Relate amplitudes, not partial waves

Not direct access to pole

- As  $\pi K$  checks: Small inconsistencies.

# Forward Dispersion Relation analysis of $\pi K$ scattering DATA up to 1.6 GeV

(not a solution of dispersion relations, but a constrained fit)

A.Rodas & JRP, PRD93,074025 (2016)

First observation:

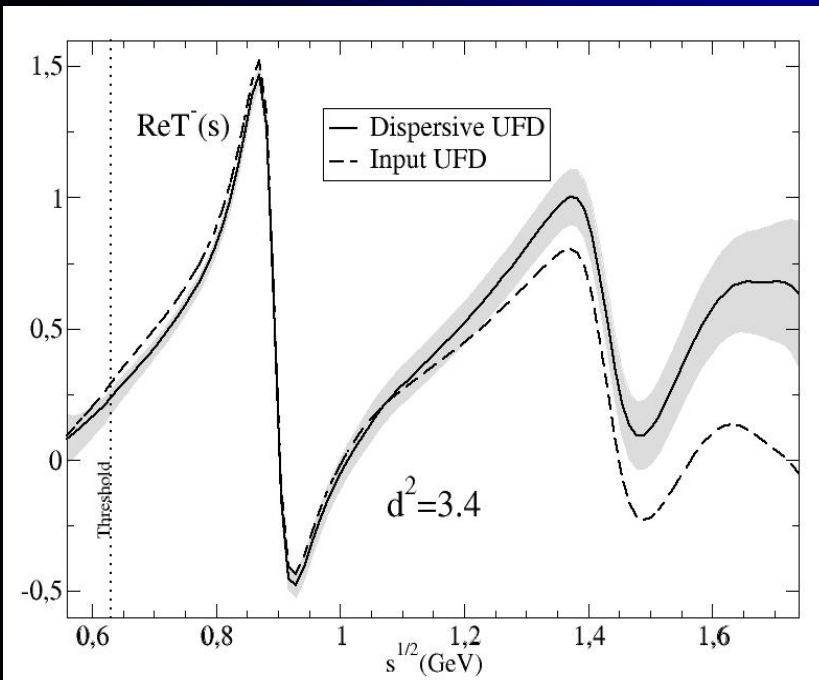
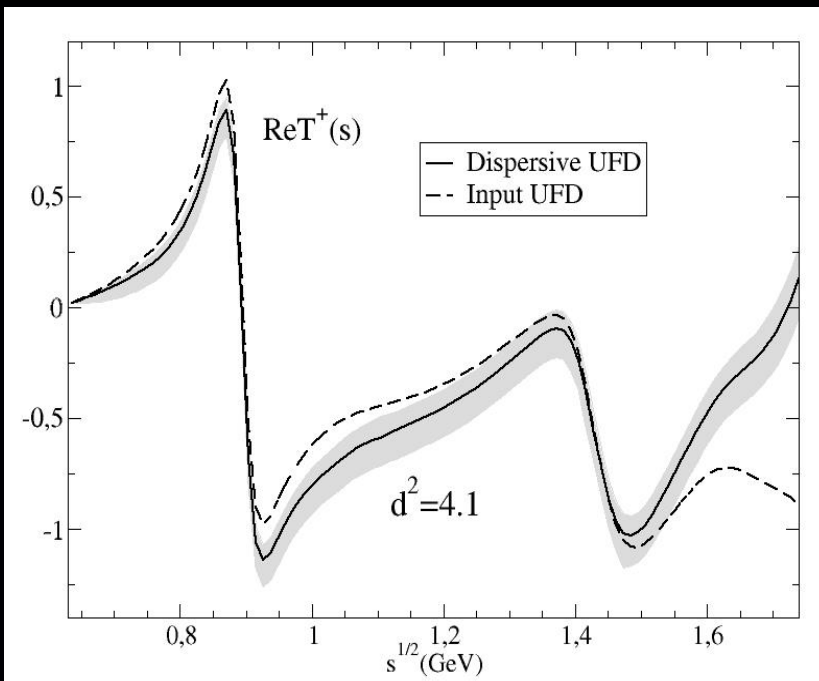
Forward Dispersion relations

Not well satisfied by data

Particularly at high energies

So we use

Forward Dispersion Relations as CONSTRAINTS on fits



Simple Unconstrained Fits to  $\pi K$  partial-wave Data (UFD).

Estimation of statistical and SYSTEMATIC errors

## Forward Dispersion Relations:

Left cut easy to rewrite

Relate amplitudes, not partial waves

Not direct access to pole

• As  $\pi K$  checks: Small inconsistencies.

• As constraints:

**$\pi K$  consistent fits up to 1.6 GeV**

JRP, A.Rodas, Phys.Rev. D93 (2016)



# How well Forward Dispersion Relations are satisfied by unconstrained fits

Every 22 MeV calculate the difference between both sides of the DR /uncertainty

Define an averaged  $\chi^2$  over these points, that we call  $d^2$

$d^2$  close to 1 means that the relation is well satisfied

$d^2 \gg 1$  means the data set is inconsistent with the relation.

This can be used to check DR

To obtain CONSTRAINED FITS TO DATA (CFD) we minimize:

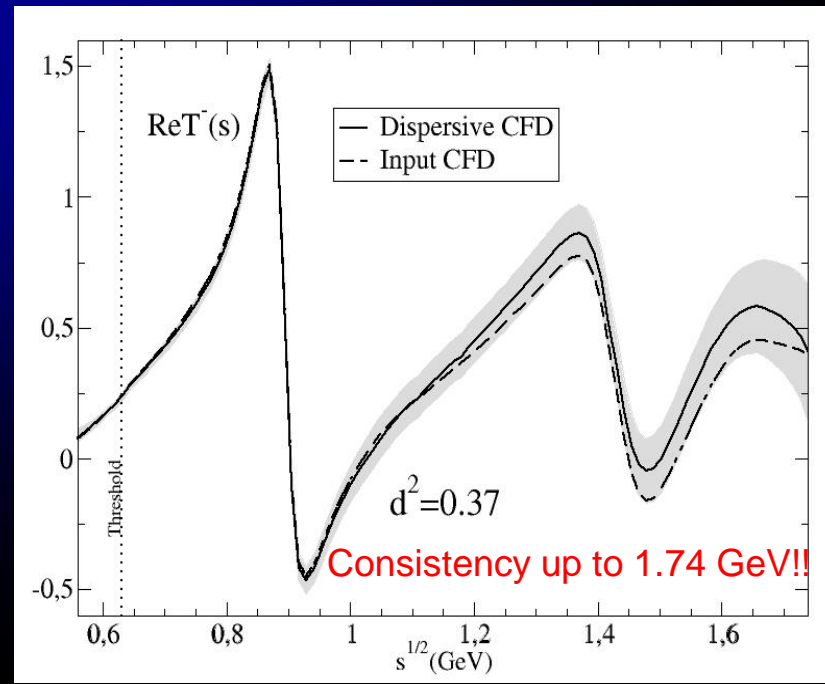
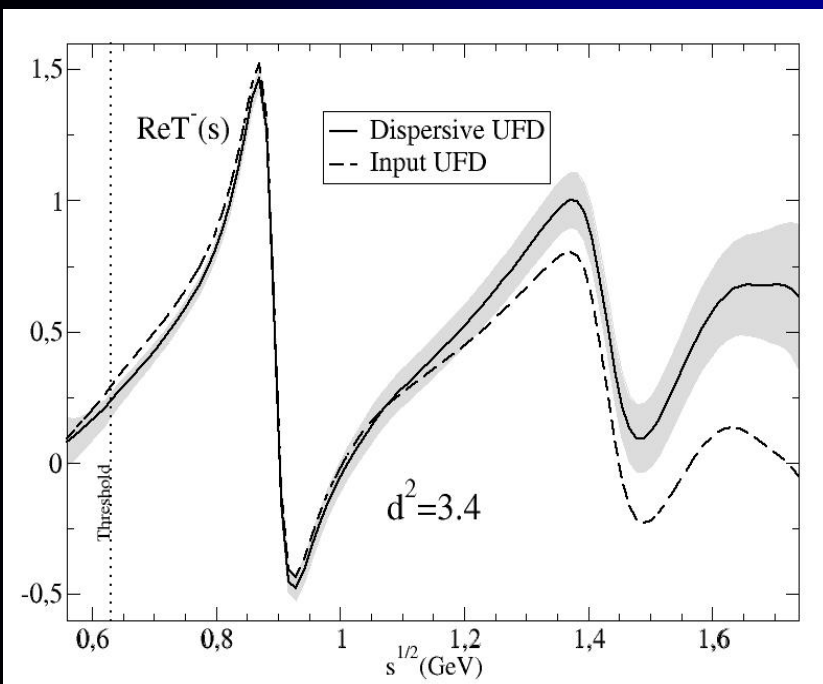
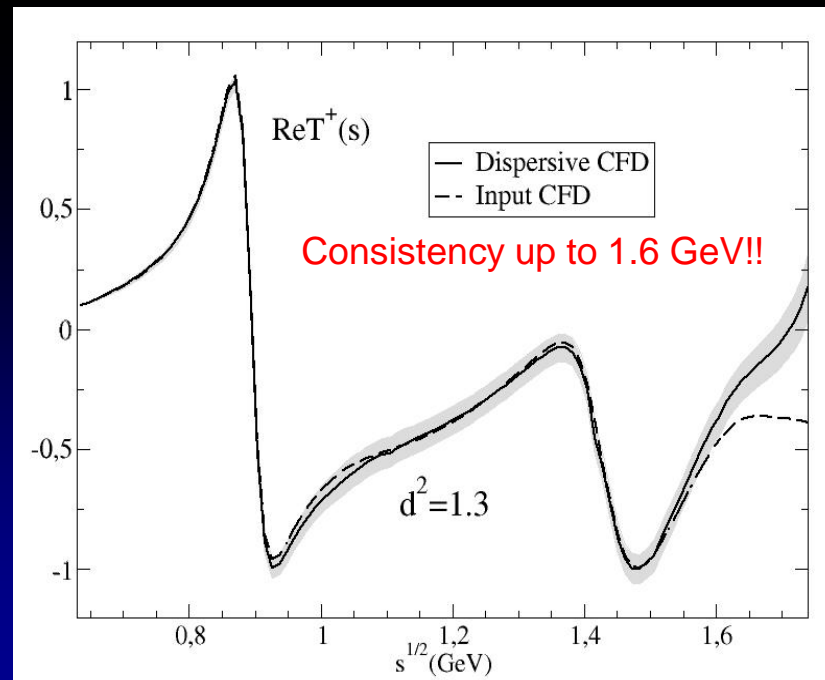
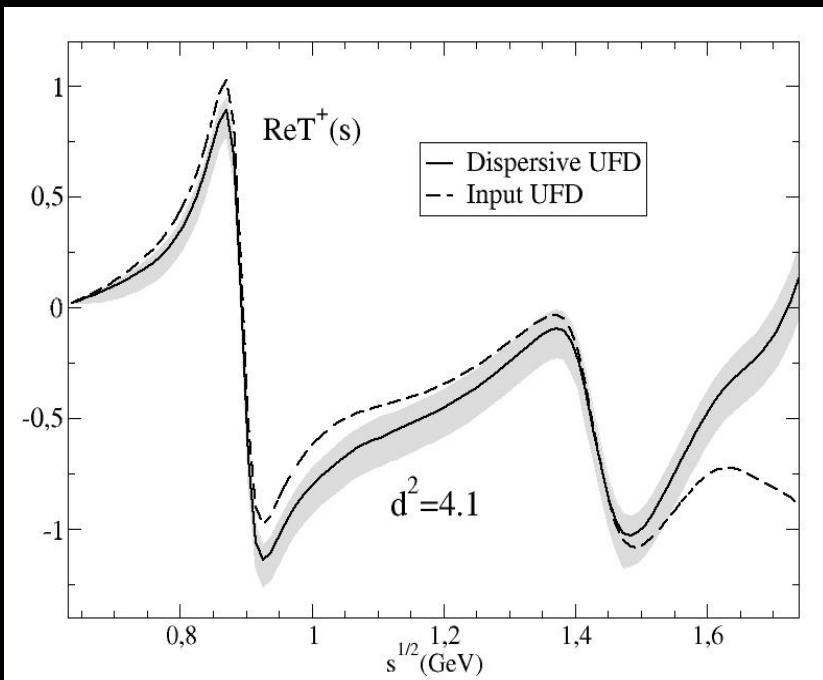
$$W^2(d_{T+}^2 + d_{T-}^2) + \sum_{I=\frac{1}{2}, \frac{3}{2}} \left( \frac{\Delta_I}{\delta\Delta_I} \right)^2 + \sum_k \left( \frac{P_k^{UFD} - P_k}{\delta P_k^{UFD}} \right)^2,$$

2 FDR's

Sum Rules  
threshold

Parameters of the  
unconstrained data fits

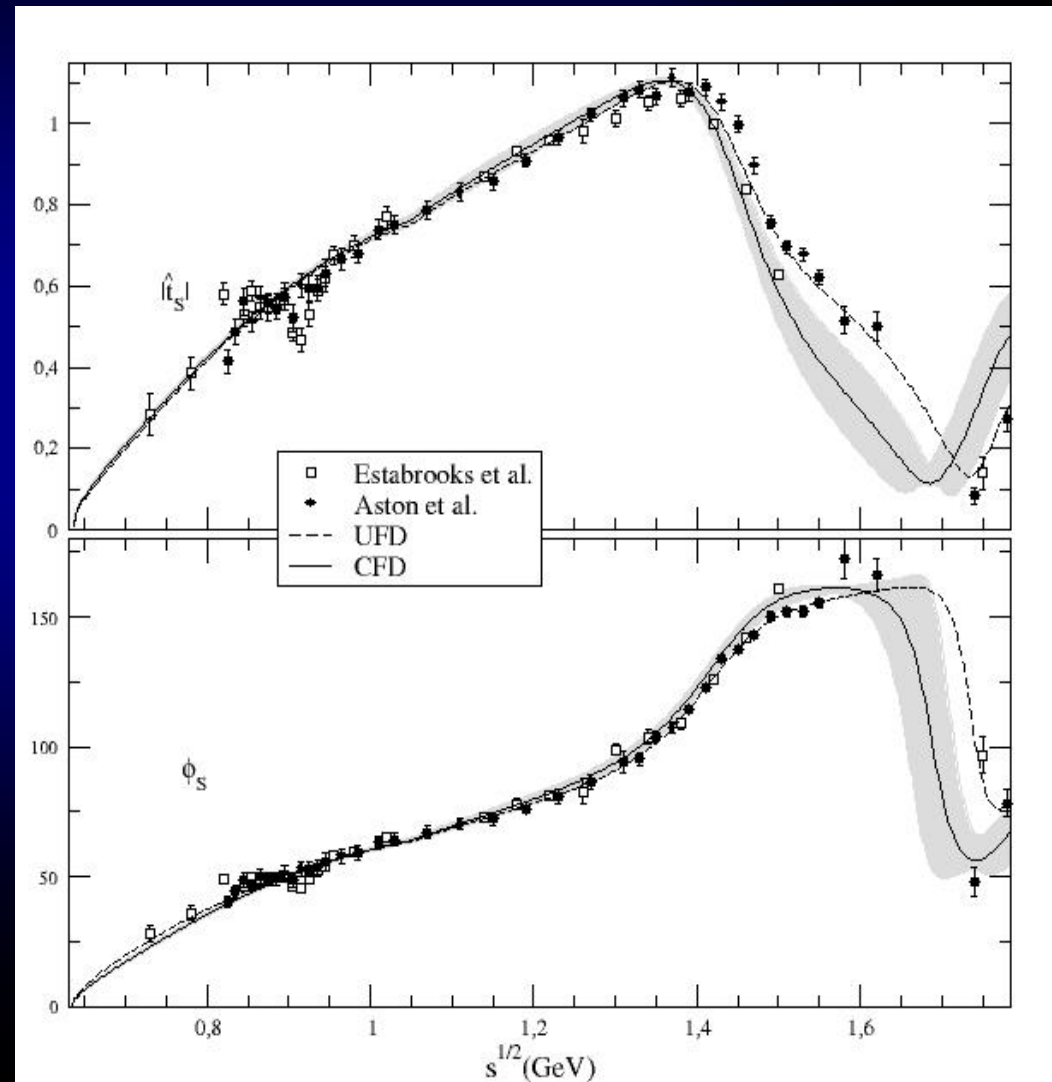
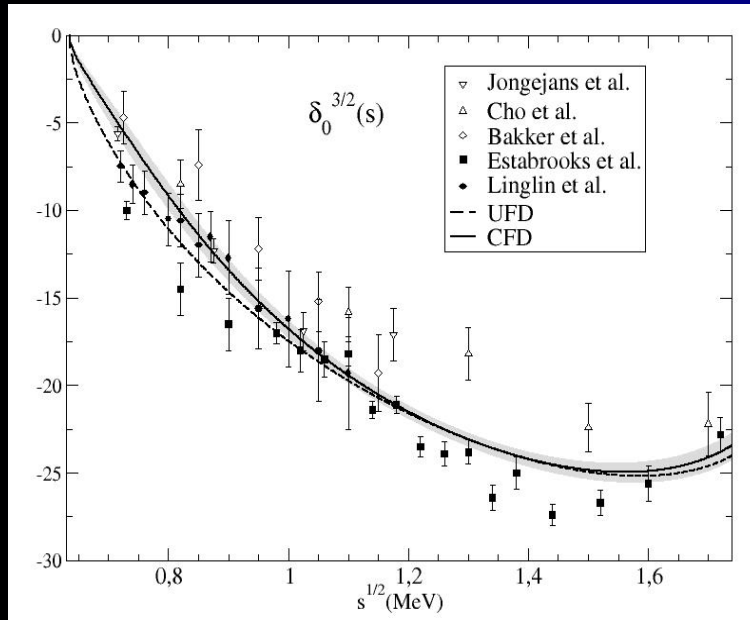
W roughly counts the number  
of effective degrees of freedom  
(sometimes we add weight on certain energy regions)



# From Unconstrained (UFD) to Constrained Fits to data (CFD)

S-waves. The most interesting for the  $K_0^*$  resonances

Largest changes from UFD to CFD  
at higher energies



Simple Unconstrained Fits to  $\pi K$  partial-wave Data (UFD).

Estimation of statistical and SYSTEMATIC errors

## Forward Dispersion Relations:

Left cut easy to rewrite

Relate amplitudes, not partial waves

Not direct access to pole

- As  $\pi K$  checks: Small inconsistencies.
- As constraints:  
 **$\pi K$  consistent fits up to 1.6 GeV** JRP, A.Rodas, Phys.Rev. D93 (2016)
- Padé Sequences to extract poles: reduced model dependence on strange resonances

JRP, A. Rodas. J. Ruiz de Elvira, Eur.Phys.J. C77 (2017)

## Partial-wave $\pi K$ Dispersion Relations

Need  $\pi\pi \rightarrow KK$  to rewrite left cut. Range optimized.

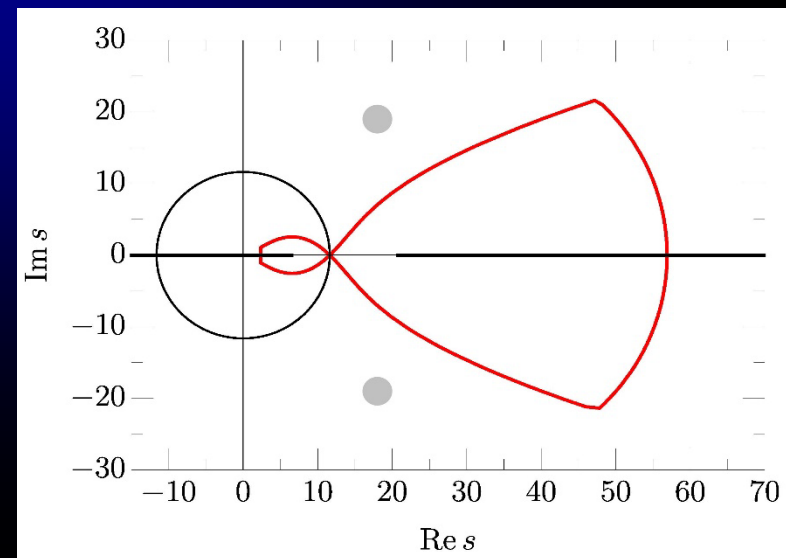
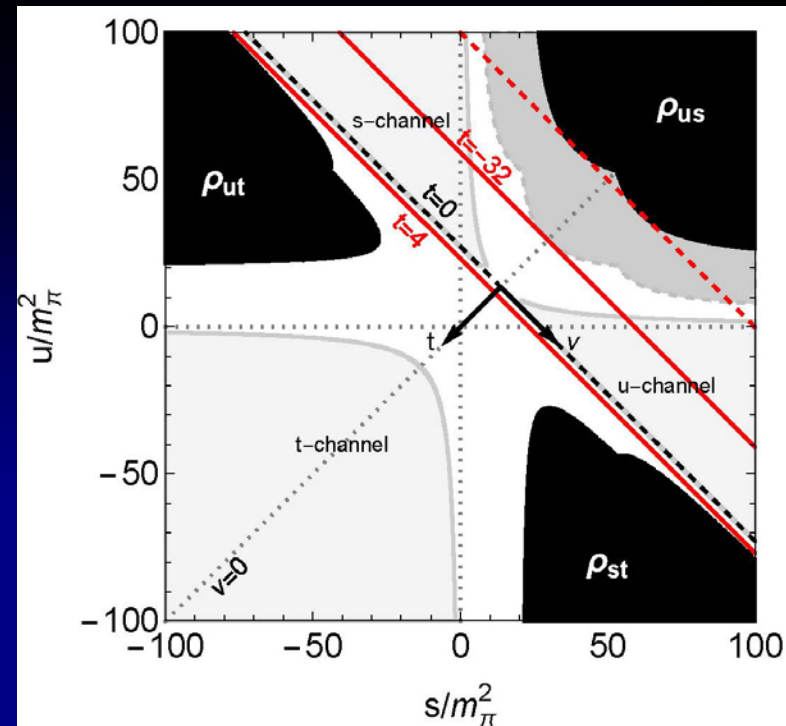
# Partial Wave $\pi K \rightarrow \pi K$ and $\pi\pi \rightarrow KK$ Dispersion Relations (Roy-Steiner eqs.)

To get a resonance pole we need  
PARTIAL-WAVE dispersion relations.

Their applicability is limited  
-by the double spectral regions  
-by the Lehmann ellipses  
(way too technical. See our appendices)

Two possibilities in the literature:

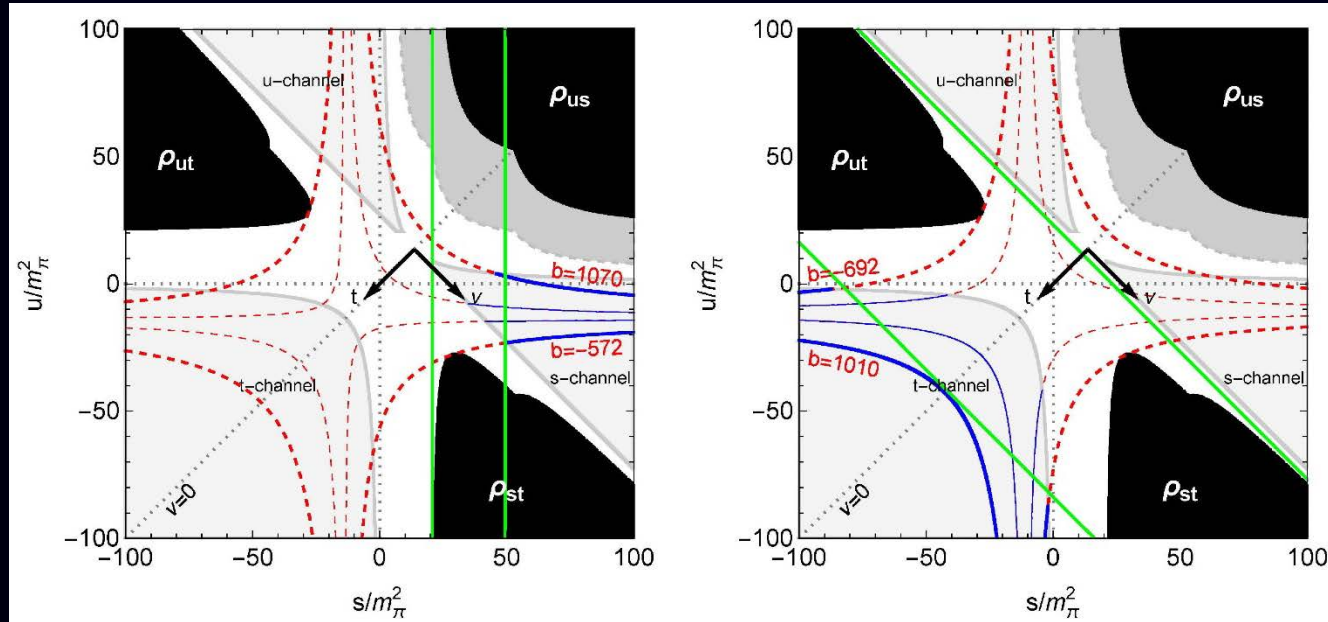
- 1) Integrate "t" for fixed-t dispersion relations.  
Fine for the real axis (1.1 GeV)  
Very mild dependence on  $\pi\pi \rightarrow KK$   
but bad to reach the pole.  
Were used to obtain solutions by the Paris Group  
We will only used them as constraints on data



# $\pi K \rightarrow \pi K$ and $\pi\pi \rightarrow KK$ Hyperbolic Dispersion Relations (HDR)

2) Integrate along  $(s-a)(u-a)=b$  hyperbolae in the Mandelstam plane

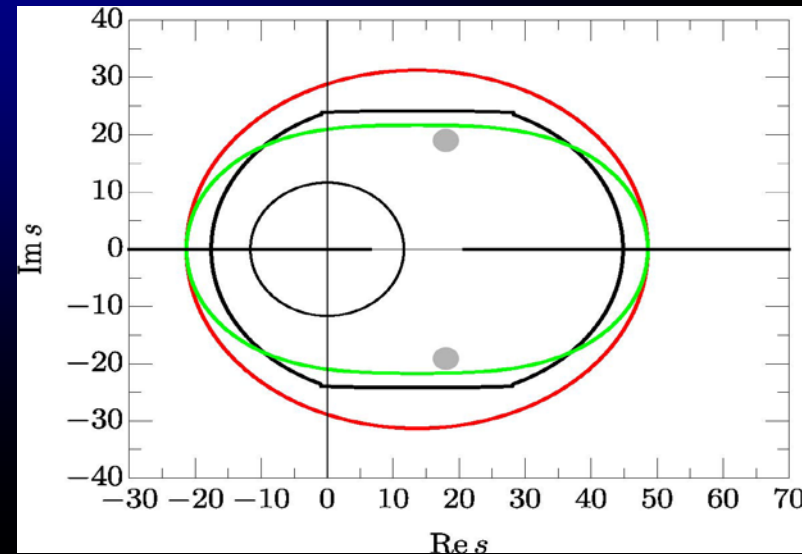
We tuned  $a=-13m_\pi^2$  to maximize applicability for  $\pi\pi \rightarrow KK$  up to 1.47 GeV.



Applicability range slightly smaller in real axis for  $\pi K$ , but covers the kappa pole if a chosen appropriately

We will use them as constraints and to get the pole.

$a=-10m_\pi^2$  chosen to include also error bars inside applicability region





# $\pi K \rightarrow \pi K$ and $\pi\pi \rightarrow KK$ Hyperbolic Dispersion Relations (HDR)

$g_J^I = \pi\pi \rightarrow KK$  partial waves. We study  $(I,J)=(0,0),(1,1),(0,2)$

$f_J^I = K\pi \rightarrow K\pi$  partial waves. Taken from previous dispersive study

JRP, A. Rodas PRD 2018

$$\begin{aligned}
 g_0^0(t) &= \frac{\sqrt{3}}{2} m_+ a_0^+ + \frac{t}{\pi} \int_{4m_\pi^2}^{\infty} \frac{\text{Im } g_0^0(t')}{t'(t'-t)} dt' - \frac{t}{\pi} \sum_{\ell \geq 2} \int_{4m_\pi^2}^{\infty} \frac{dt'}{t'} G_{0,2\ell-2}^0(t,t') \text{Im } g_{2\ell-2}^0(t') + \sum_{\ell} \int_{m_+^2}^{\infty} ds' G_{0,\ell}^+(t,s') \text{Im } f_{\ell}^+(s'), \\
 g_1^1(t) &= \frac{1}{\pi} \int_{4m_\pi^2}^{\infty} \frac{\text{Im } g_1^1(t')}{t'-t} dt' - \sum_{\ell \geq 2} \int_{4m_\pi^2}^{\infty} dt' G_{1,2\ell-1}^1(t,t') \text{Im } g_{2\ell-1}^1(t') + \sum_{\ell} \int_{m_+^2}^{\infty} ds' G_{1,\ell}^-(t,s') \text{Im } f_{\ell}^-(s'), \\
 g_2^0(t) &= \frac{t}{\pi} \int_{4m_\pi^2}^{\infty} \frac{\text{Im } g_2^0(t')}{t'(t'-t)} dt' + \sum_{\ell \geq 2} \int_{4m_\pi^2}^{\infty} \frac{dt'}{t'} G_{2,4\ell-2}^{0'}(t,t') \text{Im } g_{4\ell-2}^0(t') + \sum_{\ell} \int_{m_+^2}^{\infty} ds' G_{2,\ell}^{+'}(t,s') \text{Im } f_{\ell}^+(s').
 \end{aligned} \tag{39}$$

$G_{J,J}^I(t,t')$  = integral kernels, depend on a parameter  
 Lowest # of subtractions. Odd pw decouple from even pw.

$$g_{\ell}^0(t) = \Delta_{\ell}^0(t) + \frac{t}{\pi} \int_{4m_\pi^2}^{\infty} \frac{dt'}{t'} \frac{\text{Im } g_{\ell}^0(t')}{t'-t}, \quad \ell = 0, 2,$$

$$g_1^1(t) = \Delta_1^1(t) + \frac{1}{\pi} \int_{4m_\pi^2}^{\infty} \frac{dt'}{t'-t} \text{Im } g_1^1(t'), \tag{40}$$

$\Delta(t)$  depend on higher waves or on  $K\pi \rightarrow K\pi$ .

Integrals from  $2\pi$  threshold !

Solve in descending J order

We have used models for higher waves, but give very small contributions



For unphysical region below KK threshold, we used Omnés function

$$\Omega_\ell^I(t) = \exp \left( \frac{t}{\pi} \int_{4m_\pi^2}^{t_m} \frac{\phi_\ell^I(t') dt'}{t'(t'-t)} \right),$$

$$\Omega_\ell^I(t) \equiv \Omega_{\ell,R}^I(t) e^{i\phi_\ell^I(t)\theta(t-4m_\pi^2)\theta(t_m-t)},$$

This is the form of our HDR: Roy-Steiner+Omnés formalism

$$g_0^0(t) = \Delta_0^0(t) + \frac{t\Omega_0^0(t)}{t_m-t} \left[ \alpha + \frac{t}{\pi} \int_{4m_\pi^2}^{t_m} dt' \frac{(t_m-t')\Delta_0^0(t') \sin \phi_0^0(t')}{\Omega_{0,R}^0(t')t'^2(t'-t)} + \frac{t}{\pi} \int_{t_m}^{\infty} dt' \frac{(t_m-t')|g_0^0(t')| \sin \phi_0^0(t')}{\Omega_{0,R}^0(t')t'^2(t'-t)} \right]$$

$$g_1^1(t) = \Delta_1^1(t) + \Omega_1^1(t) \left[ \frac{1}{\pi} \int_{4m_\pi^2}^{t_m} dt' \frac{\Delta_1^1(t') \sin \phi_1^1(t')}{\Omega_{1,R}^1(t')(t'-t)} + \frac{1}{\pi} \int_{t_m}^{\infty} dt' \frac{|g_1^1(t')| \sin \phi_1^1(t')}{\Omega_{1,R}^1(t')(t'-t)} \right],$$

$$g_2^0(t) = \Delta_2^0(t) + t\Omega_2^0(t) \left[ \frac{1}{\pi} \int_{4m_\pi^2}^{t_m} dt' \frac{\Delta_2^0(t') \sin \phi_2^0(t')}{\Omega_{2,R}^0(t')t'(t'-t)} + \frac{1}{\pi} \int_{t_m}^{\infty} dt' \frac{|g_2^0(t')| \sin \phi_2^0(t')}{\Omega_{2,R}^0(t')t'(t'-t)} \right].$$

We can now check how well these HDR are satisfied

# Our Dispersive/Analytic Approach for $\pi K$ and strange resonances

Simple Unconstrained Fits to  $\pi K$  partial-wave Data (UFD).

Estimation of statistical and SYSTEMATIC errors

## Forward Dispersion Relations:

Left cut easy to rewrite

Relate amplitudes, not partial waves

Not direct access to pole

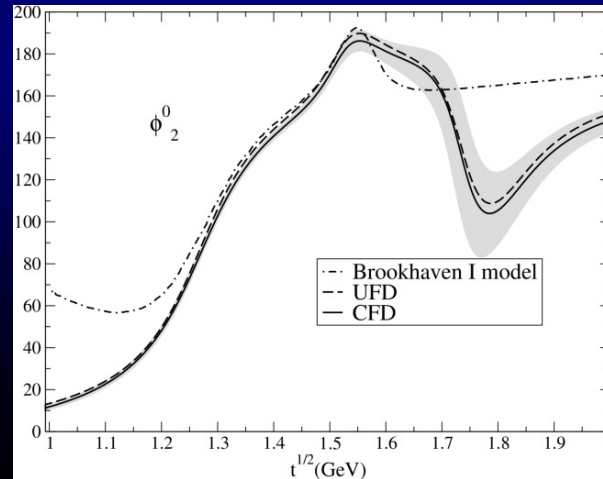
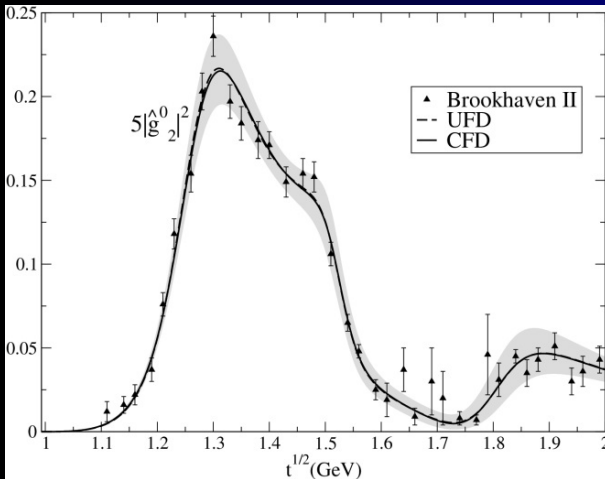
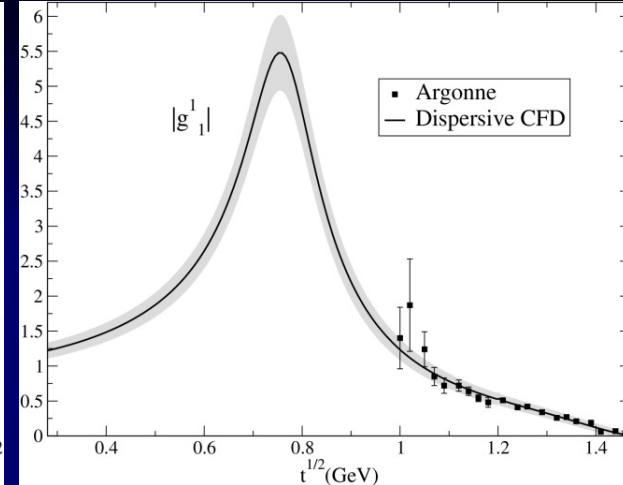
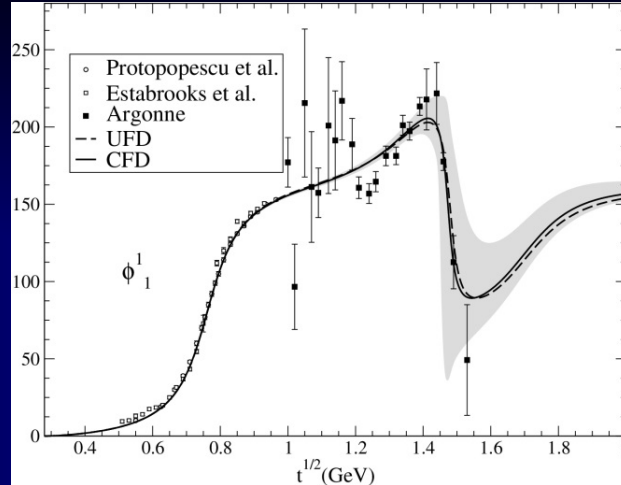
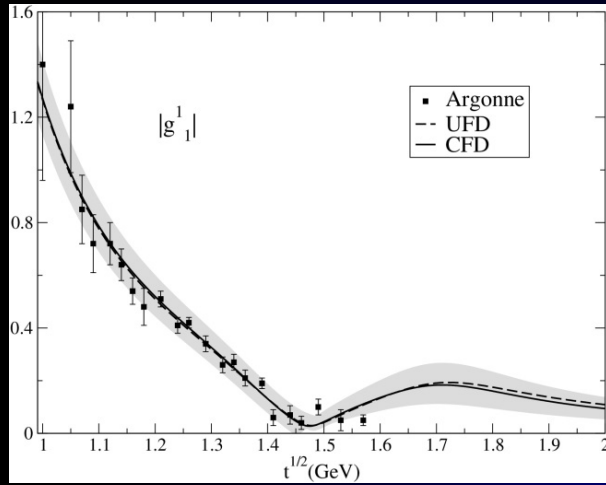
- As  $\pi K$  checks: Small inconsistencies.
- As constraints:  
 **$\pi K$  consistent fits up to 1.6 GeV** JRP, A.Rodas, Phys.Rev. D93 (2016)
- Analytic methods to extract poles: reduced model dependence on strange resonances  
JRP, A. Rodas. J. Ruiz de Elvira, Eur.Phys.J. C77 (2017)

## Partial-wave $\pi K$ Dispersion Relations

Need  $\pi\pi \rightarrow KK$  to rewrite left cut. Range optimized.

- As  $\pi\pi \rightarrow KK$  checks: Small inconsistencies.
- As constraints:  
 **$\pi\pi \rightarrow KK$  consistent fits up to 1.5 GeV** JRP, A.Rodas, Eur.Phys.J. C78 (2018)

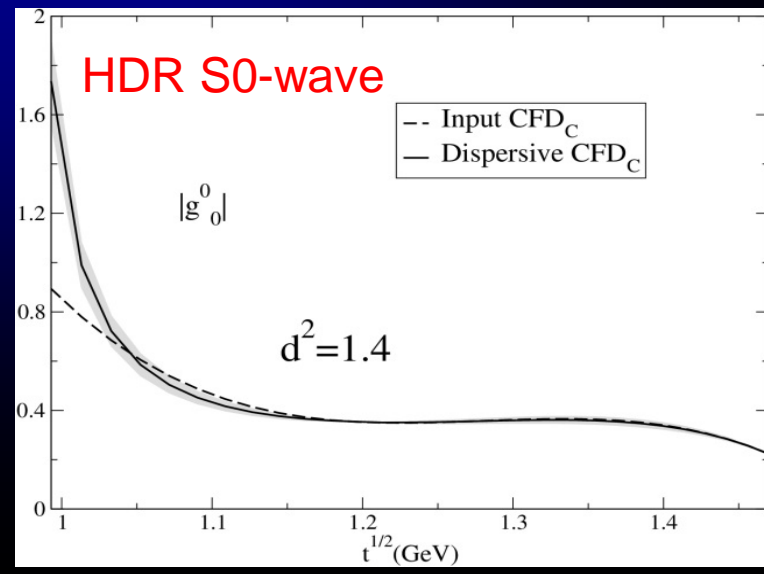
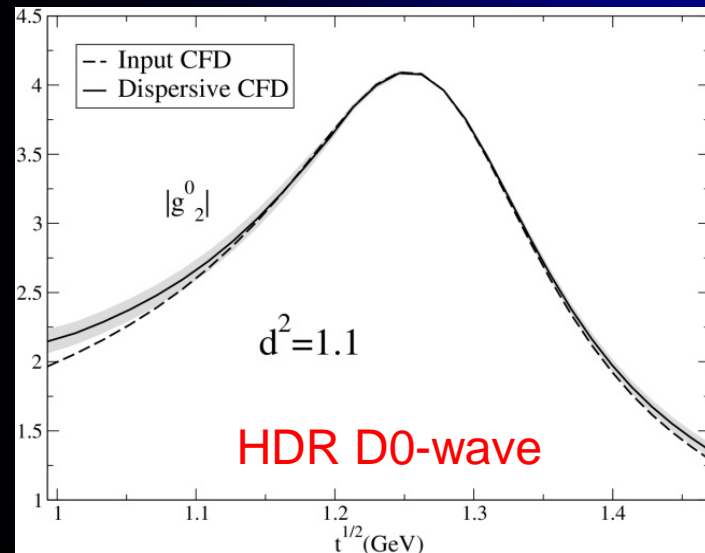
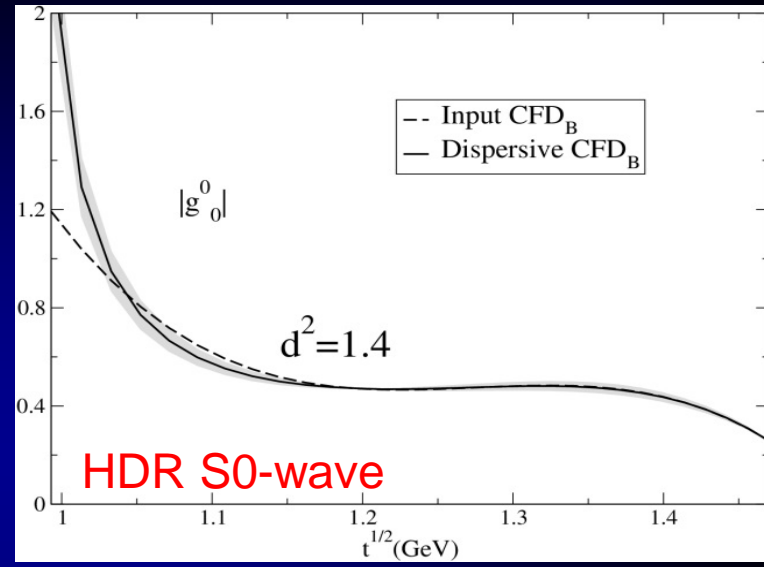
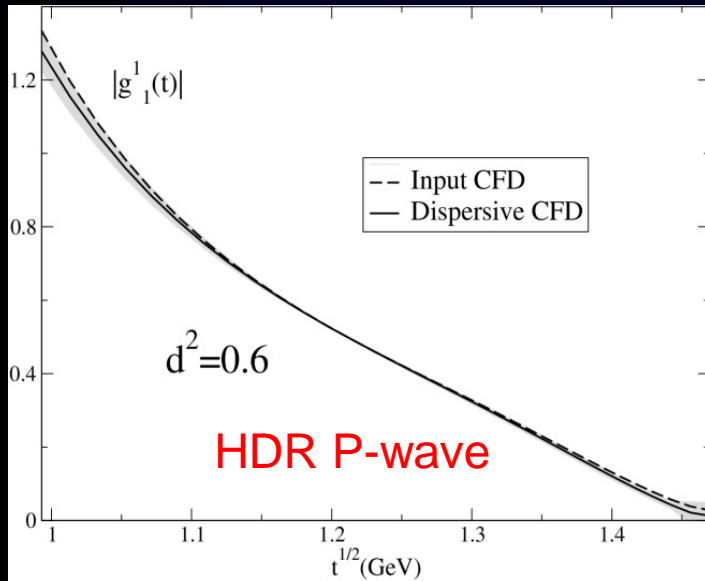
Once again we started with SIMPLE FITS TO  $\pi\pi \rightarrow KK$  DATA, including systematic uncertainties



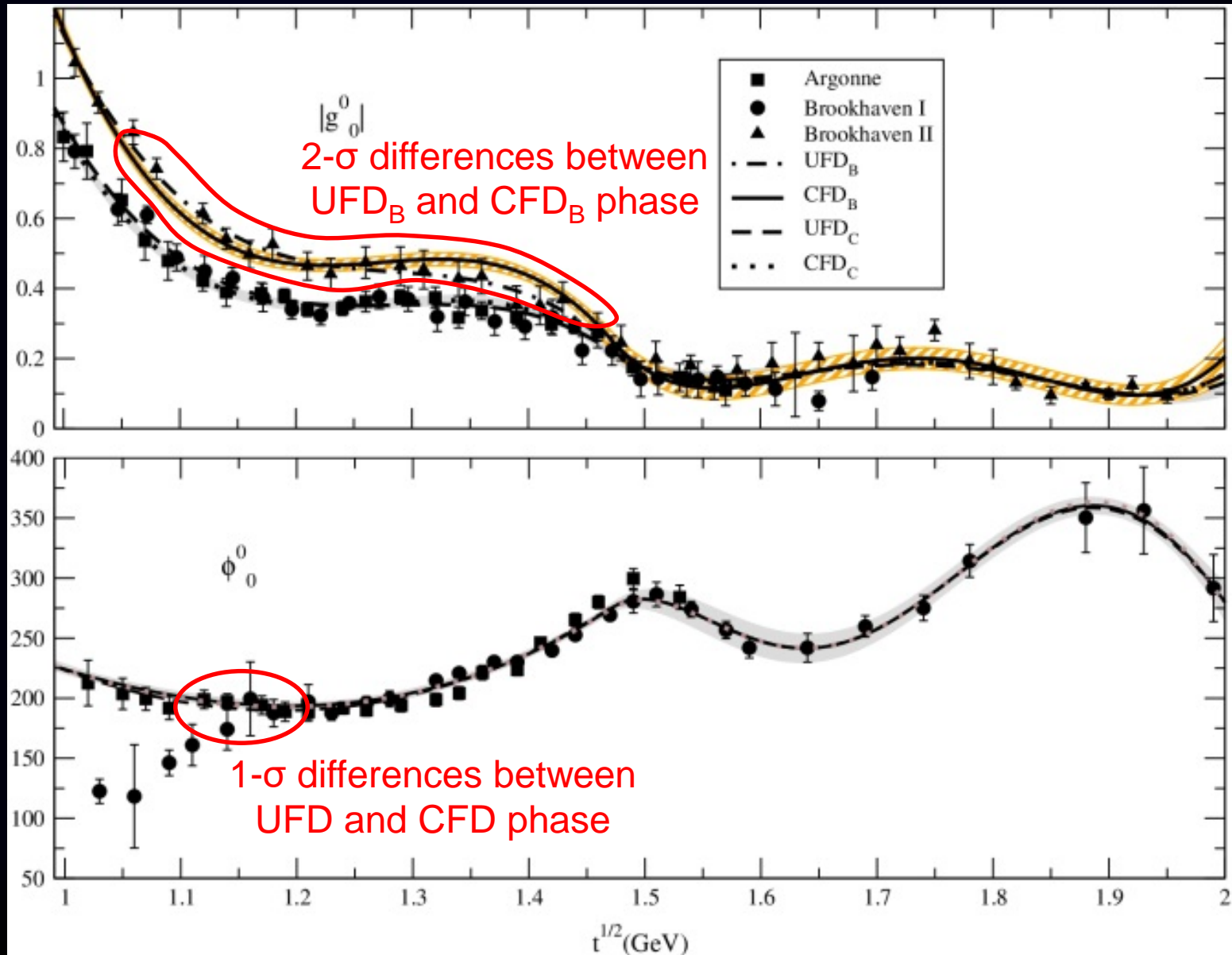
UFD Inconsistent with HDR  
If not constrained

But consistent after HDR used as constraints

Two possible solutions for S0 wave



Some  $2\text{-}\sigma$  level differences between  $\text{UFD}_B$  and  $\text{CFD}_B$  between 1.05 and 1.45 GeV  
 $\text{CFD}_C$  consistent within  $1\text{-}\sigma$  band of  $\text{UFD}_C$



# Our Dispersive/Analytic Approach for $\pi K$ and strange resonances

## Simple Unconstrained Fits to $\pi K$ partial-wave Data (UFD).

Estimation of statistical and SYSTEMATIC errors

### Forward Dispersion Relations:

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### Partial-wave $\pi K$ Dispersion Relations

Need  $\pi\pi \rightarrow KK$  to rewrite left cut. Range optimized.

- From fixed-t DR:  
 $\pi\pi \rightarrow KK$  influence small.  
 $\kappa/K_0^*(700)$  out of reach

- From Hyperbolic DR:  
 $\pi\pi \rightarrow KK$  influence important.

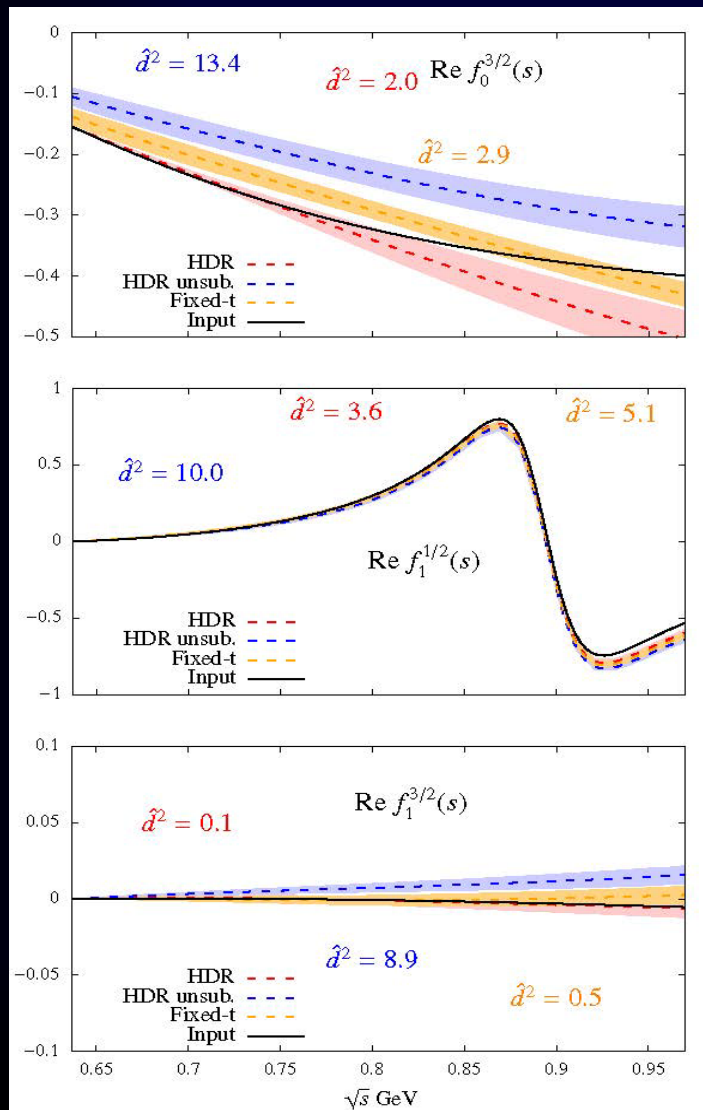
JRP, A.Rodas, in progress. PRELIMINARY results shown here

- As  $\pi\pi \rightarrow KK$  checks: Small inconsistencies.
- As constraints:  
 **$\pi\pi \rightarrow KK$  consistent fits up to 1.5 GeV**  
JRP, A.Rodas, Eur.Phys.J. C78 (2018)

- As  $\pi K$  Checks: Large inconsistencies.

# LARGE inconsistencies IF UNCONSTRAINED

## Unconstrained Fit to Data



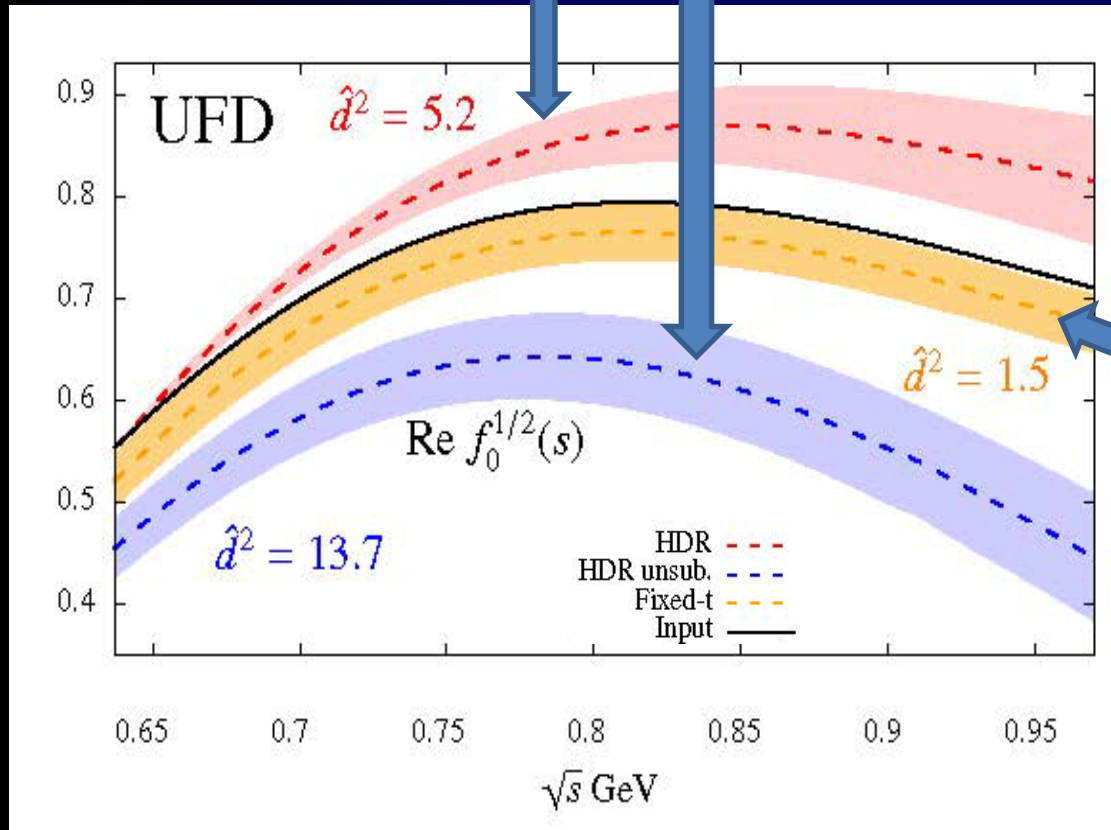


# $\pi K$ Hyperbolic Dispersion Relations $I=1/2, J=0$

The most relevant wave for the kappa resonance.

**LARGE inconsistencies with HDR Roy-Steiner from unconstrained fits (UFD)**

One or no subtraction for  $F^-$  lie on opposite sides of input



Fixed-t Roy-Steiner is fair but kappa pole outside their applicability region

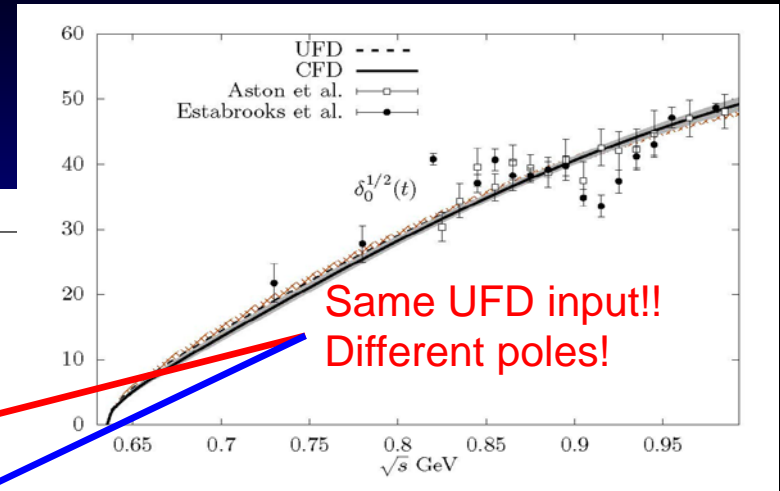
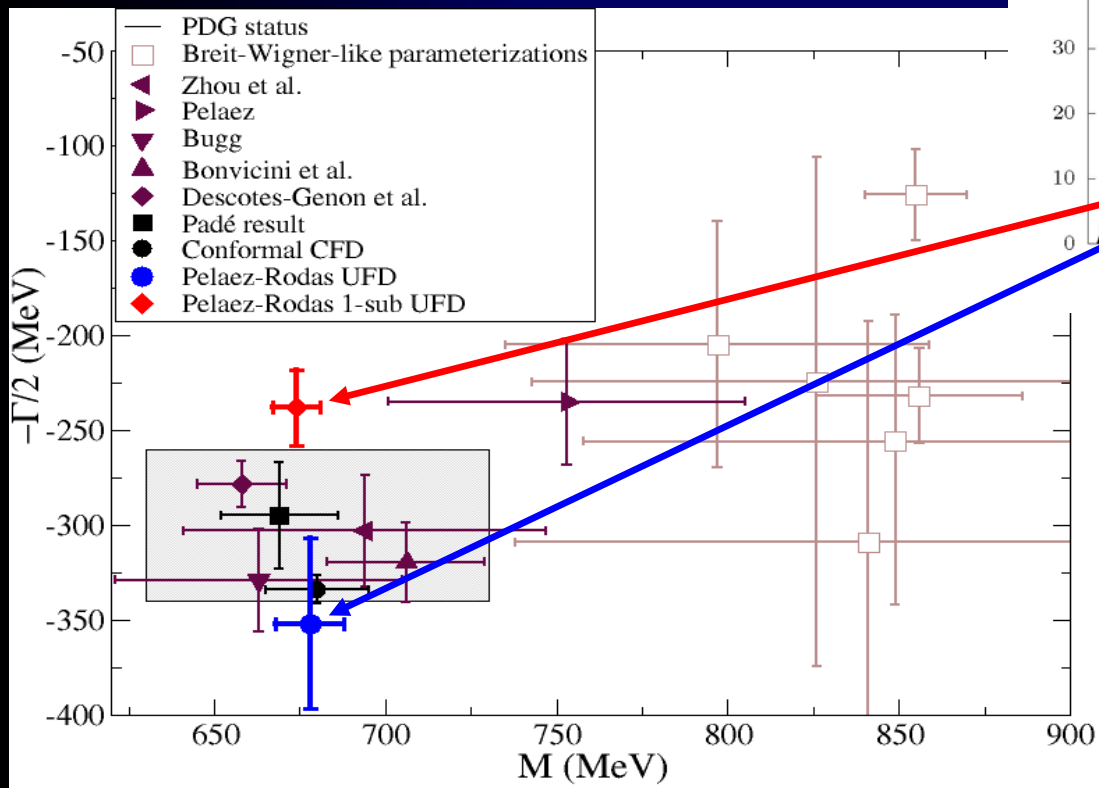
We have chosen the hyperbolae family so that the kappa pole and its uncertainties lie within their applicability region

# WARNING ABOUT THE PRECISION OF UNCONSTRAINED FITS

Before imposing Roy Eqs. incompatible results with different # of subtractions !!

This is partly due to left/circular cuts.

(Crossed Channel)



Nice-looking fits are NOT enough to get a stable and precise continuation to the complex plane

You can imagine what precision you get if you use simple models only of  $\pi K$ , without left cut or without dispersion relations...

# Our Dispersive/Analytic Approach for $\pi K$ and strange resonances

Simple Unconstrained Fits to  $\pi K$  partial-wave Data (UFD).

Estimation of statistical and SYSTEMATIC errors

## Forward Dispersion Relations:

Left cut easy to rewrite  
Relate amplitudes, not partial waves  
Not direct access to pole

- As  $\pi K$  checks: Small inconsistencies.
- As constraints:  
 **$\pi K$  consistent fits up to 1.6 GeV** JRP, A.Rodas, Phys.Rev. D93 (2016)
- Analytic methods to extract poles: reduced model dependence on strange resonances  
JRP, A. Rodas. J. Ruiz de Elvira, Eur.Phys.J. C77 (2017)

## Partial-wave $\pi K$ Dispersion Relations

Need  $\pi\pi \rightarrow KK$  to rewrite left cut. Range optimized.

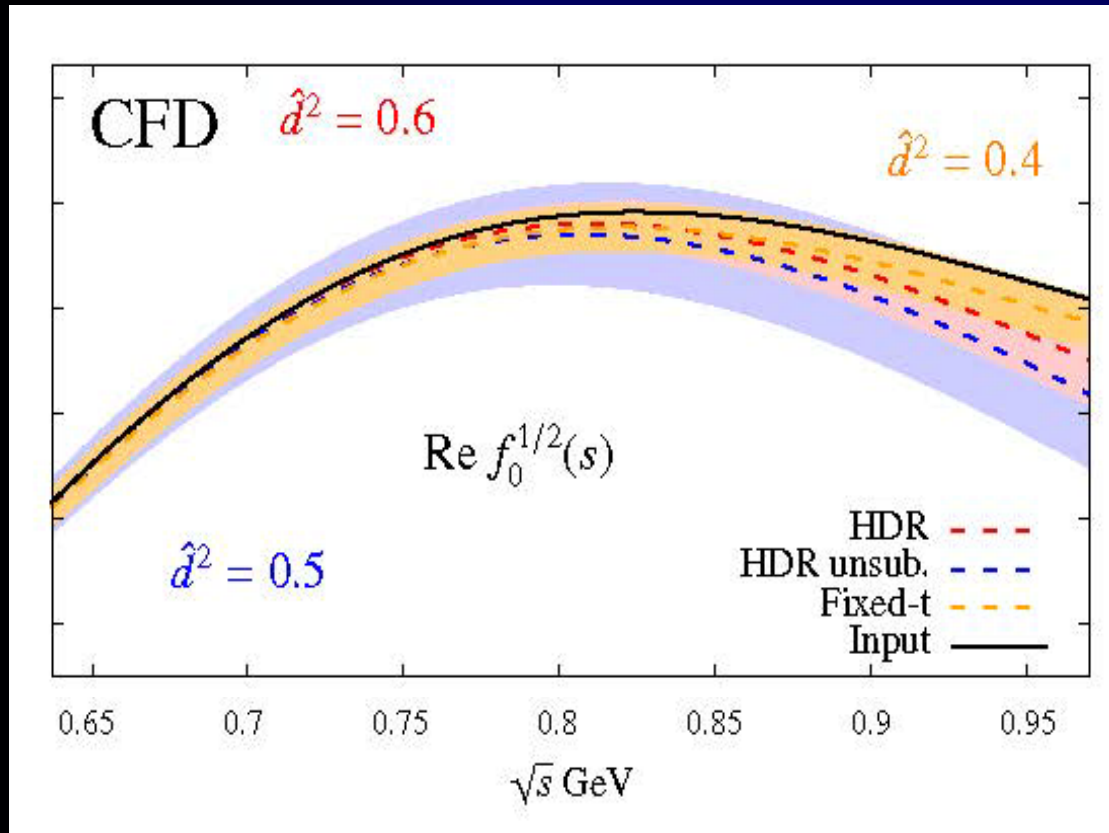
- From fixed-t DR:  
 $\pi\pi \rightarrow KK$  influence small.  
 $k/K_0^*(700)$  out of reach
- From Hyperbolic DR:  
 $\pi\pi \rightarrow KK$  influence important.

JRP, A.Rodas,  
arXiv:2010.1122.  
To appear in Physics  
Reports

- As  $\pi\pi \rightarrow KK$  checks: Small inconsistencies.
  - As constraints:  
 **$\pi\pi \rightarrow KK$  consistent fits up to 1.5 GeV**  
JRP, A.Rodas, Eur.Phys.J. C78 (2018)
- 
- As  $\pi K$  Checks: Large inconsistencies.
  - **ALL DR TOGETHER** as Constraints:  
 **$\pi K$  consistent fits up to 1.1 GeV**

We provide a constrained fit to data (CFD) satisfying 16 Dispersion relations (FDRs, fixed-t, HDR, different # subtractions)

Fairly simple and ready to use parameterizations



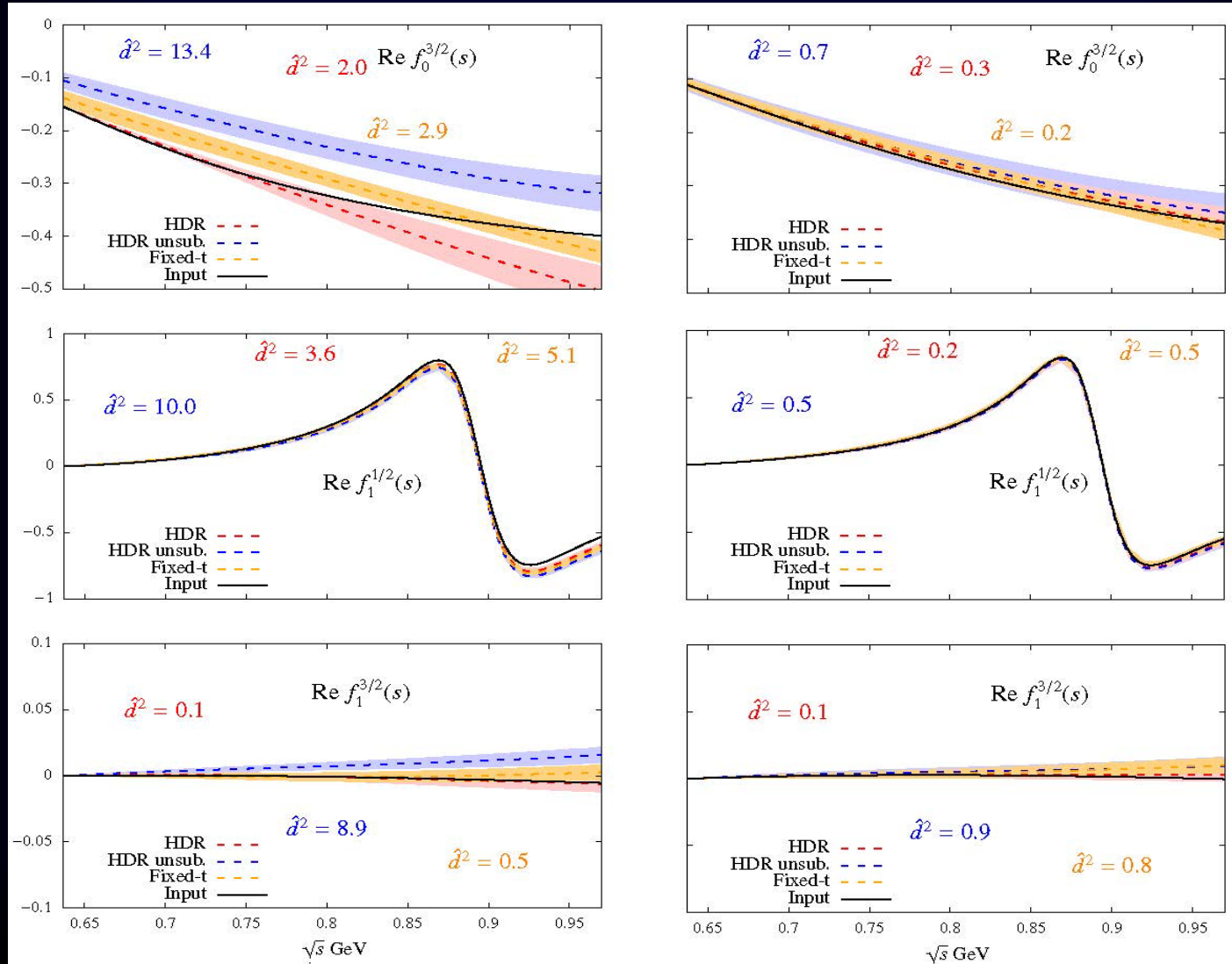
Our Constrained parameterization now yields consistent output for all Dispersion Relations

# LARGE inconsistencies FOR THE OTHER WAVES IF UNCONSTRAINED

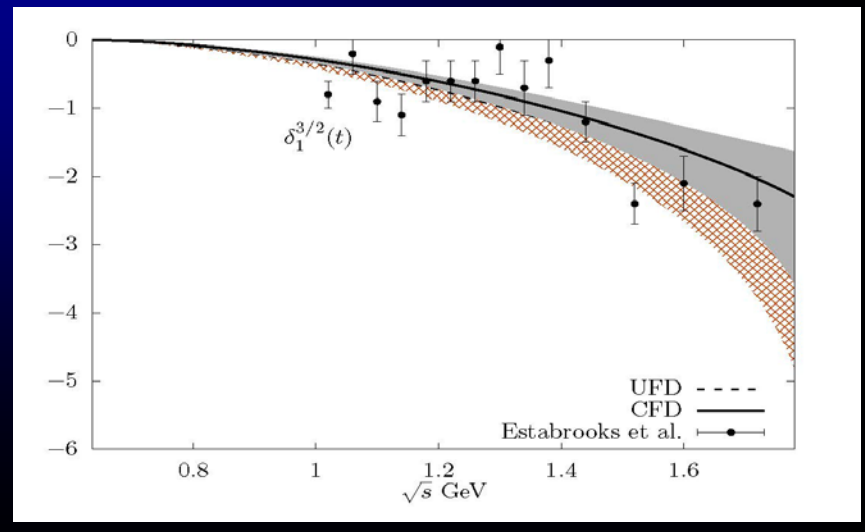
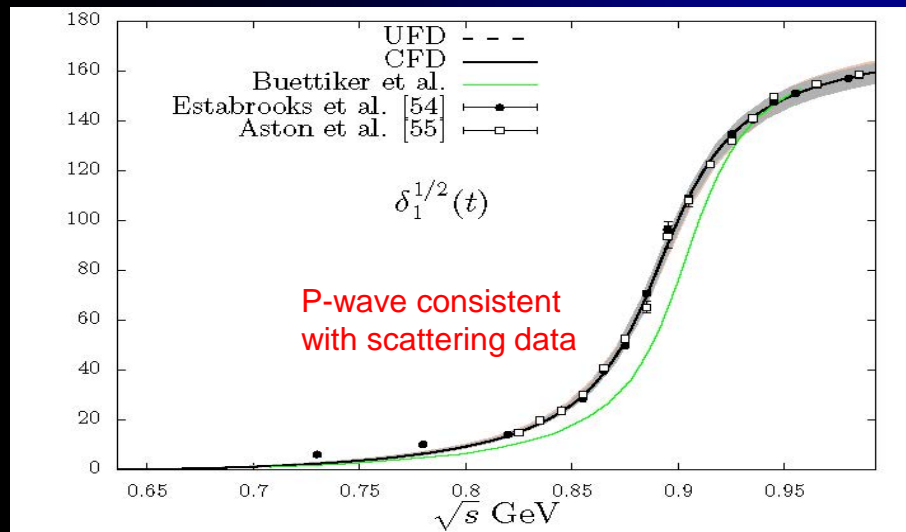
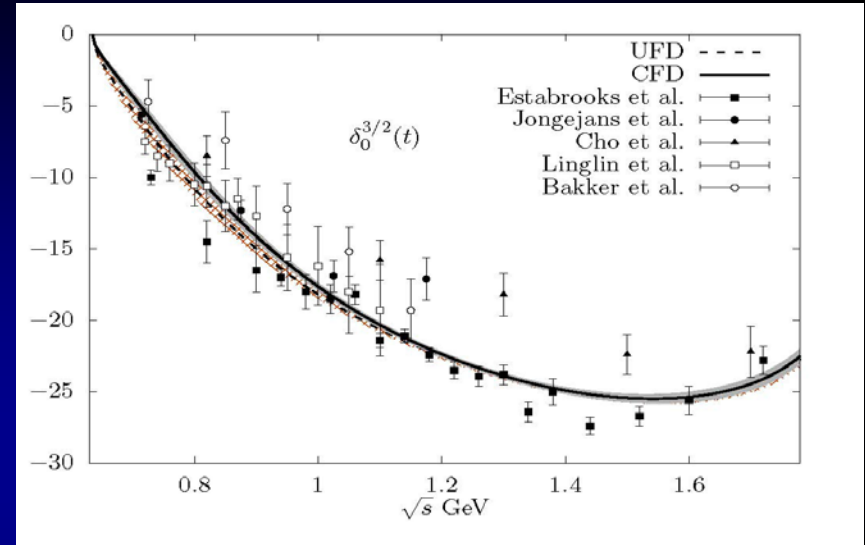
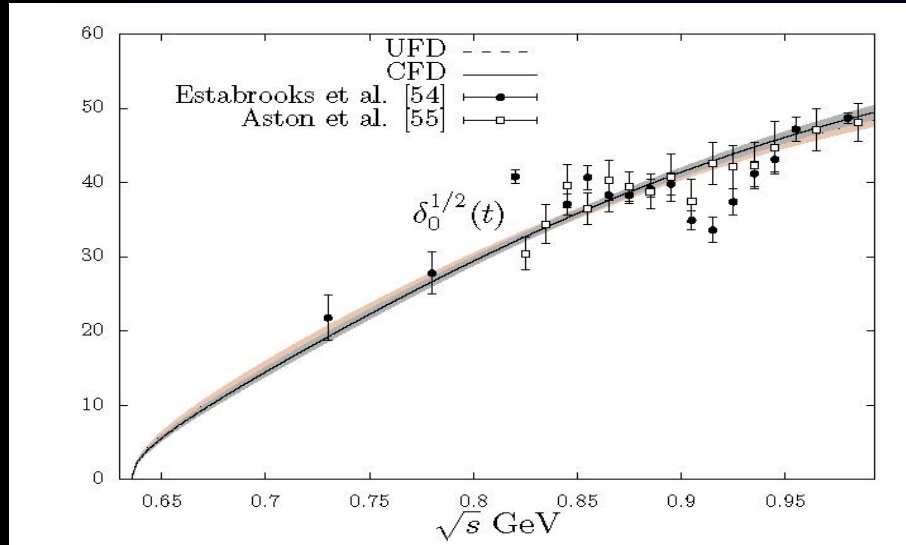
Made consistent within uncertainties for the CFD

Unconstrained Fit to Data

Constrained Fit to Data



Constrained parameterizations suffer minor changes but still describe  $\pi K$  data fairly well. Here we compare the unconstrained fits (UFD) versus the constrained ones (CFD)



The "unphysical" rho peak in  $\pi\pi \rightarrow KK$  grows by 10% from UFD to CFD



# Our Dispersive/Analytic Approach for $\pi K$ and strange resonances

Simple Unconstrained Fits to  $\pi K$  partial-wave Data (UFD).

Estimation of statistical and SYSTEMATIC errors

## Forward Dispersion Relations:

Left cut easy to rewrite

Relate amplitudes, not partial waves

Not direct access to pole

- As  $\pi K$  checks: Small inconsistencies.
- As constraints:  
 **$\pi K$  consistent fits up to 1.6 GeV**
- Padé sequences to extract poles from local information: reduced model dependence on strange resonances

JRP, A.Rodas,  
Phys.Rev. D93 (2016)

JRP, A. Rodas. J. Ruiz de Elvira, Eur.Phys.J. C77 (2017)

## Partial-wave $\pi K$ Dispersion Relations (PWDR)

Need  $\pi\pi \rightarrow KK$  to rewrite left cut. Range optimized.

- As  $\pi\pi \rightarrow KK$  checks: Small inconsistencies.
- As constraints:  
 **$\pi\pi \rightarrow KK$  consistent fits from KK threshold to 1.5 GeV**

JRP, A.Rodas,  
Eur.Phys.J. C78 (2018)

- From fixed-t DR:  
 $\pi\pi \rightarrow KK$  influence small.  
 $\kappa/K_0^*(700)$  pole out of reach
- From Hyperbolic DR:  
 $\pi\pi \rightarrow KK$  influence important.  
As  $\pi K$  Checks:  
Large inconsistencies

- **ALL DR TOGETHER** as Constraints:  
 $\pi K$  consistent fits up to 1.1 GeV for PWDR,  
up to 1.6 for FDRs,  
 $\pi\pi \rightarrow KK$  up to 1.5 GeV and unphysical region
- **Precise  $\pi K$  threshold parameters**

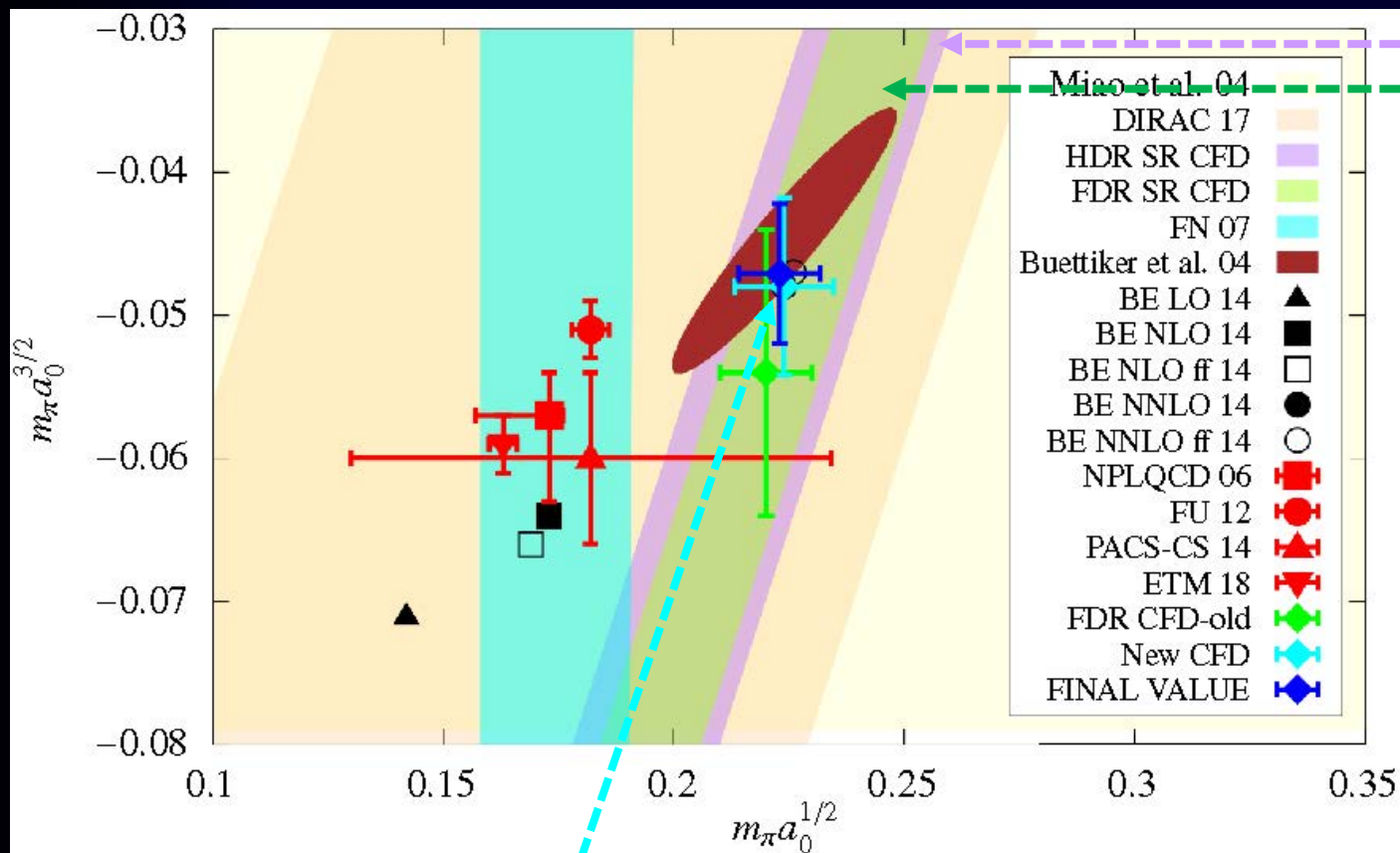
JRP, A.Rodas,

arXiv:2010.1122.

To appear in Physics  
Reports



- Threshold parameters relevant to test ChPT (NNLO at present).
- Present tension between lattice and dispersive results



• Our Dispersive  
SUM RULES  
for  $a_0^-$

Our dispersively Constrained  
Fit to DATA (CFD)

Table 25:  $S$ -wave scattering lengths ( $m_\pi$  units).

	UFD	CFD	Ref. [43]
$a_0^{1/2}$	$0.241 \pm 0.012$	<b><math>0.224 \pm 0.011</math></b>	$0.224 \pm 0.022$
$a_0^{3/2}$	$-0.067 \pm 0.012$	<b><math>-0.048 \pm 0.006</math></b>	$-0.0448 \pm 0.0077$

- We provide sum rule values for scattering lengths and slopes up to D-waves.
- Good consistency with CFD for S,P waves (constrained) and D-wave lengths

	This work sum rules with CFD input				This work direct CFD	Sum rules [43] Fixed- $t$	NNLO ChPT [85] and [86]*
	Fixed- $t$	HDR	HDR <sub>sub</sub>	Final Value			
$m_\pi a_0^{1/2}$	0.224±0.009	0.221± 0.012	like CFD	<b>0.223±0.009</b>	0.224±0.011	0.224±0.022	0.224*
$m_\pi^3 b_0^{1/2} \times 10$	1.04± 0.04	1.05±0.07	1.15± 0.04	<b>1.08±0.08</b>	0.95±0.04	0.85±0.04	1.278
$m_\pi a_0^{3/2} \times 10$	-0.478± 0.052	-0.460±0.064	like CFD	<b>-0.471±0.049</b>	-0.48±0.06	-0.448±0.077	-0.471*
$m_\pi^3 b_0^{3/2} \times 10$	-0.42±0.02	-0.41±0.03	-0.44±0.02	<b>-0.43±0.03</b>	-0.36±0.04	-0.37±0.03	-0.326
$m_\pi^3 a_1^{1/2} \times 10$	0.228±0.010	0.218±0.008	0.222±0.006	<b>0.222±0.009</b>	0.20±0.04	0.19±0.01	0.152
$m_\pi^5 b_1^{1/2} \times 10^2$	0.58±0.03	0.59±0.03	0.60±0.03	<b>0.59±0.02</b>	0.5±0.2	0.18±0.02	0.032
$m_\pi^3 a_1^{3/2} \times 10^2$	0.15±0.05	0.19±0.05	0.17±0.04	<b>0.17±0.05</b>	0.15±0.11	0.065±0.044	0.293
$m_\pi^5 b_1^{3/2} \times 10^3$	-0.94±0.09	-0.97±0.08	-1.03±0.07	<b>-0.99±0.09</b>	-1.04±0.8	-0.92±0.17	0.544
$m_\pi^5 a_2^{1/2} \times 10^3$	0.60±0.13	0.54±0.03	0.55±0.02	<b>0.55±0.05</b>	0.53±0.05	0.47±0.03	0.142
$m_\pi^7 b_2^{1/2} \times 10^4$	-0.89±0.10	-0.96±0.09	-0.95±0.09	<b>-0.94±0.09</b>	0.20±0.02	-1.4±0.3	-1.98
$m_\pi^5 a_2^{3/2} \times 10^4$	-0.05±0.60	-0.11±0.16	-0.18±0.15	<b>-0.14±0.17</b>	-0.09±0.03	-0.11±0.27	-0.45
$m_\pi^7 b_2^{3/2} \times 10^4$	-1.12±0.10	-1.13±0.09	-1.14±0.09	<b>-1.13±0.06</b>	-0.03±0.01	-0.96±0.26	0.61

## Simple Unconstrained Fits to $\pi K$ partial-wave Data (UFD).

Estimation of statistical and SYSTEMATIC errors

### Forward Dispersion Relations:

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Not direct access to pole

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JRP, A.Rodas,  
Phys.Rev. D93 (2016)

JRP, A. Rodas. J. Ruiz de Elvira, Eur.Phys.J. C77 (2017)

### Partial-wave $\pi K$ Dispersion Relations (PWDR)

Need  $\pi\pi \rightarrow KK$  to rewrite left cut. Range optimized.

- As  $\pi\pi \rightarrow KK$  checks: Small inconsistencies.
- As constraints:  
 **$\pi\pi \rightarrow KK$  consistent fits from KK threshold to 1.5 GeV**

JRP, A.Rodas,  
Eur.Phys.J. C78 (2018)

- From fixed-t DR:  
 $\pi\pi \rightarrow KK$  influence small.  
 $\kappa/K_0^*(700)$  pole out of reach
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 $\pi\pi \rightarrow KK$  influence important.  
As  $\pi K$  Checks:  
Large inconsistencies

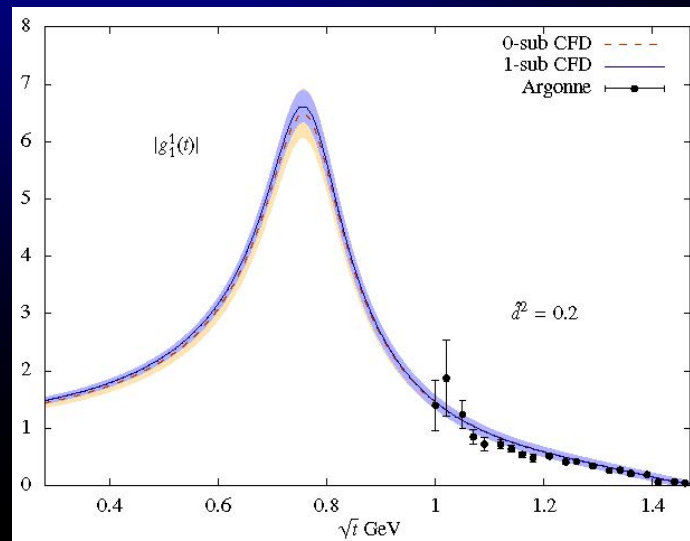
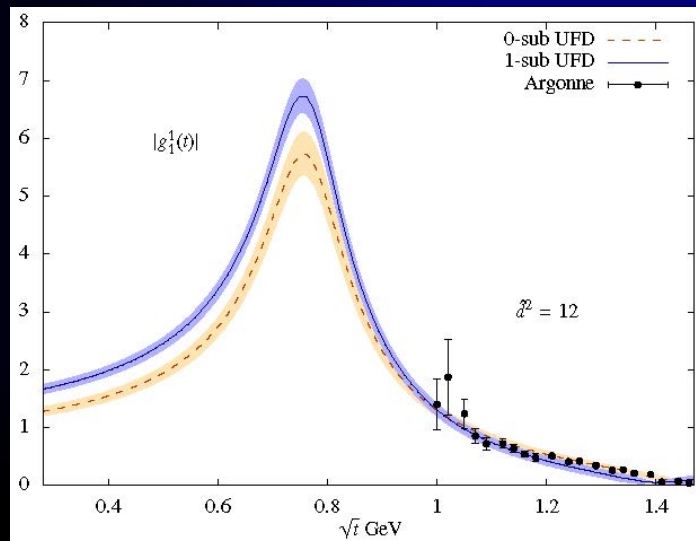
- **ALL DR TOGETHER** as Constraints:  
 $\pi K$  consistent fits up to 1.1 GeV for PWDR, up to 1.6 for FDRs,  
 $\pi\pi \rightarrow KK$  up to 1.5 GeV and unphysical region
- **Precise  $\pi K$  threshold parameters**
- **Rigorous  $\kappa/K_0^*(700)$  pole**

JRP, A.Rodas,  
PRL. 124 (2020) 17, 172001

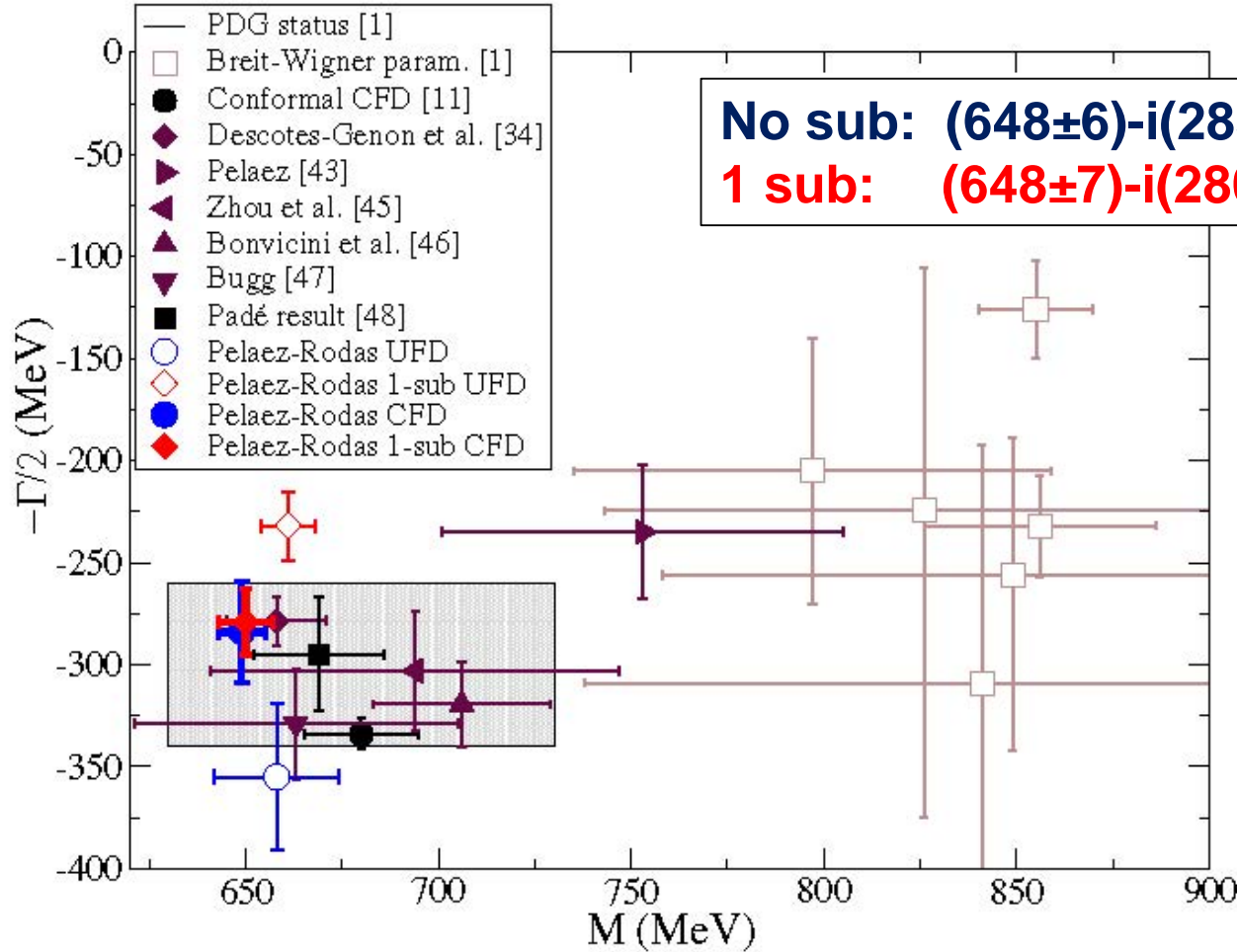
JRP, A.Rodas,  
arXiv:2010.1122.  
To appear in Physics  
Reports

Now we have:

- **FIT TO DATA** (not solution but fit) **CONSTRAINED WITH 16 DR**
- Improved  $P^{1/2}$ -wave (consistent with data) and  $P^{3/2}$
- Improved Pomeron
- Realistic  $\pi\pi \rightarrow KK$  uncertainties (none before)
- Constrained  $\pi\pi \rightarrow KK$  input with DR
- FDR up to 1.6 GeV
- Fixed-t Roy-Steiner Eqs.
- Hyperbolic Roy Steiner Eqs.
  - Both one and no-subtractions for F- HDR (only the subtracted one before)
  - both in real axis (not HDR before) and complex plane
  - Unphysical P-wave  $\pi\pi \rightarrow KK$  region VERY RELEVANT



When using the constrained fit to data both poles come out nicely compatible



Compatible with Paris group

Decotes-Genon-Moussallam 2006  
 $(658 \pm 13) - i(278.5 \pm 12)$  MeV

And with our previous “Pade sequence” determination  
 $(670 \pm 18) - i(295 \pm 28)$  MeV

JRP, A. Rodas. J. Ruiz de Elvira, Eur.Phys.J. C77 (2017)

## Summary

- $\pi K$  and  $\pi\pi \rightarrow KK$  data do not satisfy well basic dispersive constraints
- Using dispersion relations as constraints we provide **simple** and consistent data parameterizations.
- We have implemented partial-wave dispersion relations whose applicability range reaches the kappa pole.
- We have also derived and used SUM RULES to obtain precise threshold parameters
- We confirm previous studies and provide a precise determination of the  $\kappa/K_0^*(700)$  parameters **FROM DATA. A good control on the left/circular cuts is needed to claim this precision.**
- This resonance will be considered “well-established” in next RPP, completing the nonet of lightest scalars.



## Long way since 1966 TO DO LIST

1. The  $\kappa(725)$  (Lynch, Rittenberg, Rosenfeld, Söding, Dec. 1966)

We are beginning to think that  $\kappa$  should be classified along with flying saucers, the Loch Ness Monster, and the Abominable Snowman. We have heard of several experiments which were supposed to confirm it, and each

Confirm the  $\kappa/K_0^*(700)$



At last @PDG 2021\* !!

\* C. Hanhart, private communication

## OUTLOOK

Confirm flying saucers

Confirm Nessie

Abominable Snowman

Work in progress.... stay tuned!

Thank you!