

Accretion disk models in diverse environments

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Star and planet formation

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6th Summer School, Life Cycle of Dust



INSTITUTO DE
ASTROFÍSICA DE
ANDALUCÍA



EXCELENCIA
SEVERO
OCHOA

spfe star planet formation and evolution

By studying the first and last phases of a star, one of our objectives is to understand which the limits are as regards planetary formation conditions and their survival, and how this can be explained within the diversity of the exoplanets observed in the surroundings of the sun.

Research

People

Funded by AYA2017-84390-C2-1-R grant of the AEI, Spain (co-funded by FEDER)

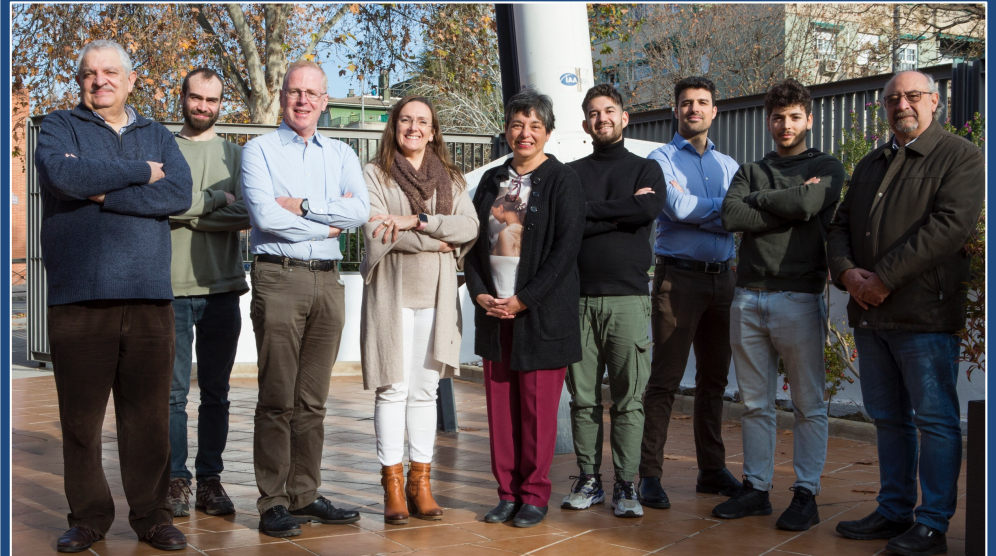
Funded by Junta de Andalucía (Spain) grant P20-00880 (FEDER, EU)



Star and Planet Formation teams IAA+UB+ICE(Institute of Space Sciences)



Barcelona



Granada

Outline

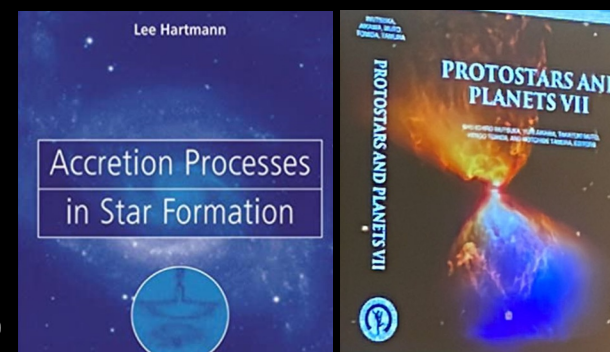
- **First part:** Fundamental properties of accretion disks, applications of α disk models to the observations.
- **Second part:** Examples of models in peculiar accretion disks: Planet formation in extreme conditions.

Bibliography

Accretion processes in Star Formation (Hartmann's book)

Protostars and Planets VII, 2023.

Planet formation in extreme conditions (M. Osorio), Proceedings of the XV Scientific Meeting of the Spanish Astronomical Society



Star Formation

Interstellar Molecular Cloud

Stars

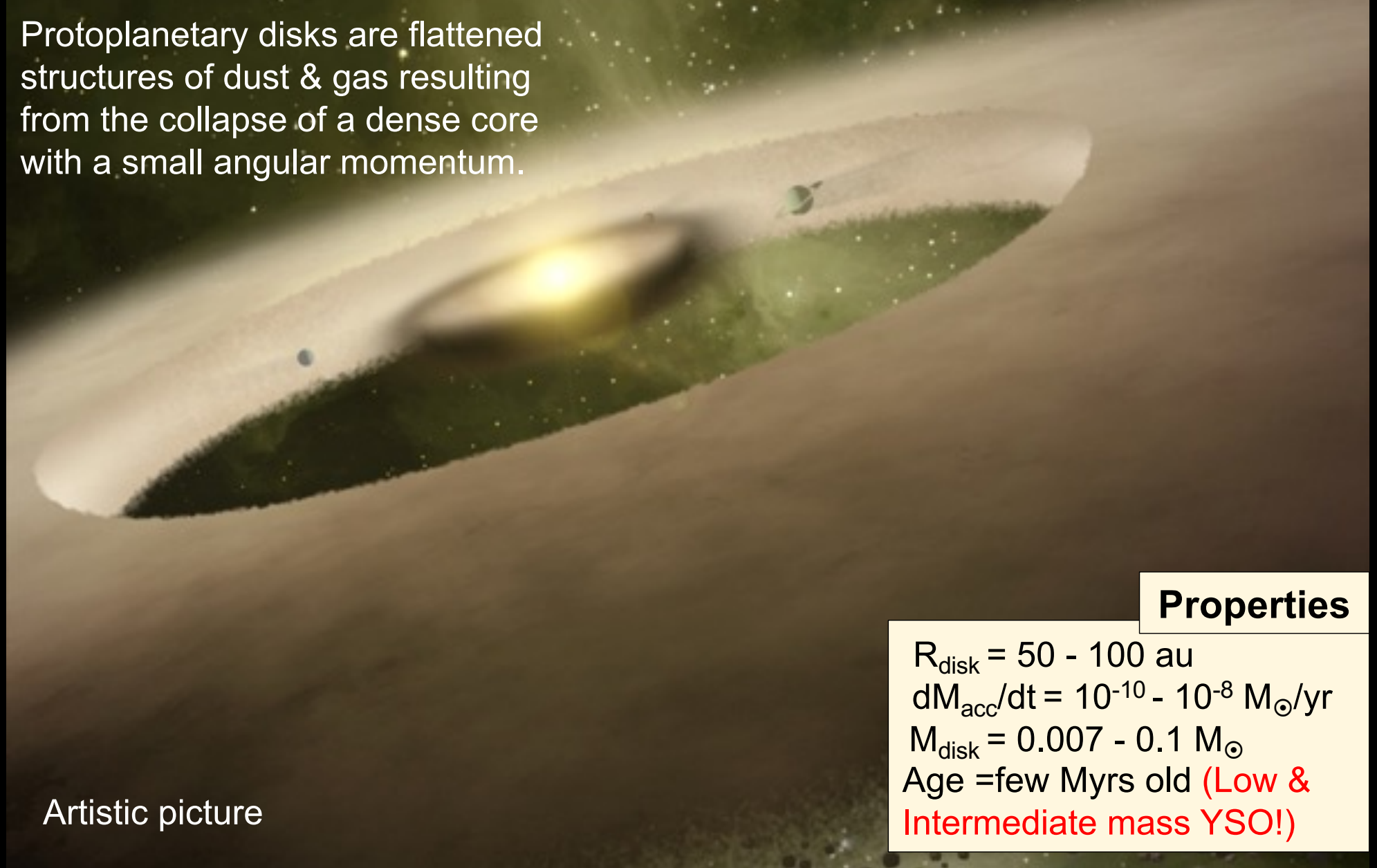
Size: $\times 10^{-7}$
Density: $\times 10^{20}$
Temperature: $\times 10^6$

Planets



Progenitors of the exoplanetary systems: Protoplanetary disks.

Protoplanetary disks are flattened structures of dust & gas resulting from the collapse of a dense core with a small angular momentum.



Properties

$$R_{\text{disk}} = 50 - 100 \text{ au}$$

$$dM_{\text{acc}}/dt = 10^{-10} - 10^{-8} M_{\odot}/\text{yr}$$

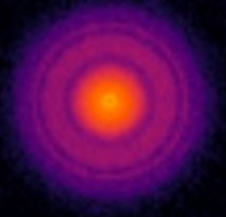
$$M_{\text{disk}} = 0.007 - 0.1 M_{\odot}$$

Age = few Myrs old (**Low & Intermediate mass YSO!**)

Artistic picture

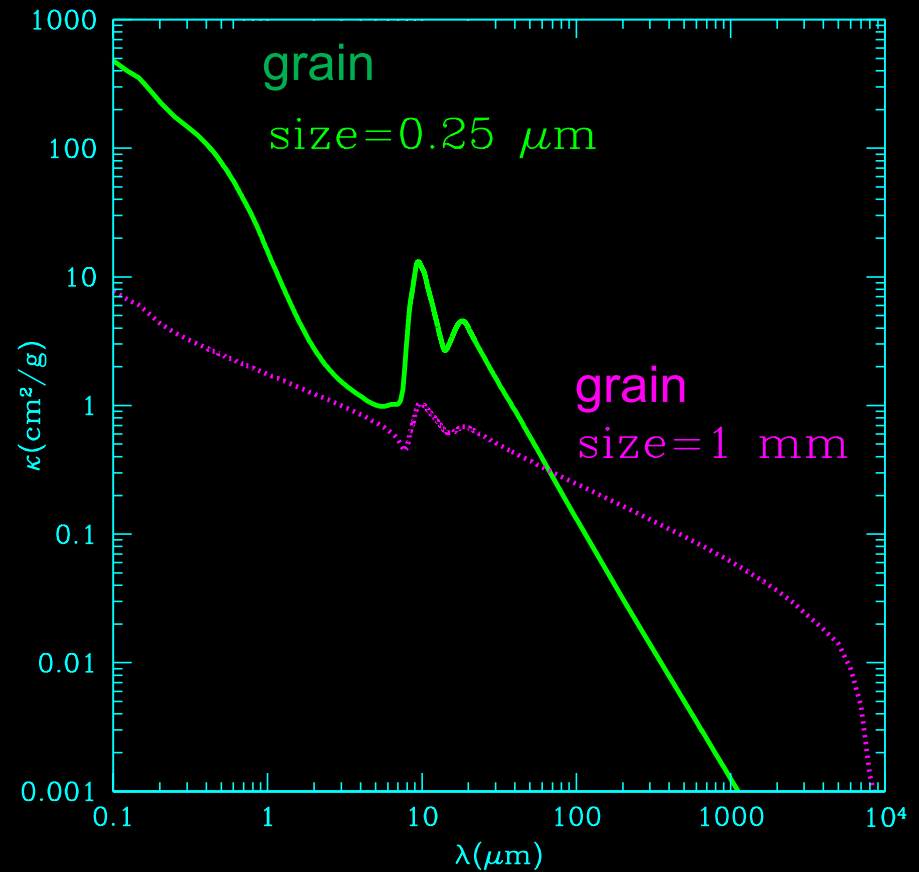
Disks are observed in dust thermal emission (radio observations)

TW Hydra
0.9 mm

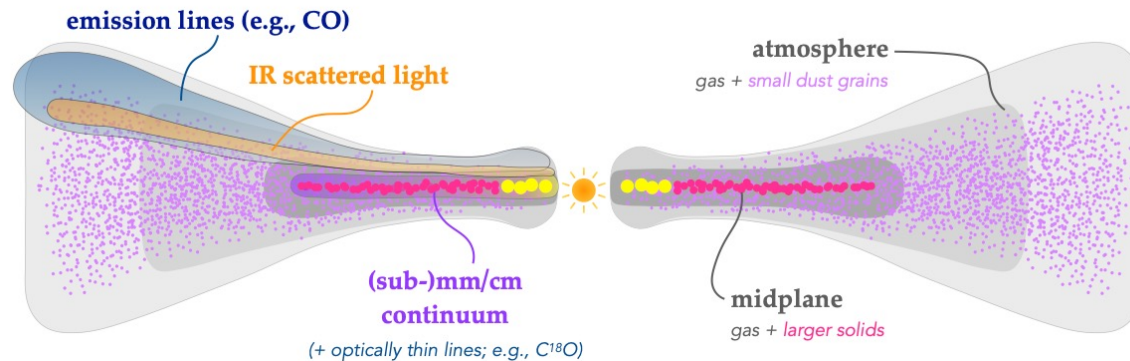


2"

120 au

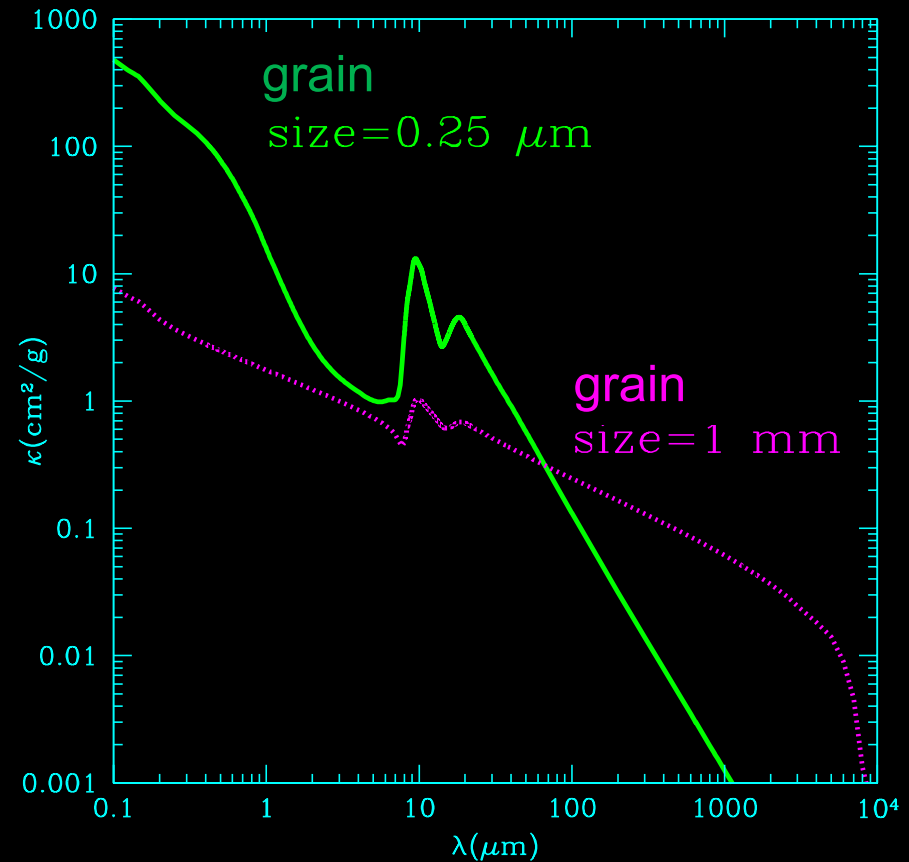
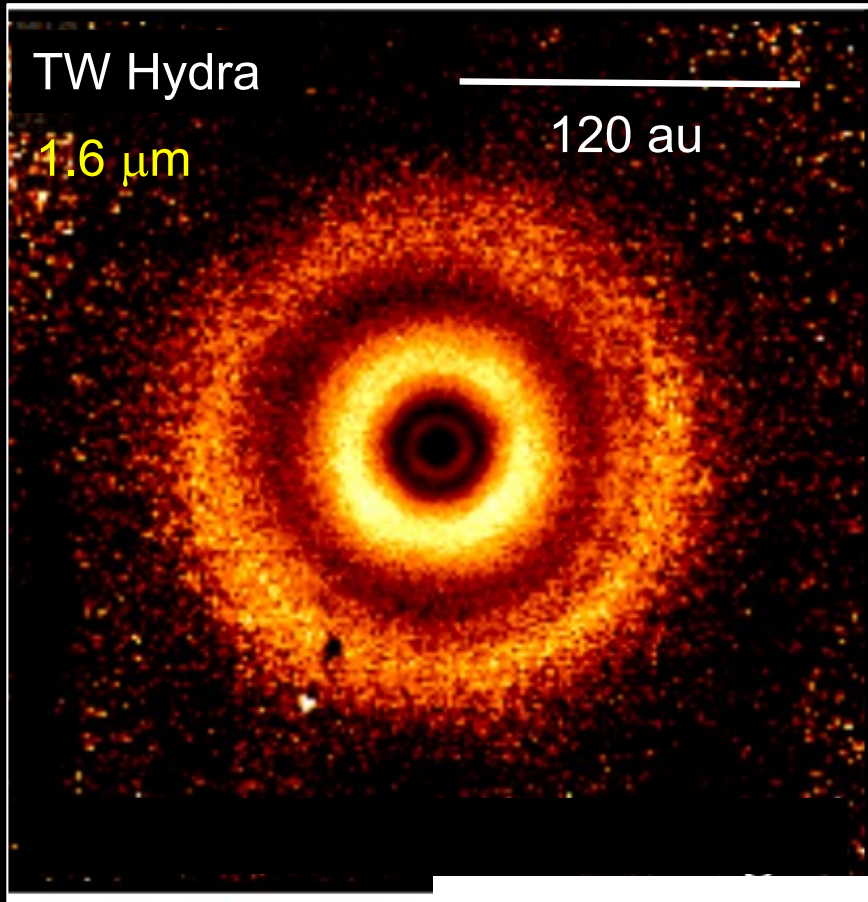


Andrews+2016

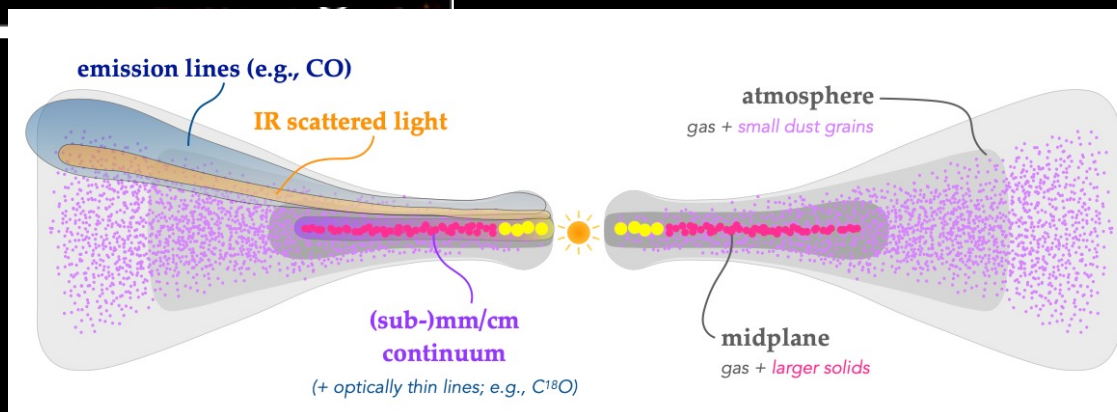
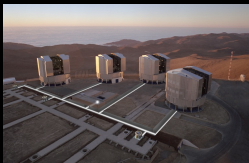


schematic cross section of the disk

Disks are observed in IR scattered light (Near Infrared observations)



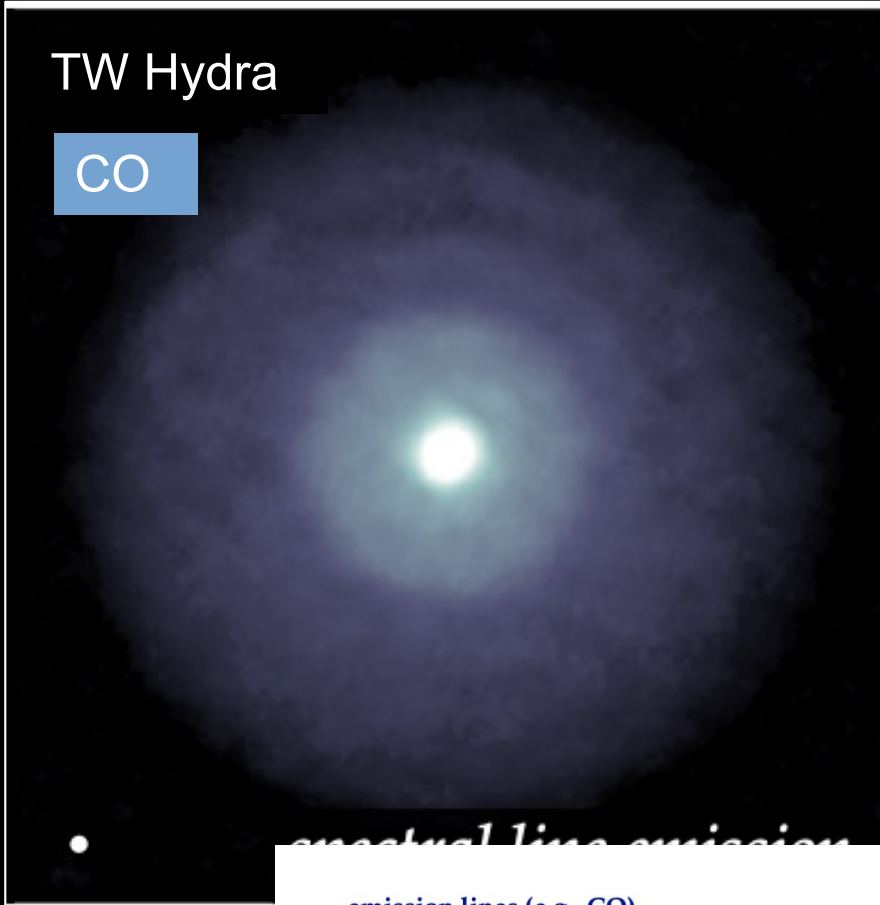
Van Boekel+2018



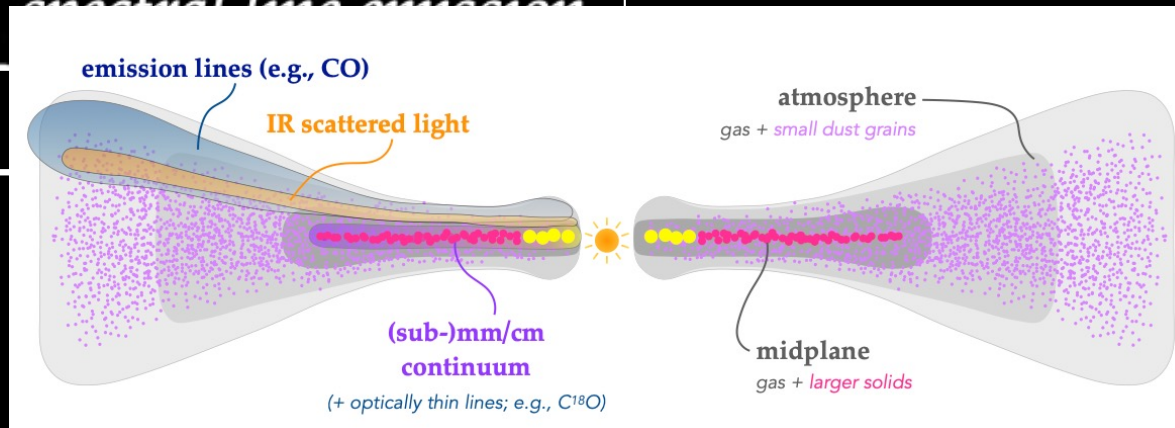
Disks are observed in molecular emission (radio observations)

TW Hydra

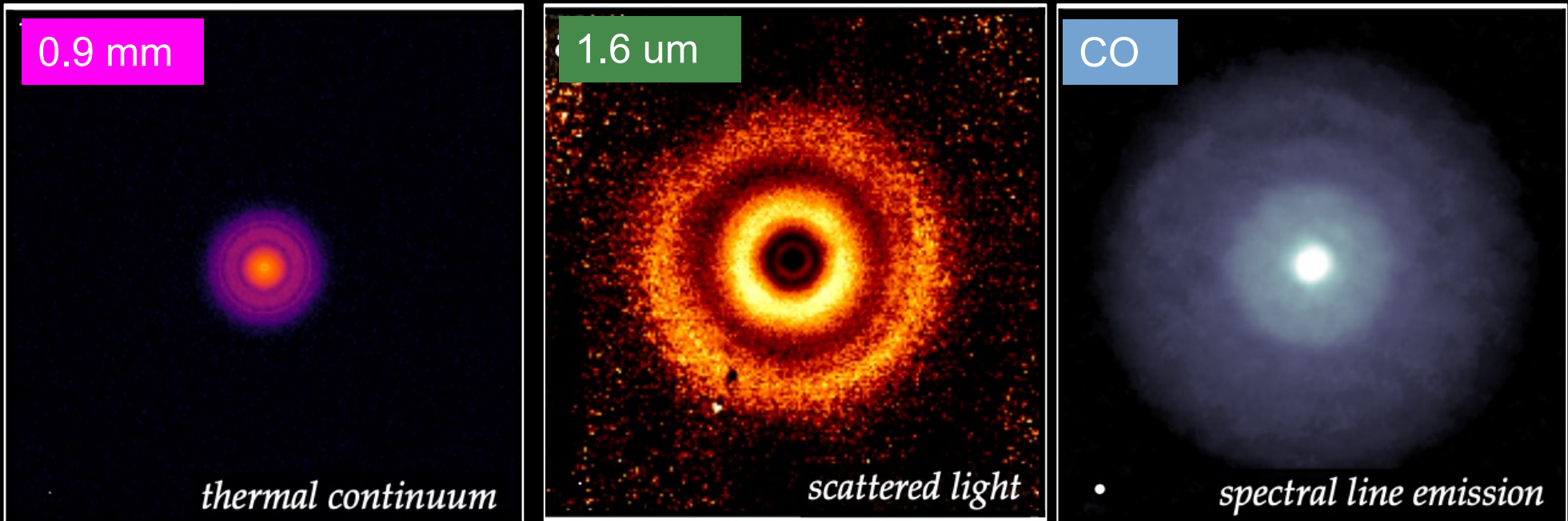
CO



Huang+2018



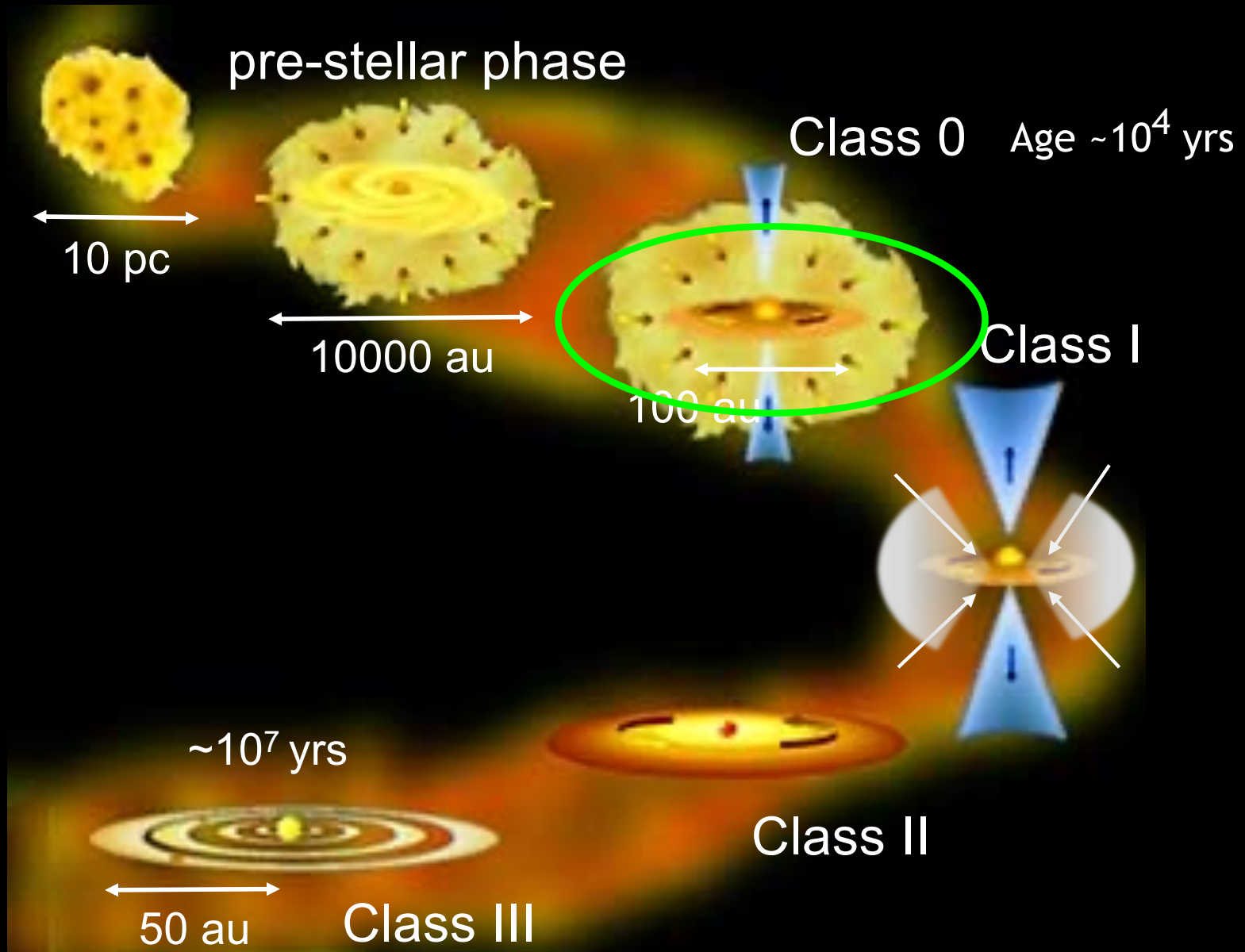
Protoplanetary disks (IR and radio observations)



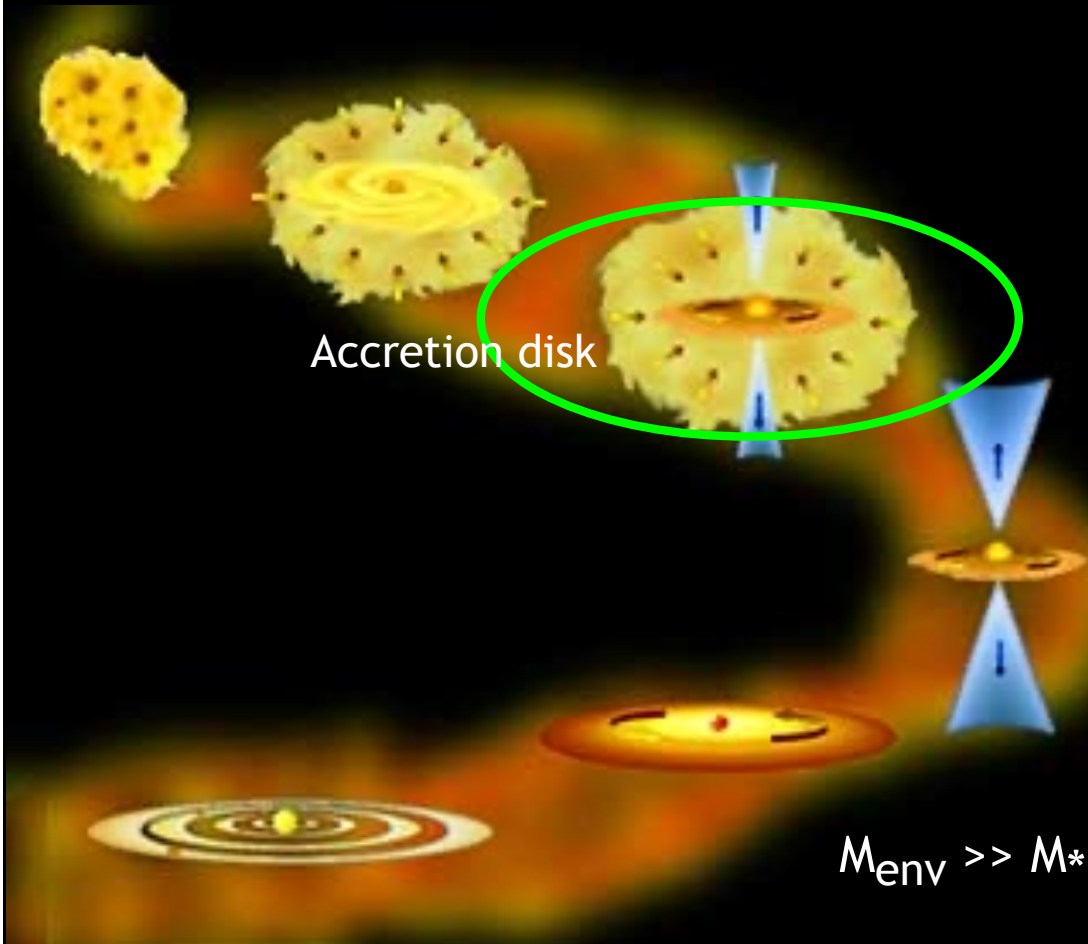
The disk is much smaller, traced by millimeter emission than when it is detected by scattered light and molecular emission.

Radiation pressure pushes small grains which are coupled to the gas outwards, in this process small grains drag large grains toward the center. (radial drift, Widenschilling 1977).

Star formation paradigm (Shu, Adams, Lizano 1987)



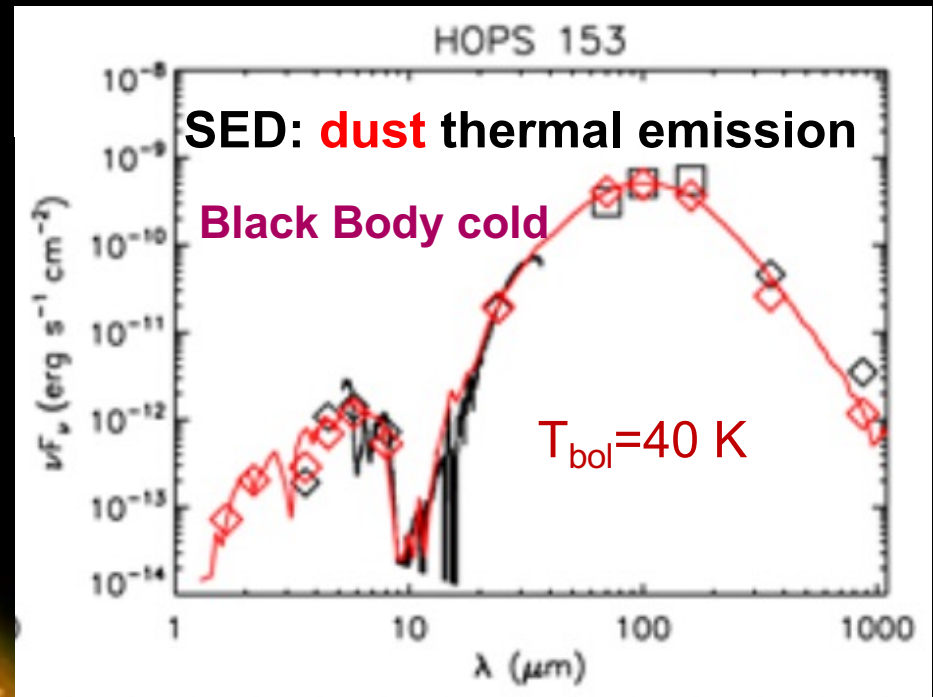
Class 0 (dominated by the envelope)



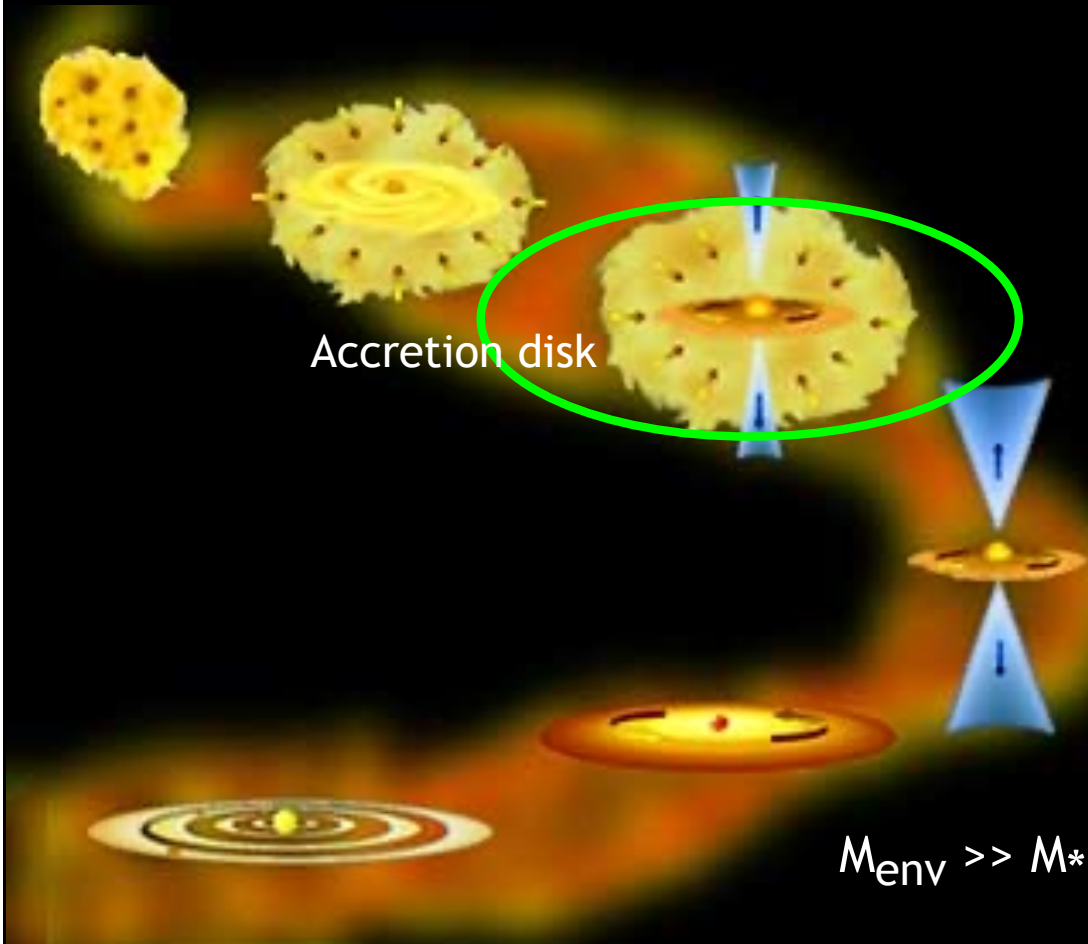
$$M_{\text{env}} \gg M_*$$

$$T_{\text{bol}} < 70 \text{ K}$$

$$\text{Age} \sim 10^4 \text{ yrs}$$



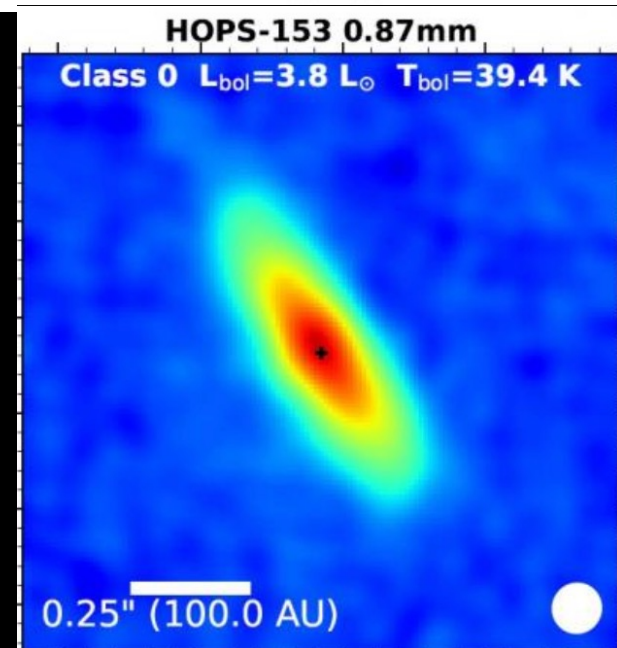
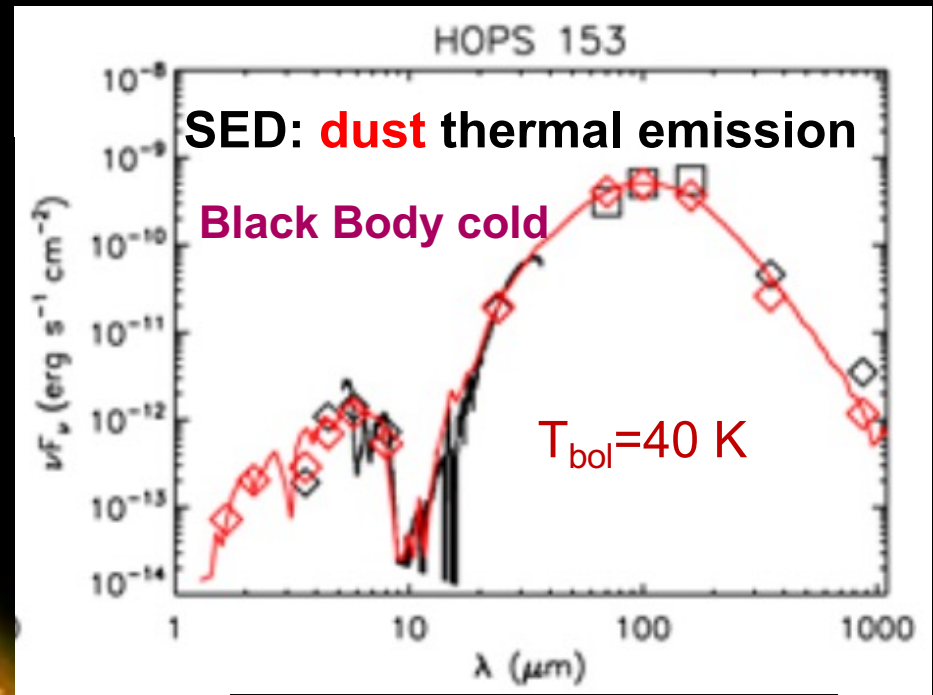
Class 0 (dominated by the envelope)



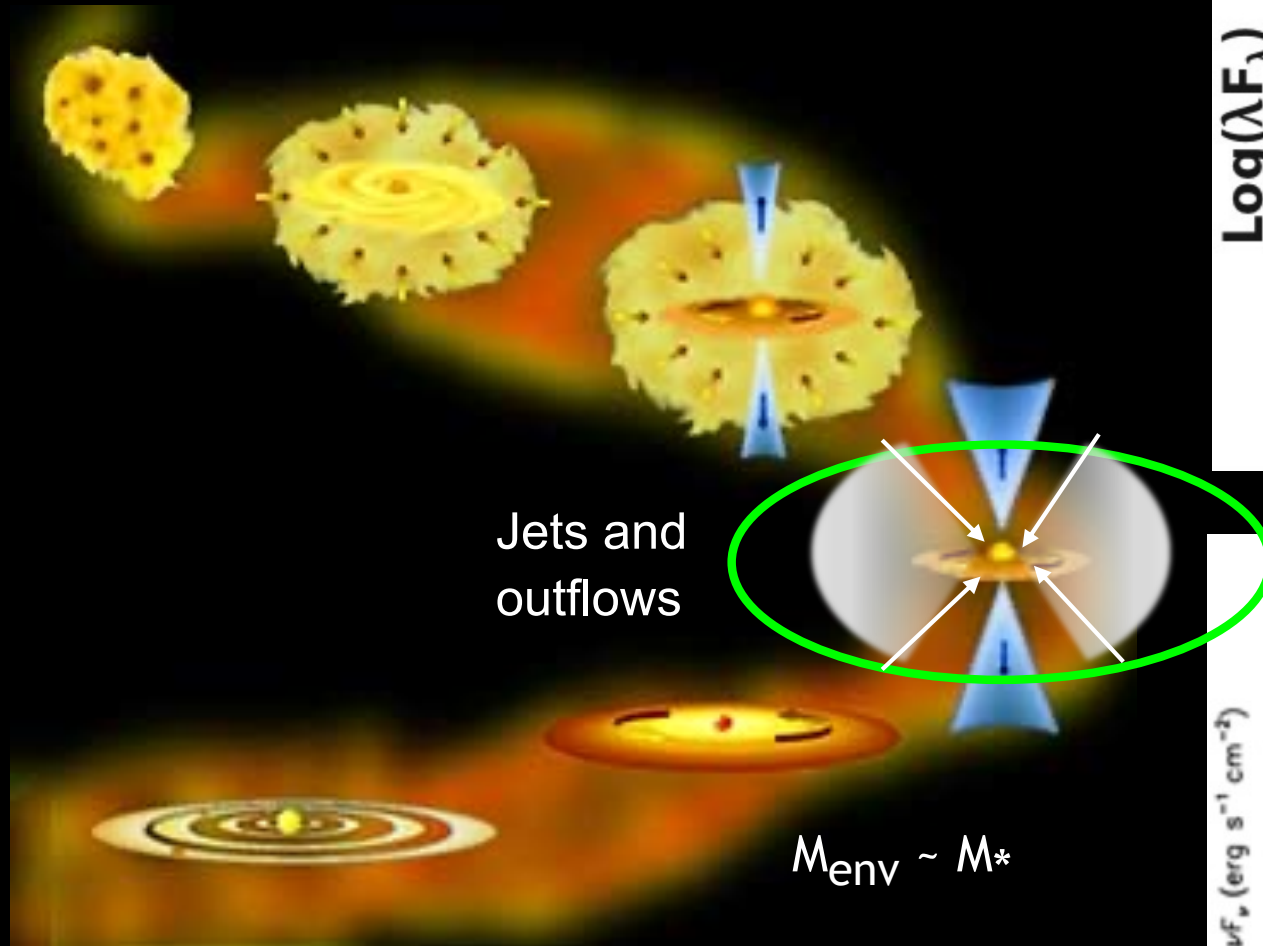
$$M_{\text{env}} \gg M_*$$

$$T_{\text{bol}} < 70 \text{ K}$$

$$\text{Age} \sim 10^4 \text{ yrs}$$



Class I (disk + envelope)

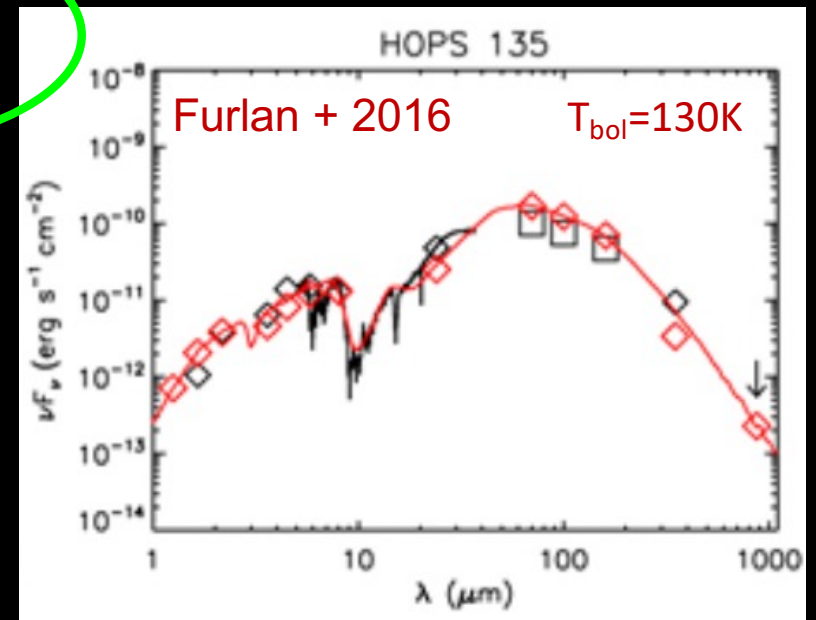
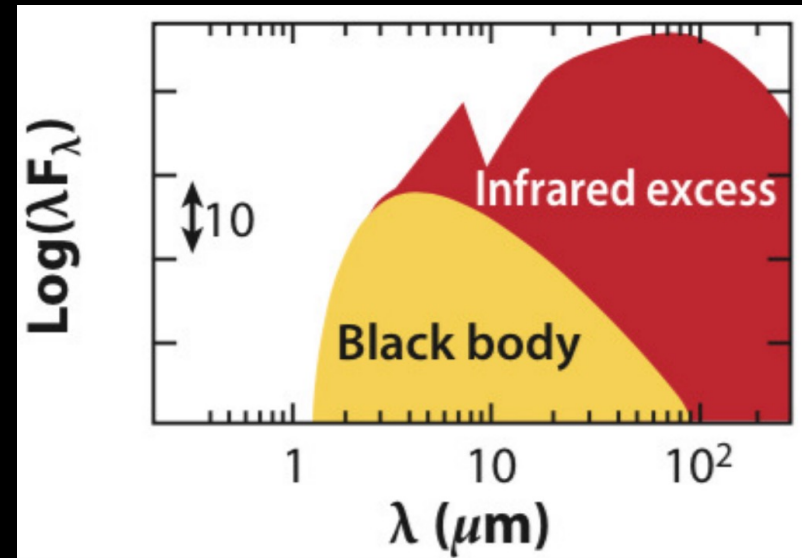


Jets and outflows

$$M_{\text{env}} \sim M_*$$

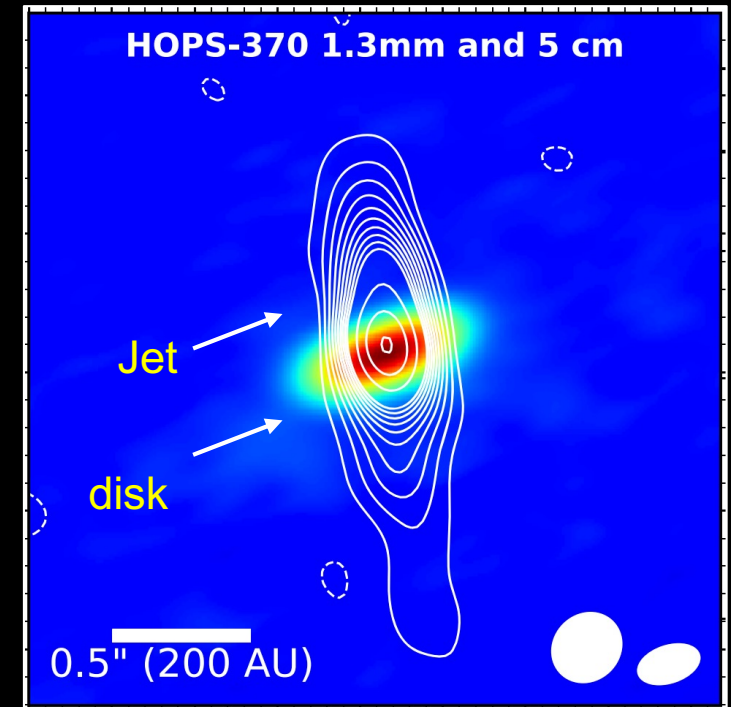
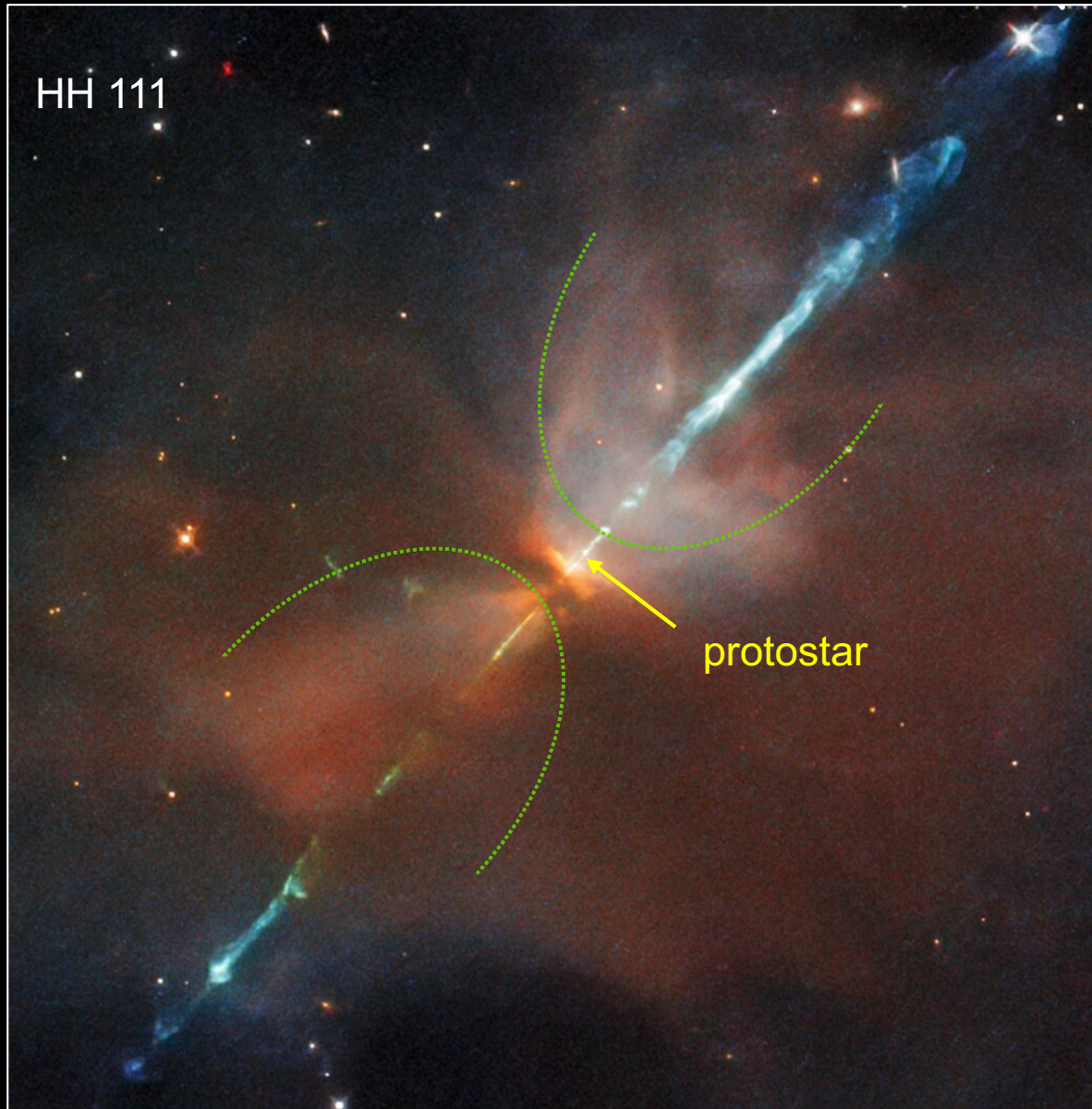
$$T_{\text{bol}} \sim 100\text{-}700 \text{ K}$$

$$\text{Age} \sim 10^5 \text{ yrs}$$



Class I (disk + envelope)

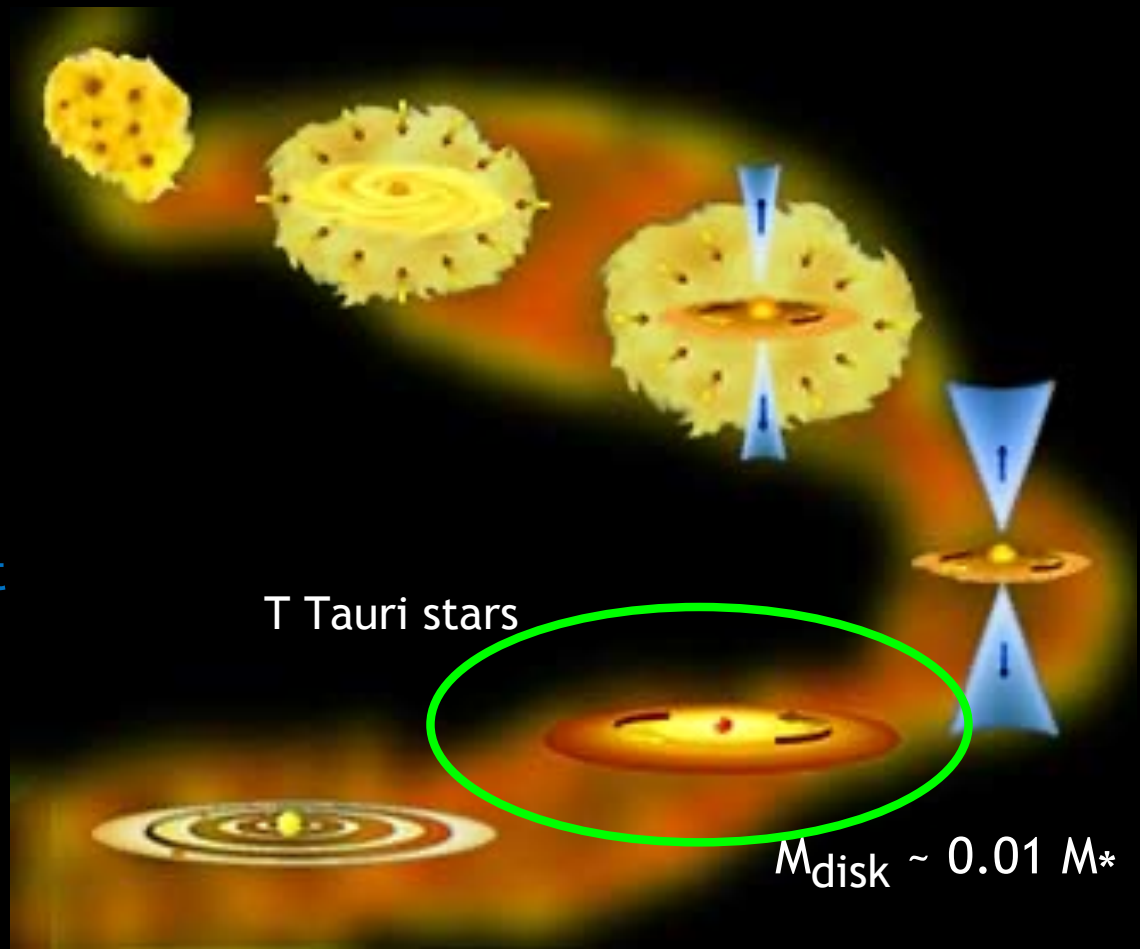
The disk plays a crucial role not only in feeding material the protostar, but also in collimating these spectacular optical jets.



Tobin+2020

<https://spfe.es/en/catalogs/>

Class II (dominated by the disk)



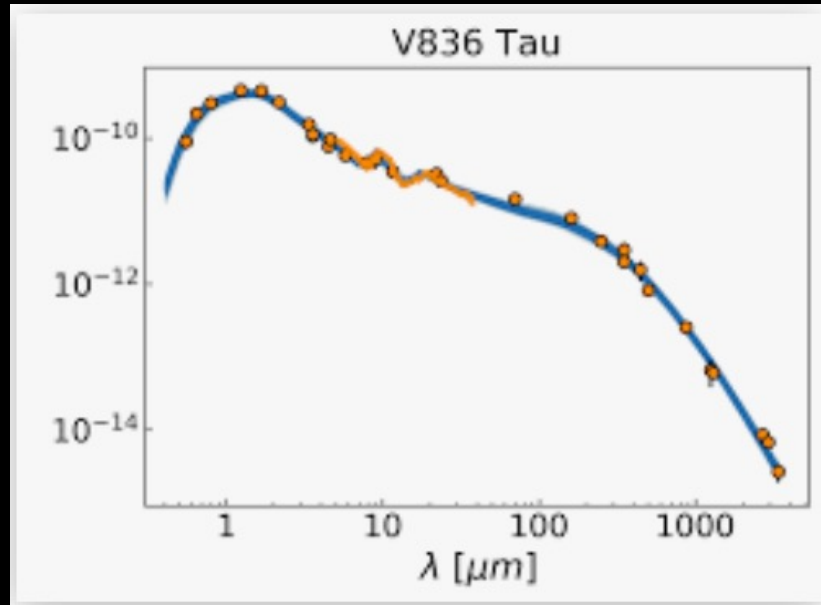
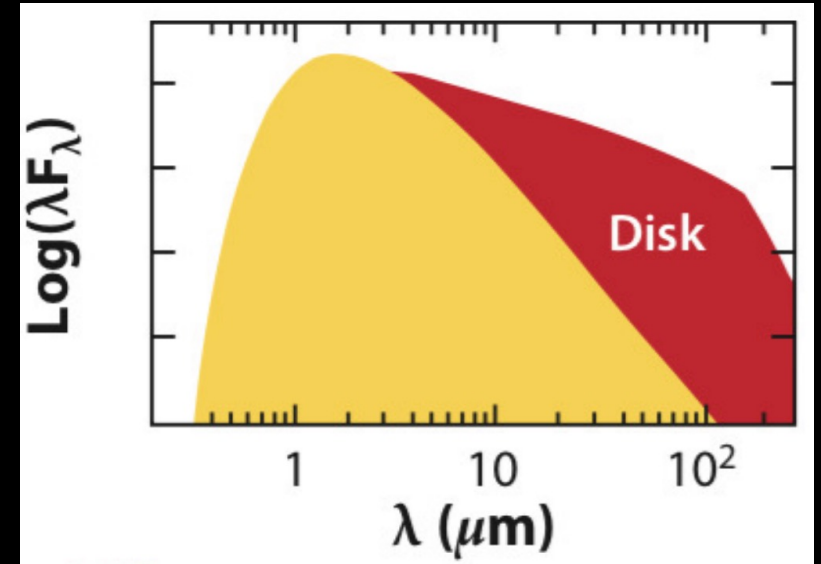
T Tauri stars

$$M_{\text{disk}} \sim 0.01 M_*$$

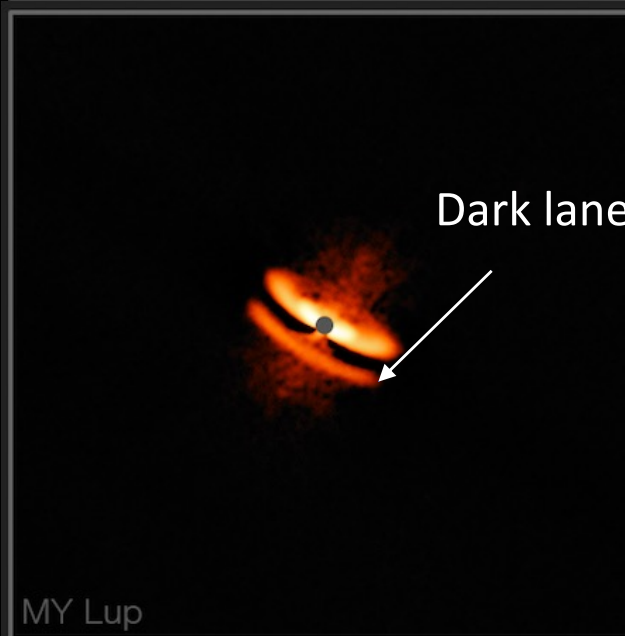
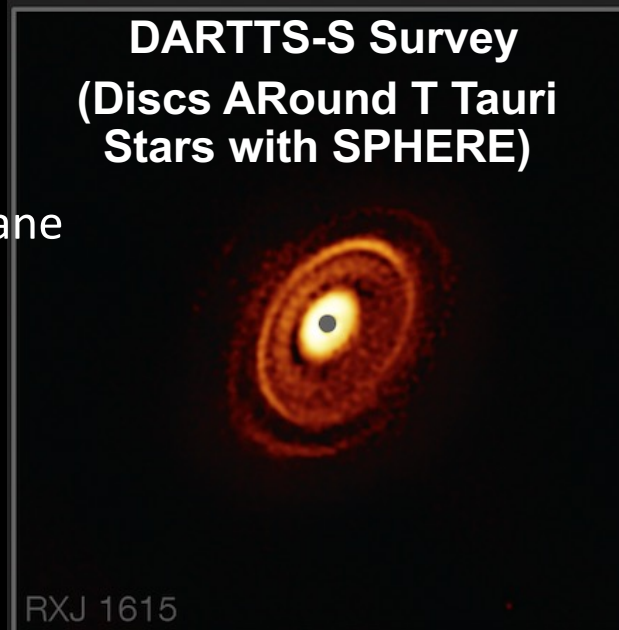
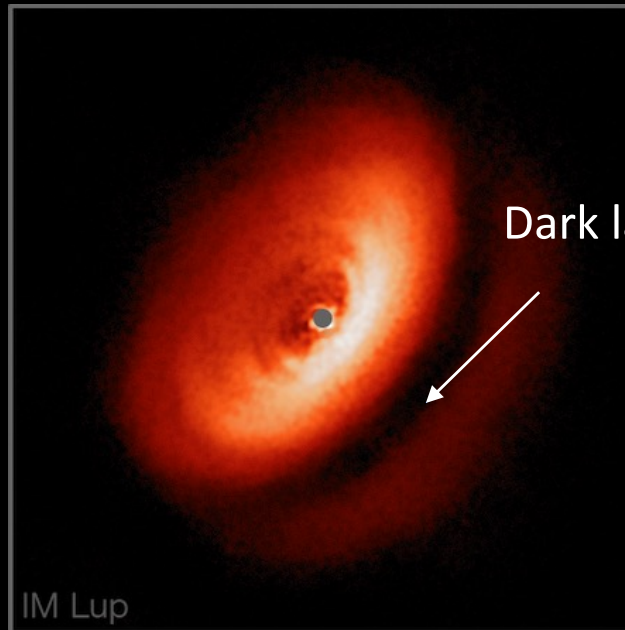
$$T_{\text{bol}} \sim 700\text{-}3000 \text{ K}$$

$$\text{Age} \sim 10^6 \text{ yrs}$$

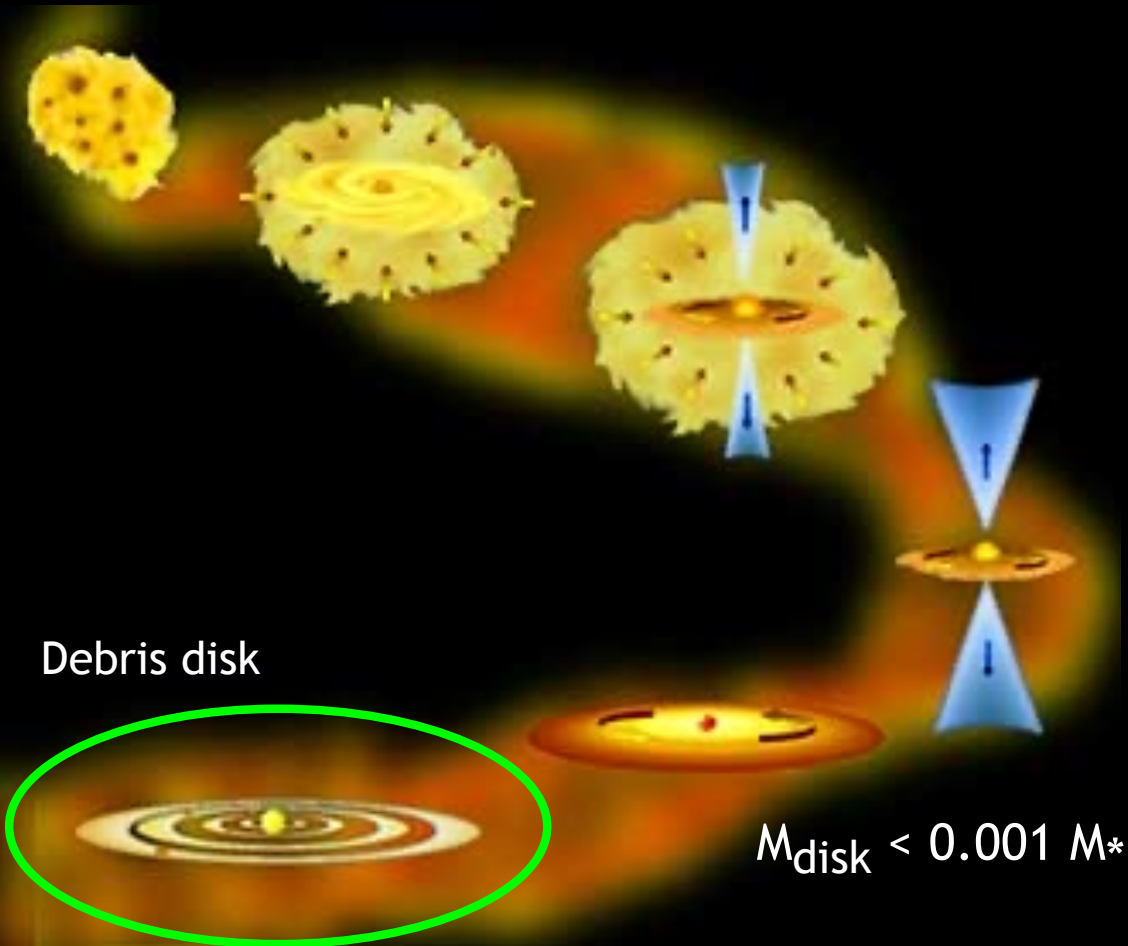
Since the circumstellar disk becomes a planetary system, it is called protoplanetary disk in this stage.



Class II (observations)



Class III (star + debris disk)

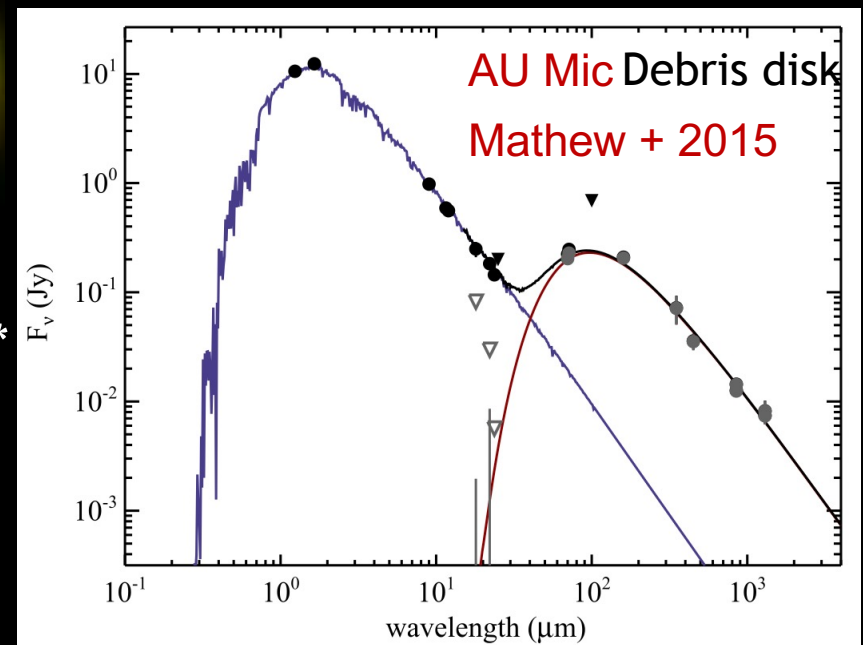
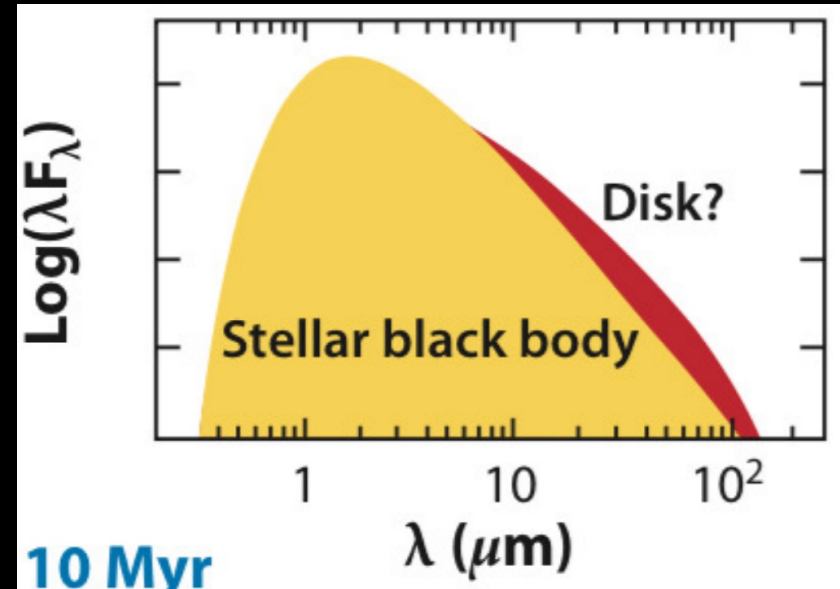


Debris disk

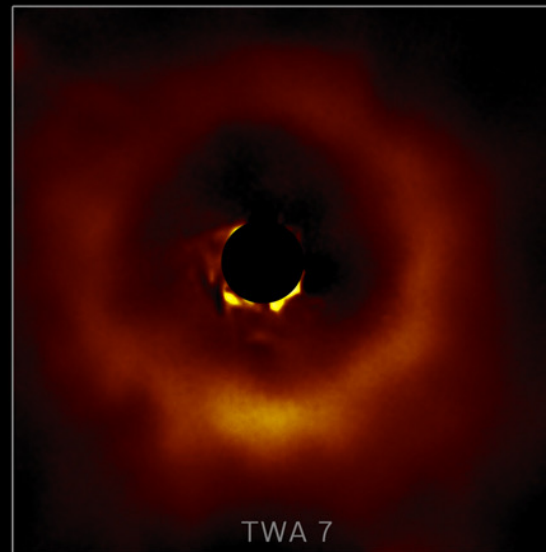
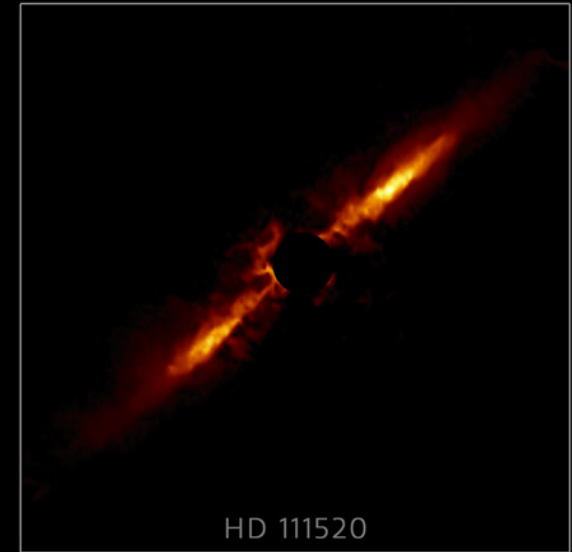
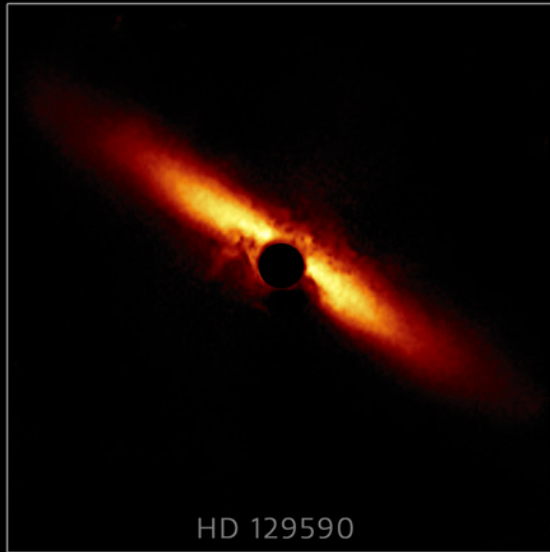
$$M_{\text{disk}} < 0.001 M_{\odot}$$

$$T_{\text{bol}} > 3000 \text{ K}$$

$$\text{Age} \sim 10^7 \text{ yrs}$$



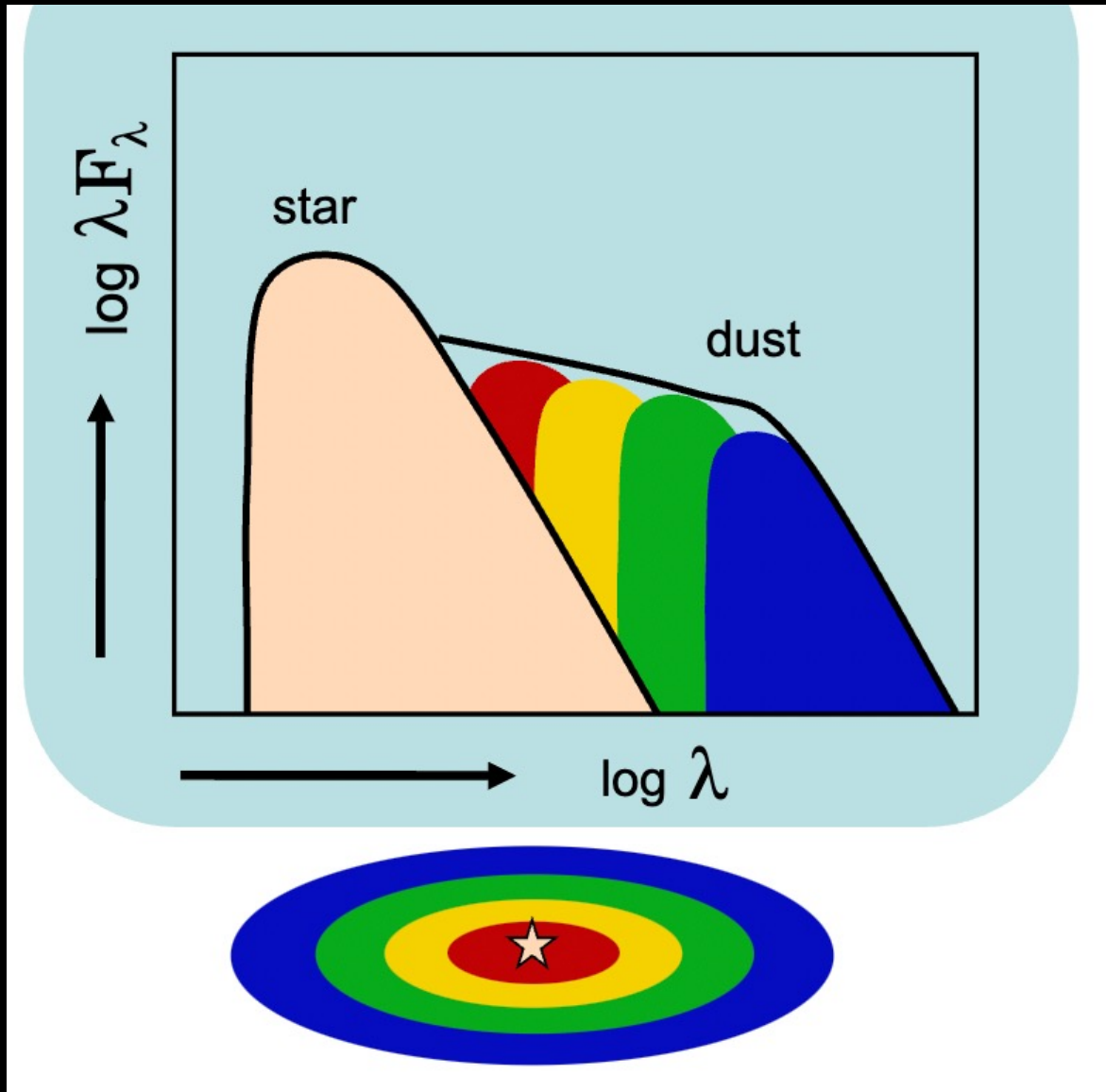
Class III (observations)



GPI (Gemini planet image)

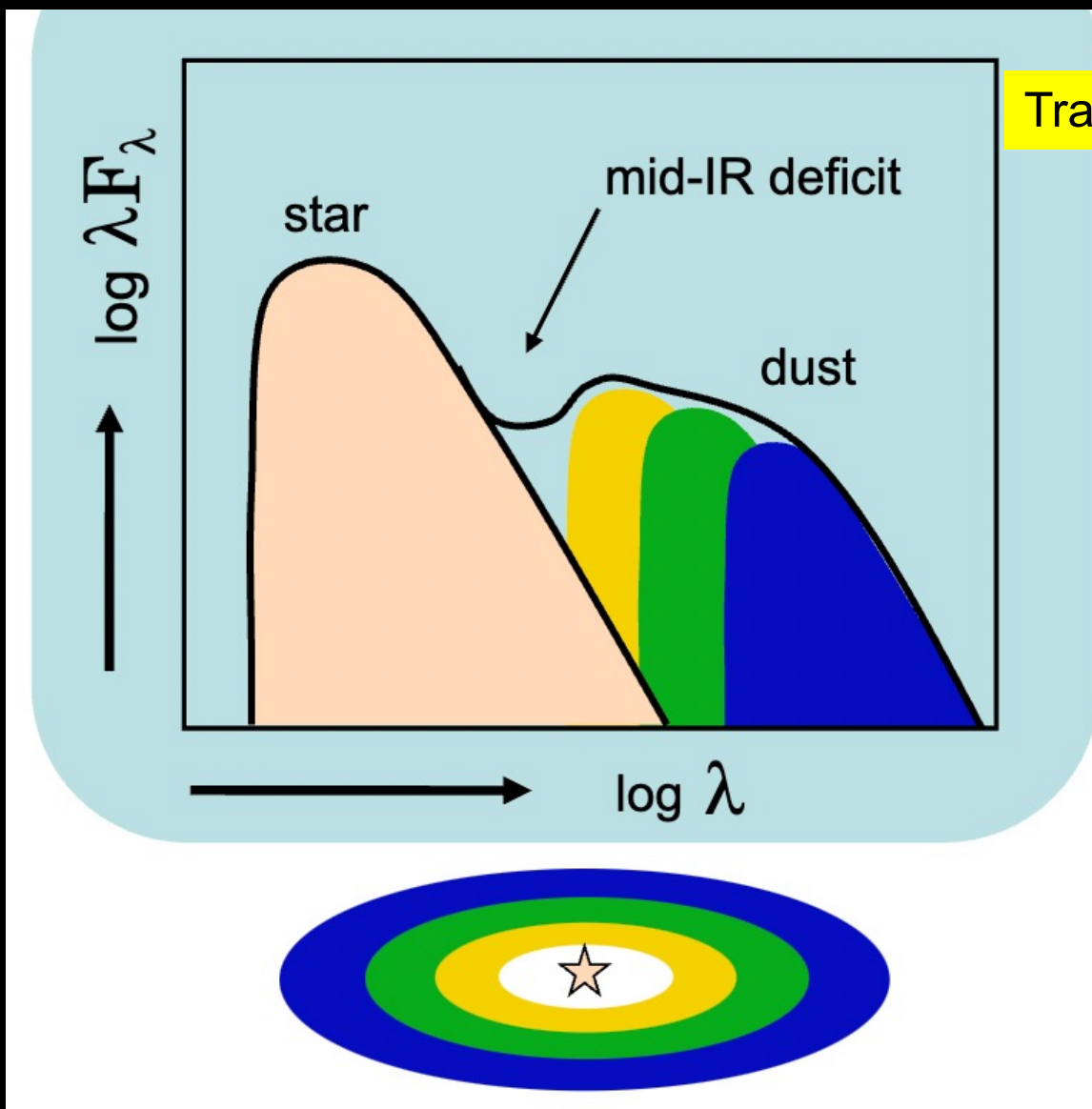
Flat spectrum (class II)

Full disk



The disk SED shows a flattened shape that can be explained for a superposition of several disk layers emitting as a BB at different temperatures.

Flat spectrum

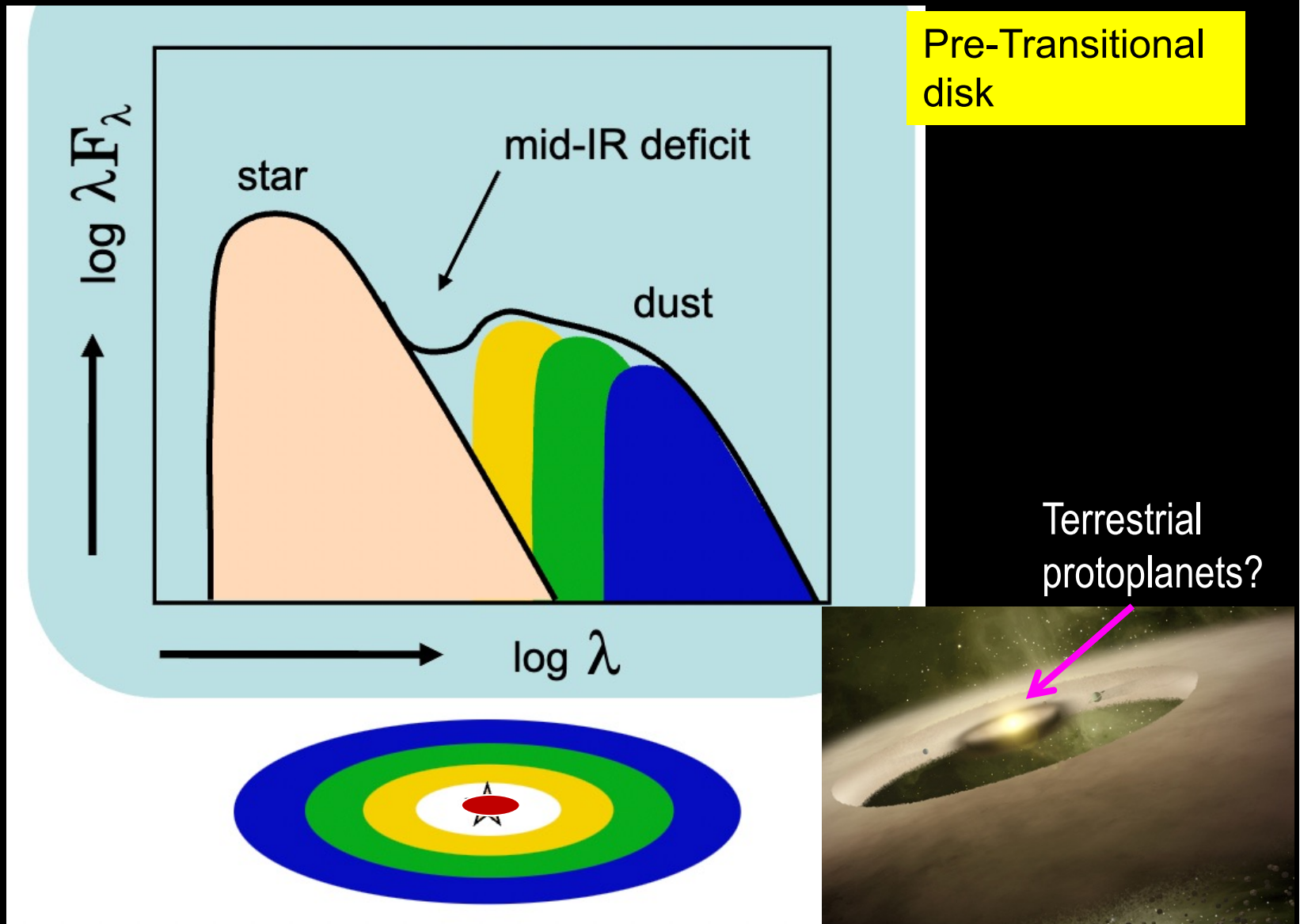


Transitional disk

Based on the modelling of the spectra observed by Spitzer telescope, Calvet + identified many of these protostellar disks. ALMA has confirmed their central cavities



Flat spectrum

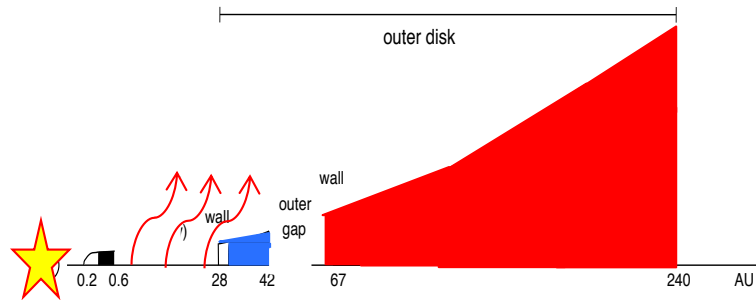


A subset of these disks shows still a significant amount of dust inside the cavity

Origin of cavities in Transitional disks

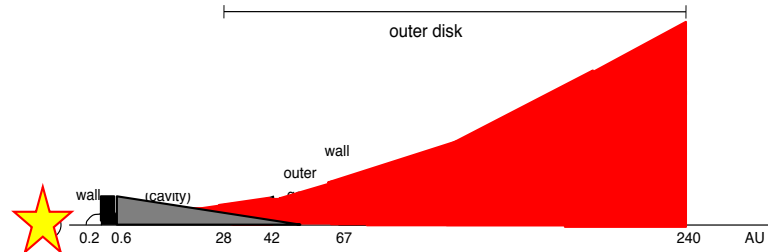
Photoevaporation

EUV or X rays from the Star

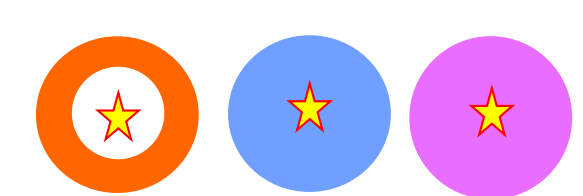
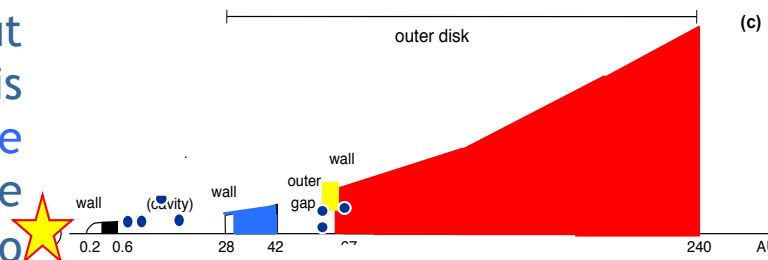


Geometric shadowing.

Shadow casts by a wall or residual disk

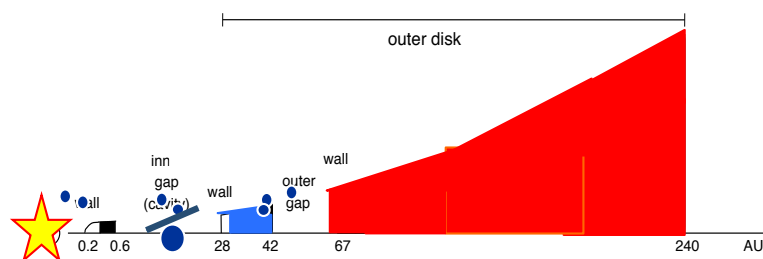


Grain growth Some cavities are seen in IR scattered light but not in mm wavelengths. This suggests that they are not true cavities, but regions where dust grains may have grown to a size whose IR emission is undetected.



Companions

The presence of another body that dynamically creates a gap



Adapted from Van der Marel

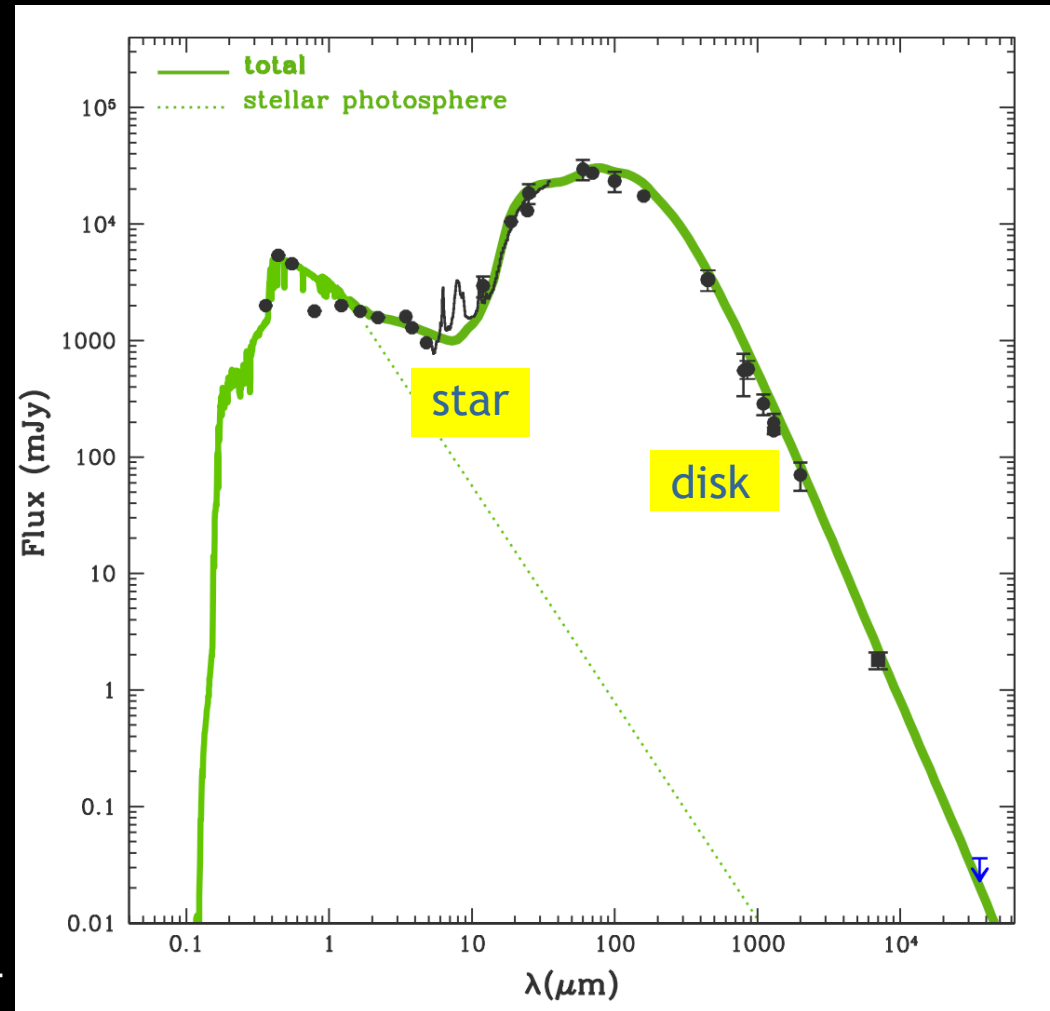
The protoplanetary disk of HD169142

Relatively isolated intermediate mass protostar with an age $\sim 6\text{Myr}$ (D=140pc Gaia)

This source has been observed from ultraviolet to cm wavelengths.

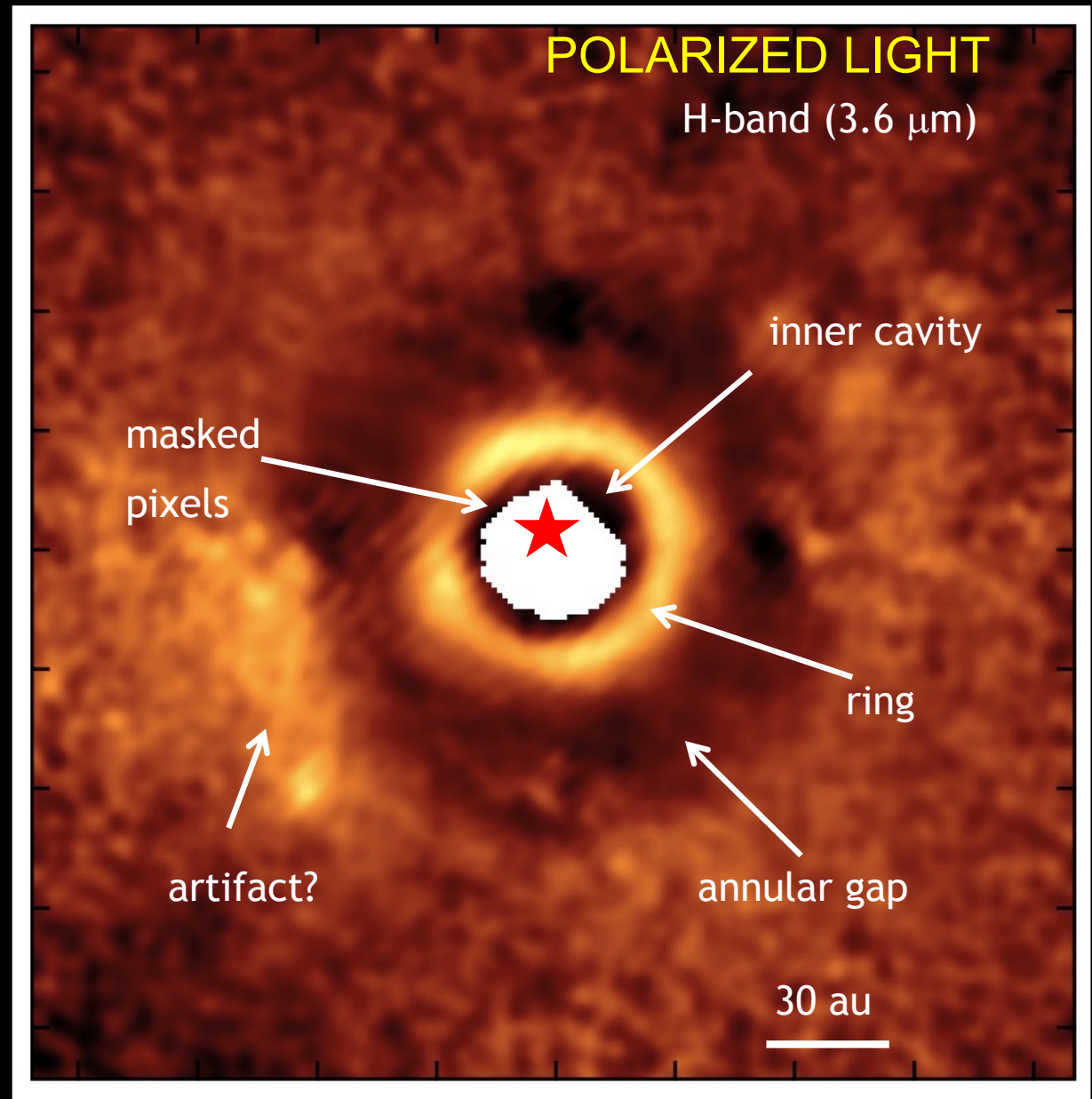
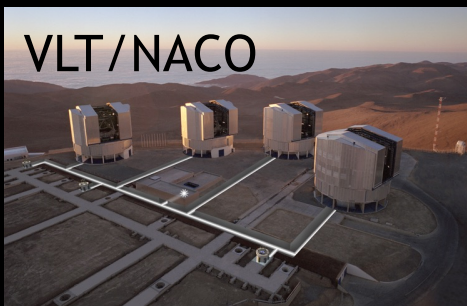
Based on the SED modelling (deficit in the near IR) Grady+(2007), Meeus+(2010), Honda+(2012) predicted the existence of a central cavity in this protoplanetary disk=> **Transitional disk!**

Dent + 2006, Osorio + 2014



Near-IR polarized light imaging of HD169142

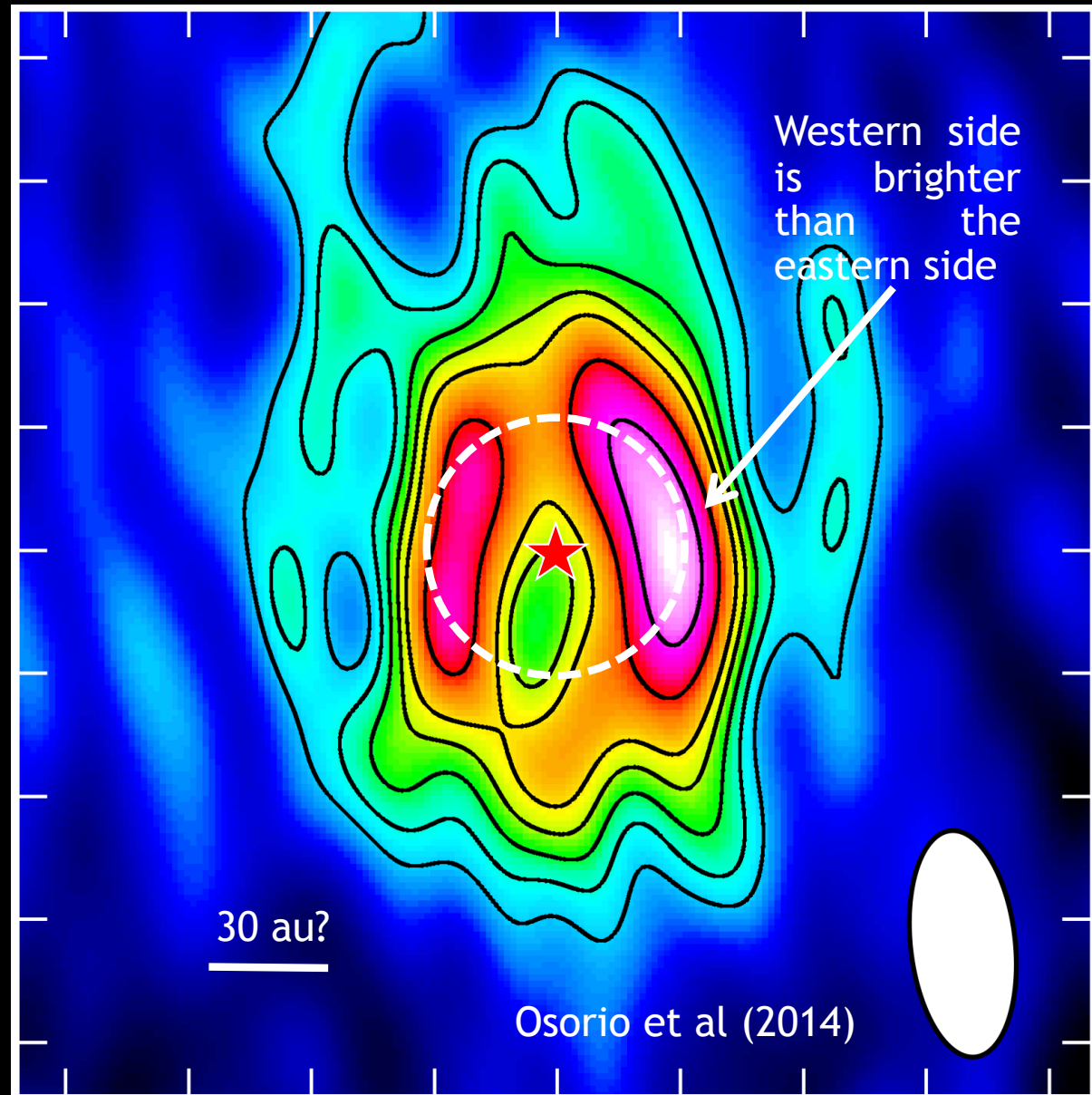
- Face-on disk, $R=200$ au
- Central cavity
- Ring at 30 au (cavity's wall)
- Gap spanning from 40 to 70 au, where the surface brightness decreases by a factor of 3.



(Quanz et al 2013)

VLA 7 mm imaging of HD169142

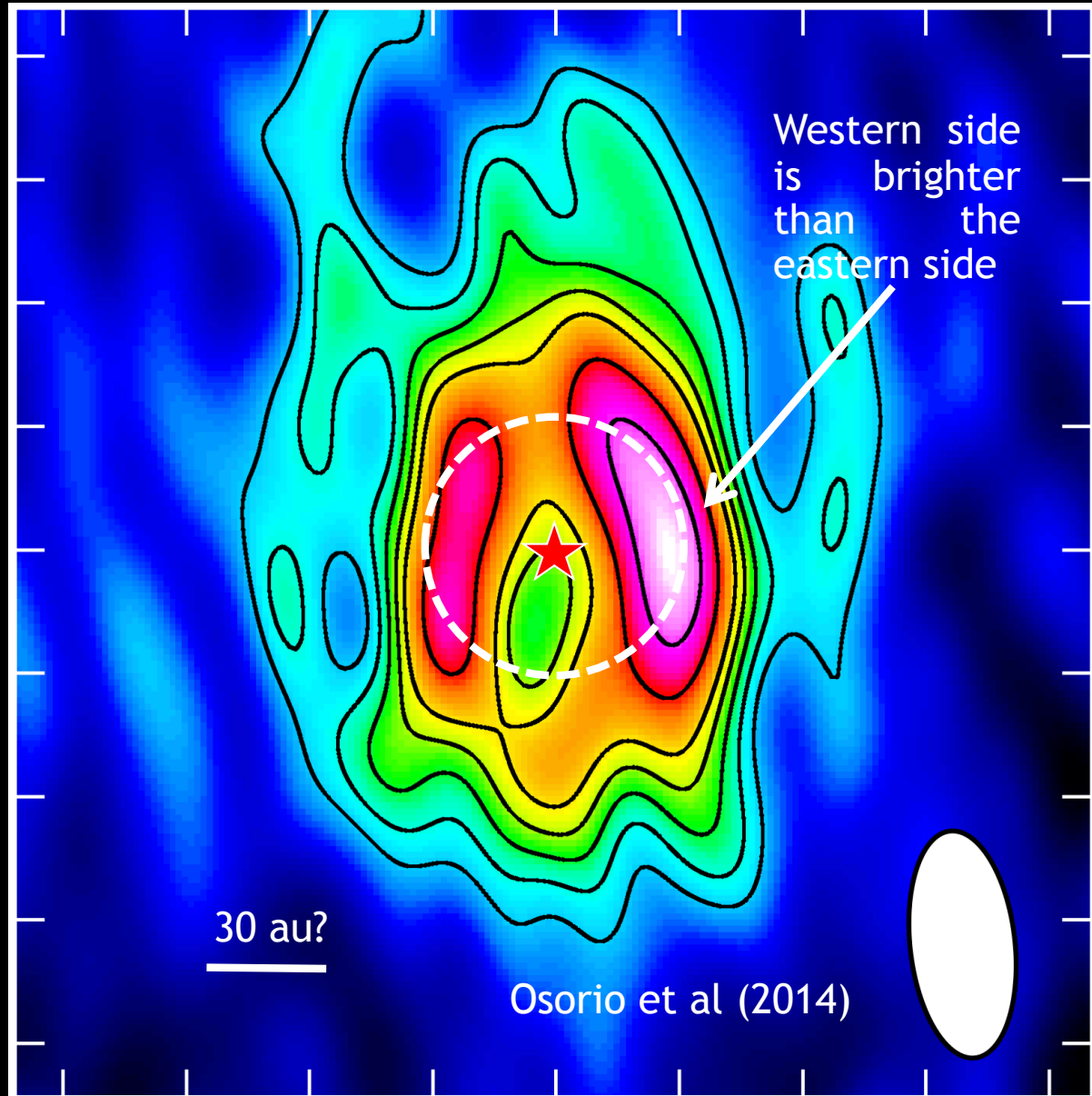
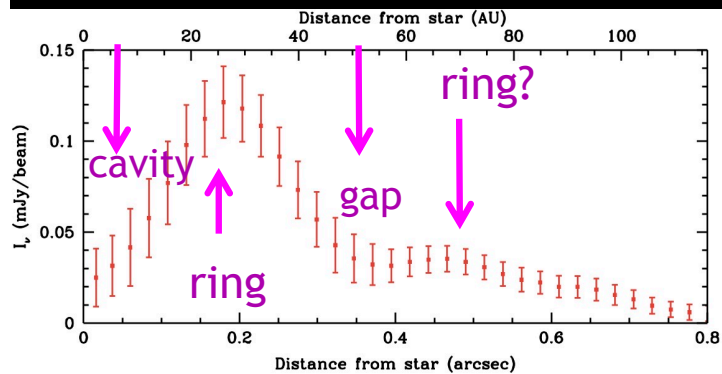
The **cavity** and **ring** were confirmed (may be a second ring) by **dust thermal emission**.



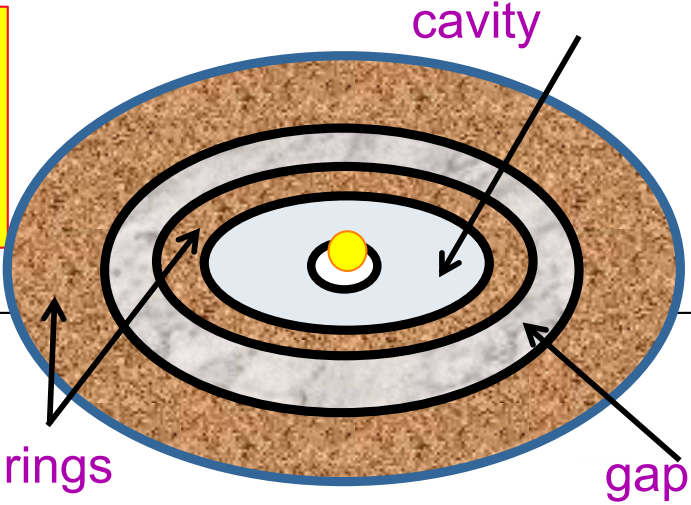
VLA 7 mm imaging of HD169142

The **cavity** and **ring** were confirmed by **dust thermal emission**.

7mm azimuthally-averaged radial profile



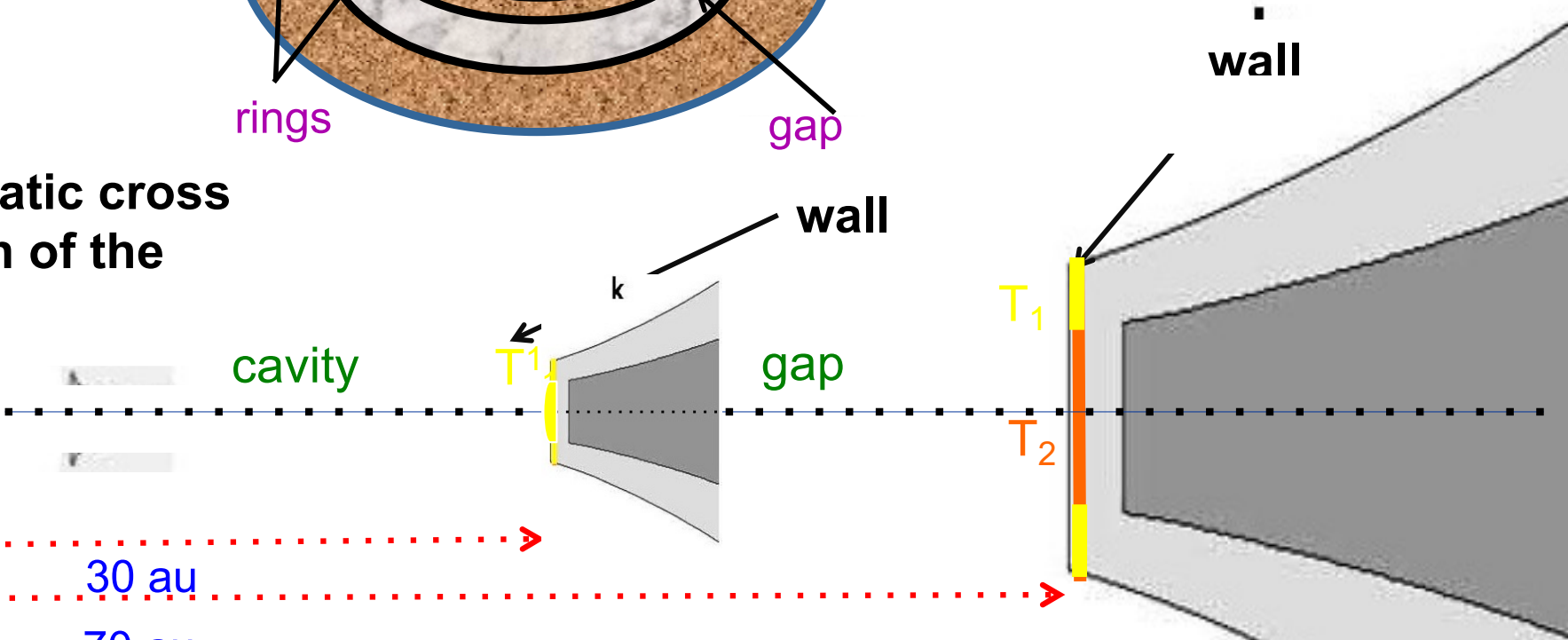
Double-ringed disk



schematic cross section of the disk



star



← →

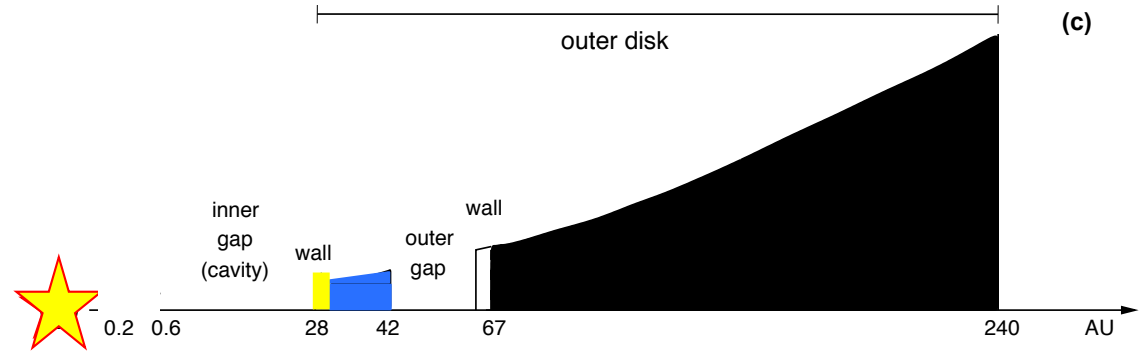
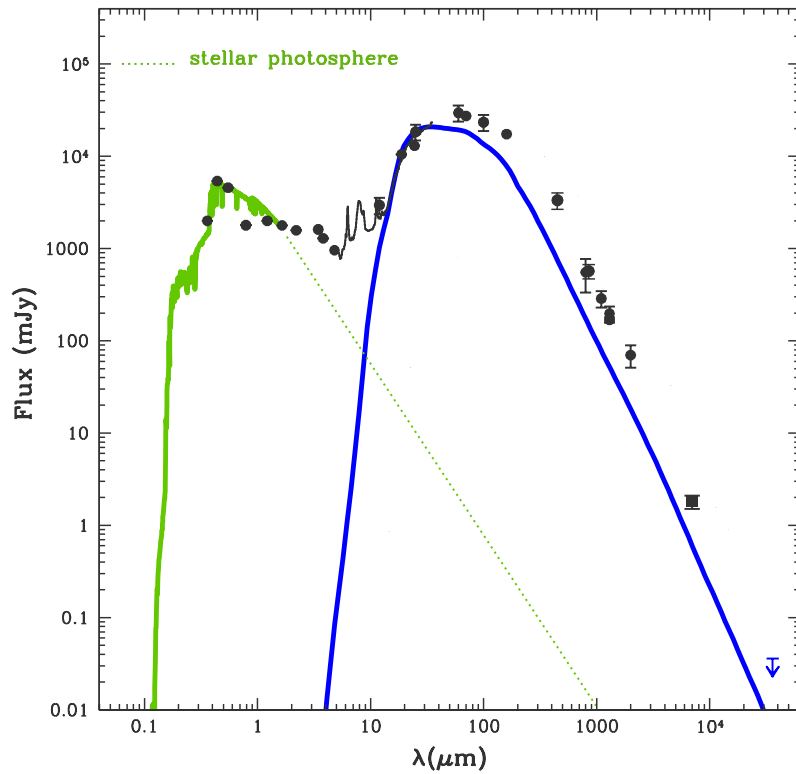
30 au

← →

70 au

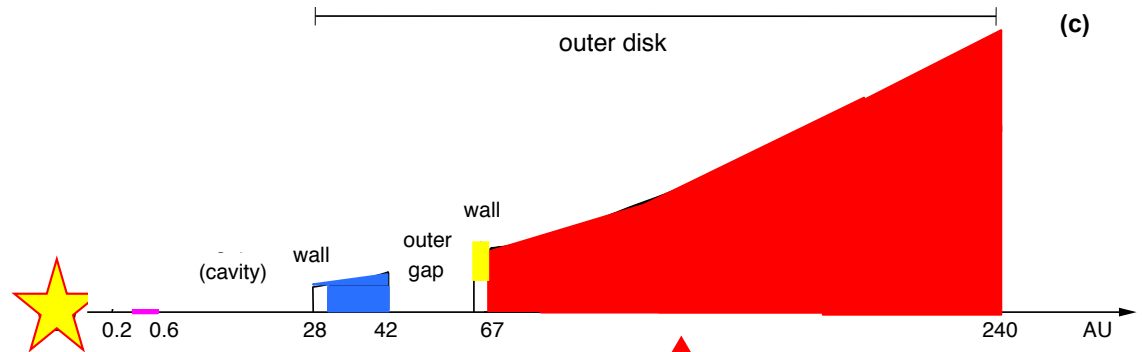
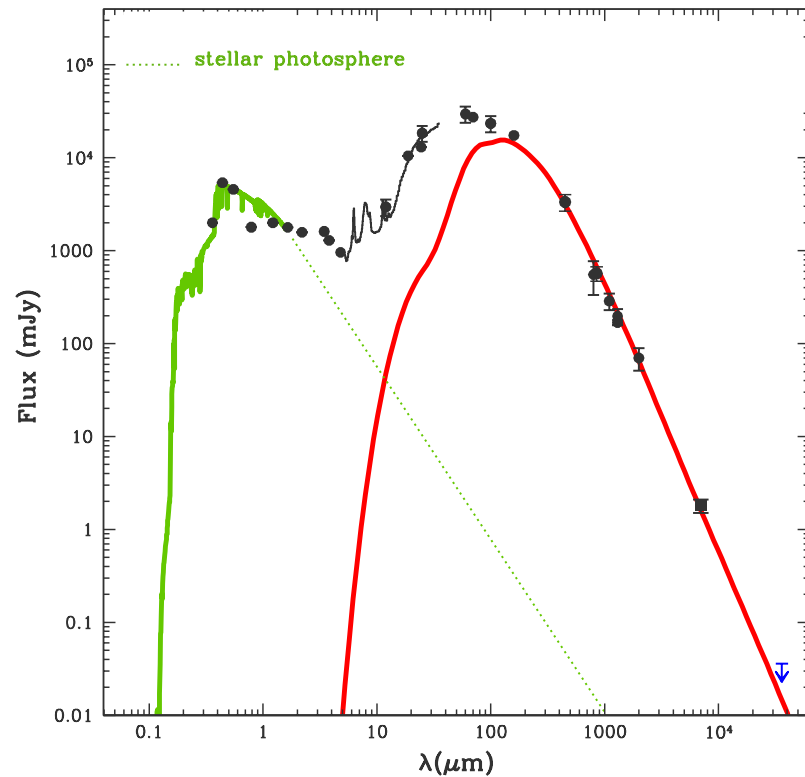
Walls of the gaps:
 T_1 : illuminated zone
 $T_2 < T_1$: shadowed zone
 large accumulation of big grains in the outer edges of the walls of each gap (relevant to explain 7mm image)

Contribution to the SED of the disk components



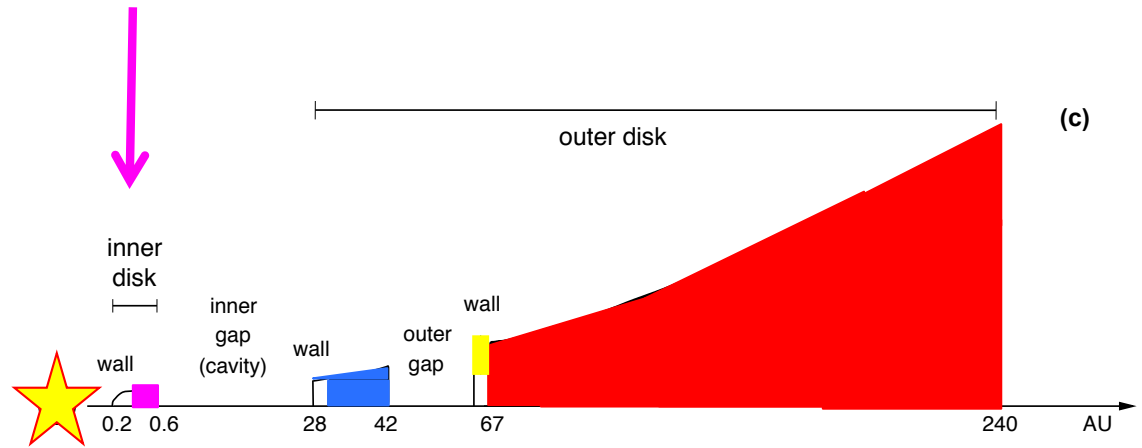
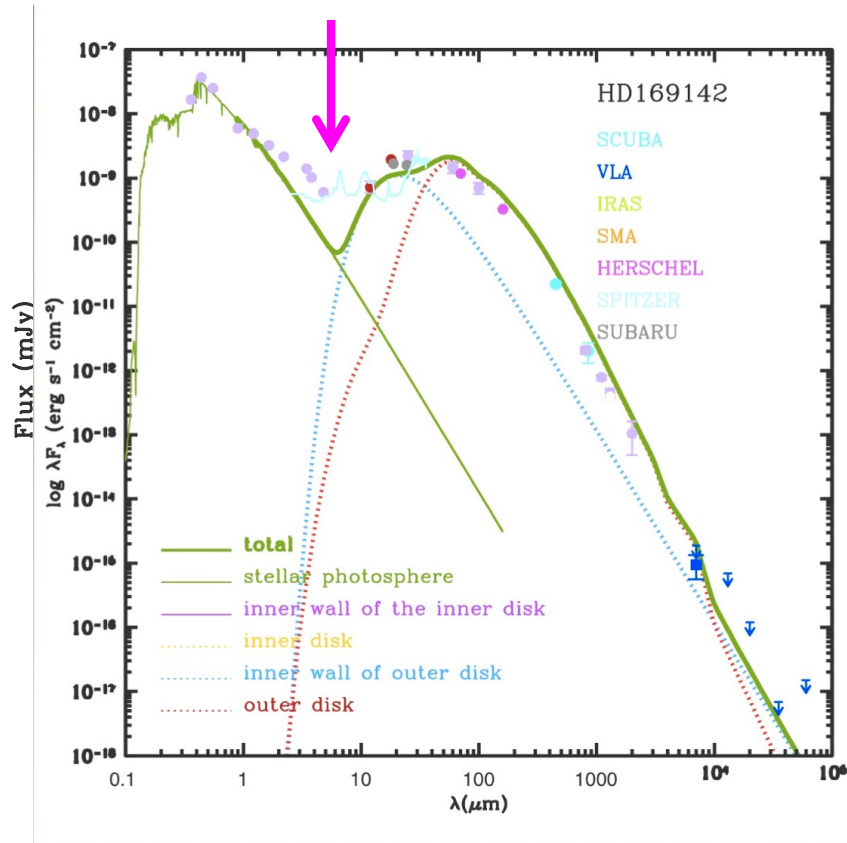
Disk + Wall of the inner Gap
 $H=9\text{au}$, $T_1=110\text{K}$,

Contribution to the SED of the disk components

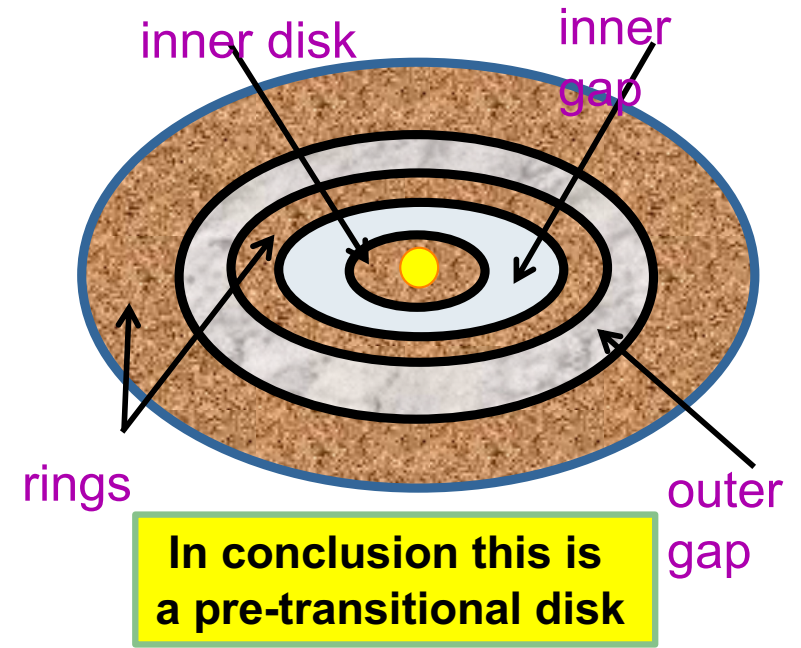
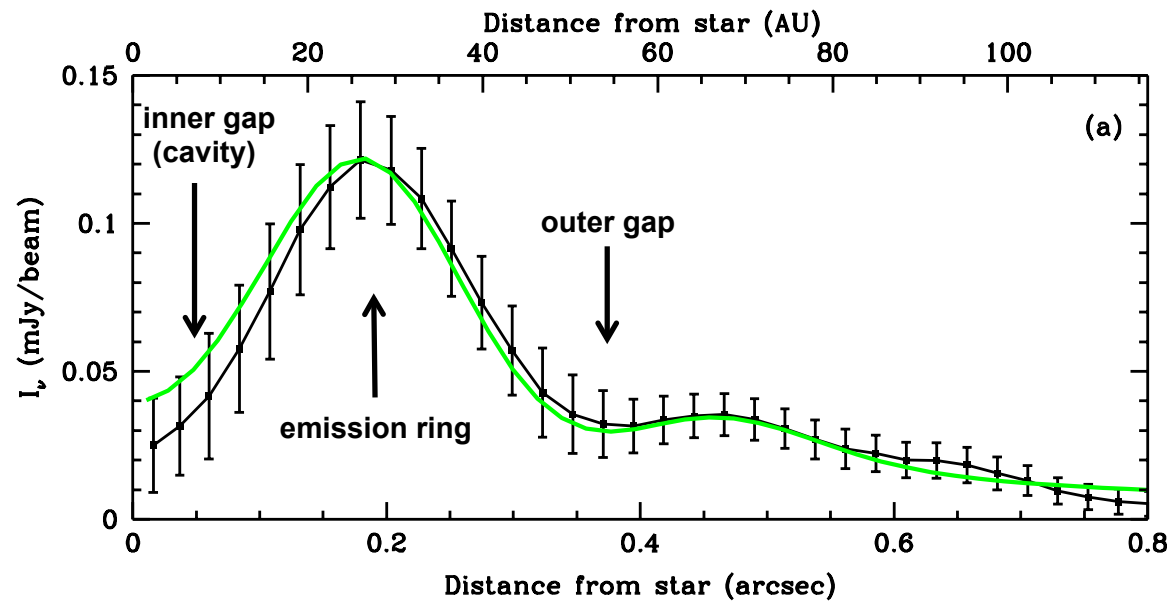
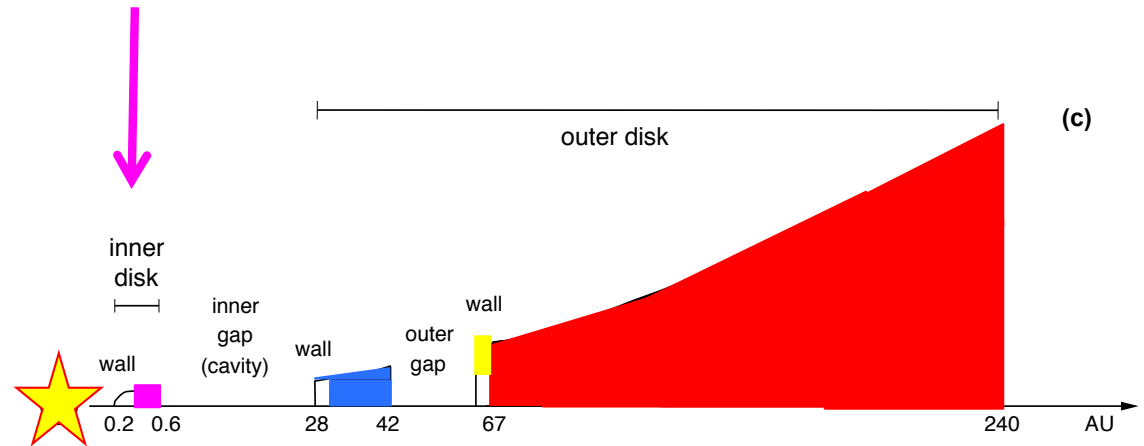
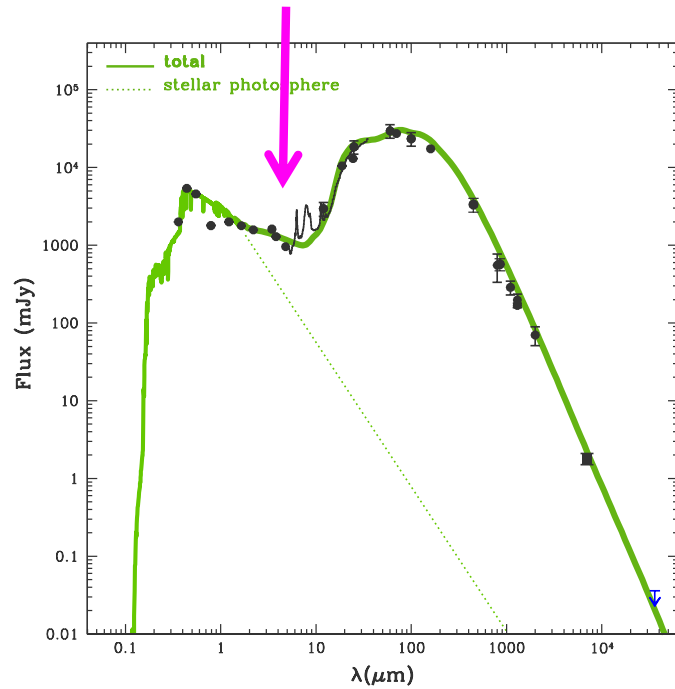


- Outer disk + Wall of the Outer gap
($H=30$ au,
upper layers: $T_1=70\text{K}$,
lower layers: $T_2=40\text{K}$)
- Illuminated fraction: 10%

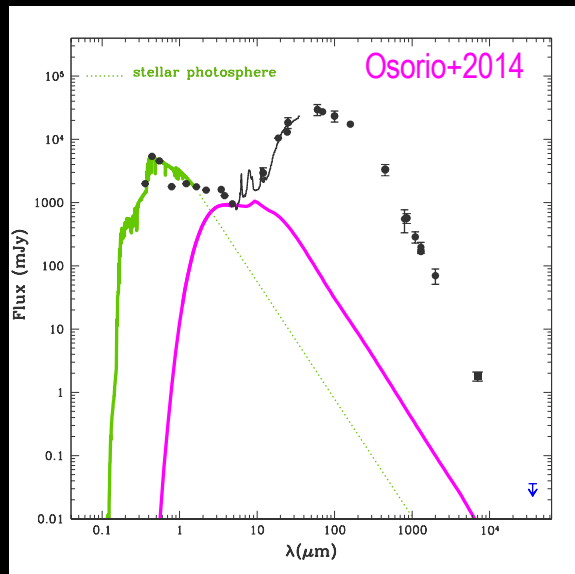
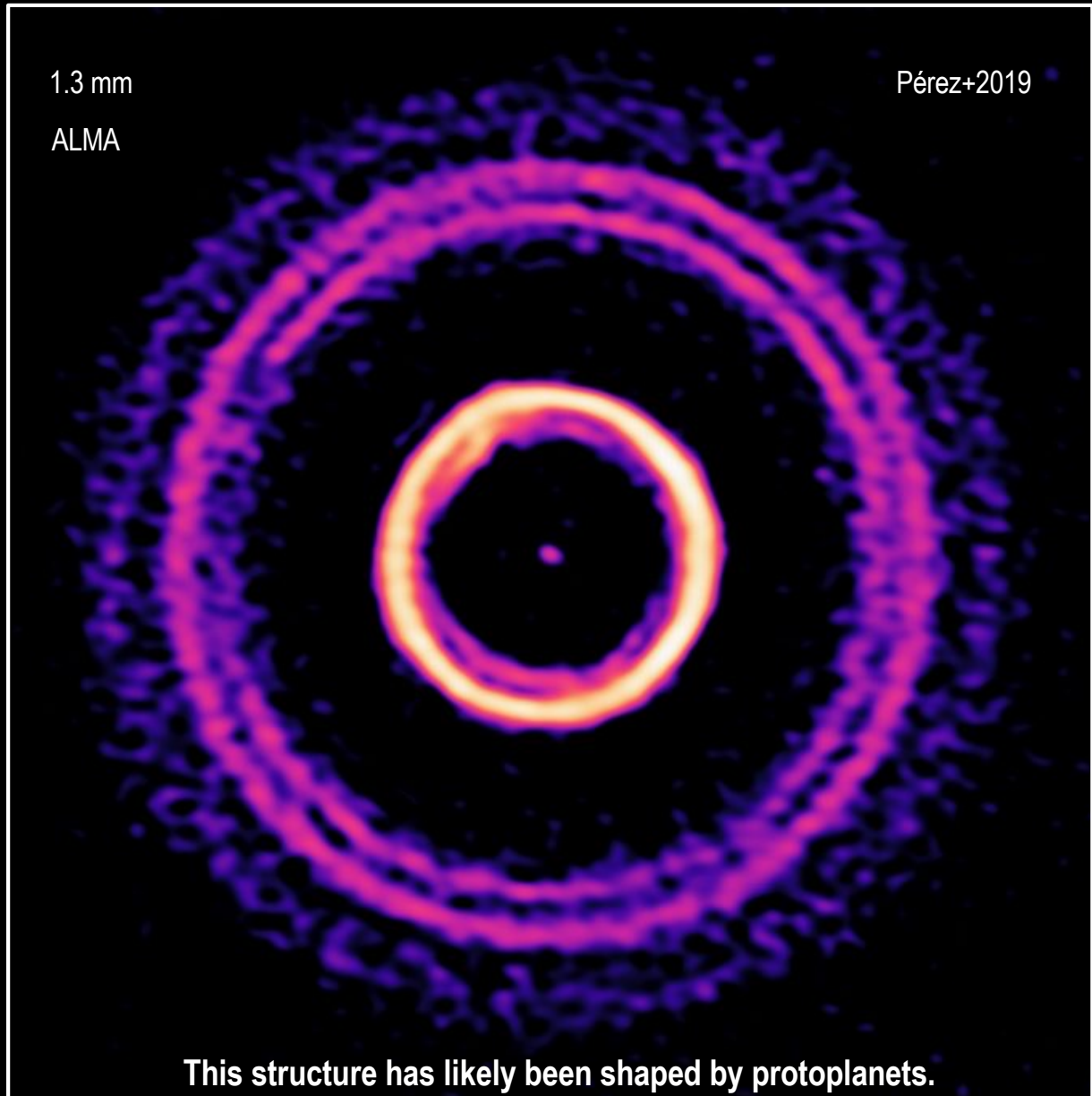
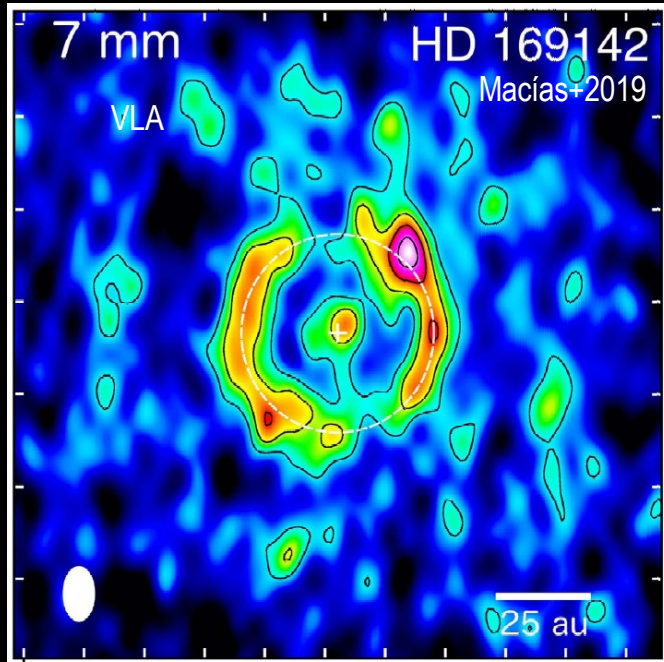
Contribution to the SED of the disk components



Contribution to the SED of the disk components



ALMA and VLA observations of HD169142 disk



Signs of planet formation in the HD169142 disk

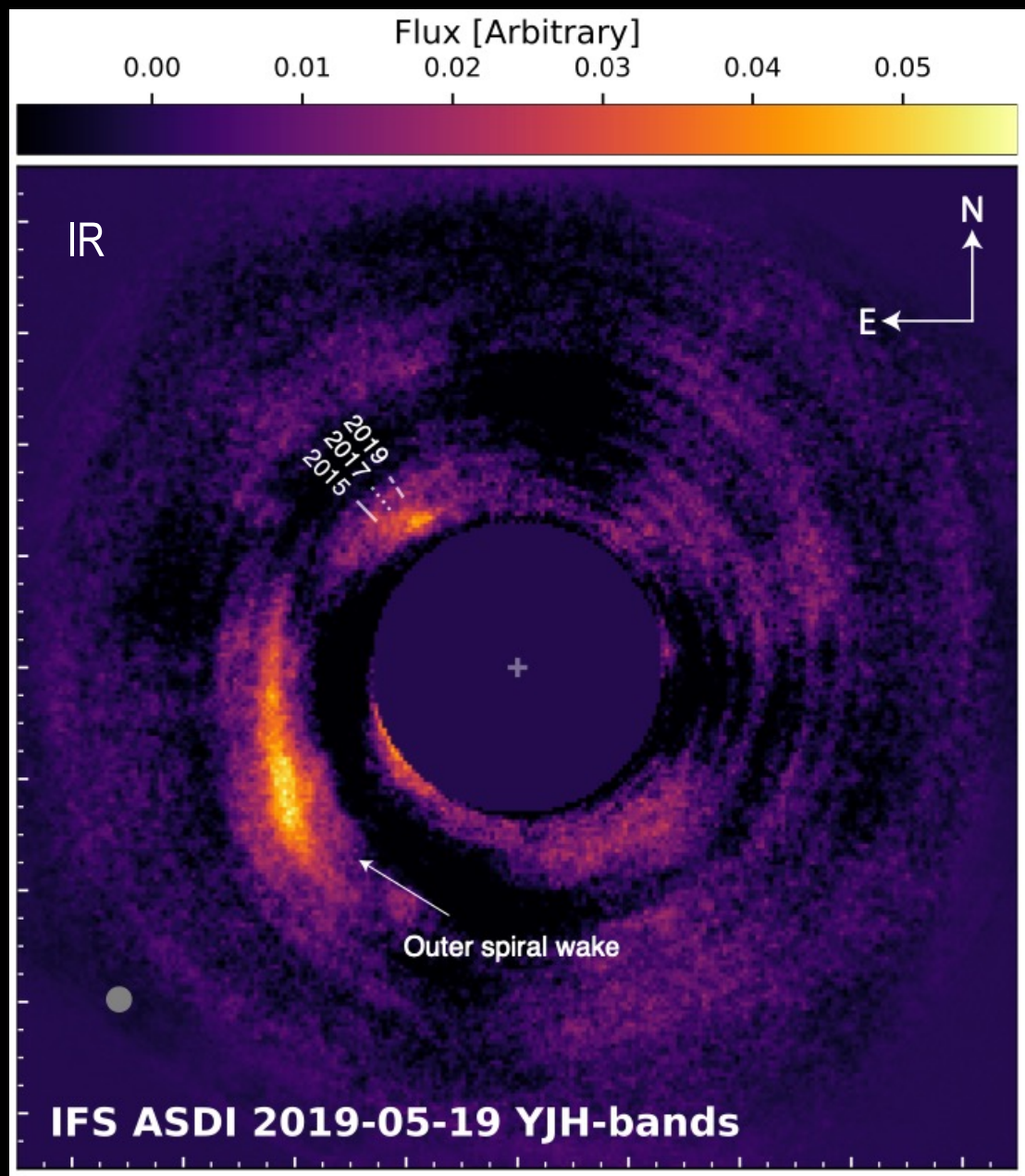
Sequential planet formation?

Protoplanet candidates inside of the cavity and gap (one of the most compelling cases).

Orbital proper (Keplerian) motions, measured during 4 years, showing the candidate is orbiting the star.

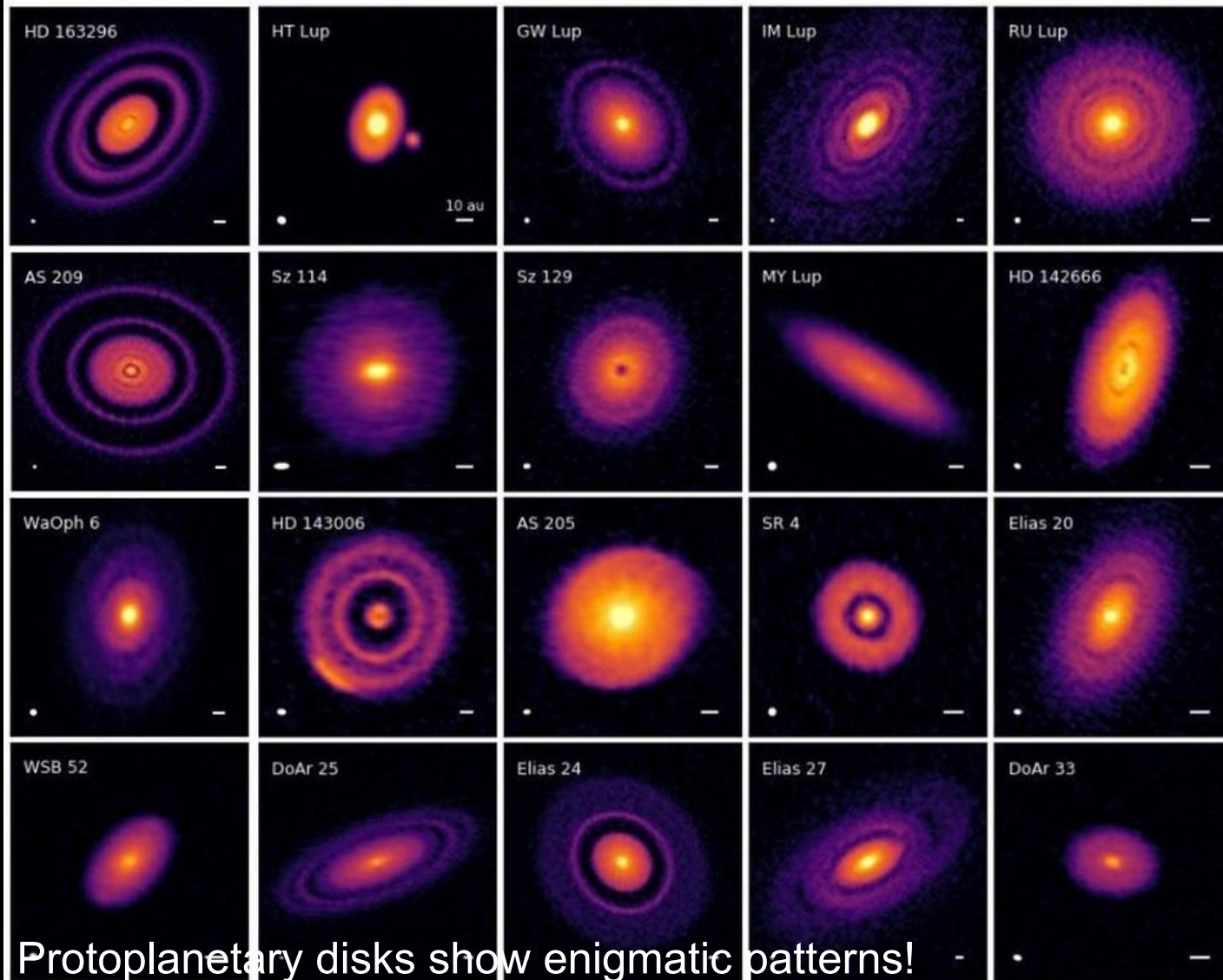
This would be the first baby planet where researchers have been able to follow its movement.

Grattom+2019,
Hammond+2023



Transitional disks: evidence of planet formation?

In the last years the focus on the study of star Formation have shifted to planet Formation



Signspots of planet formation:

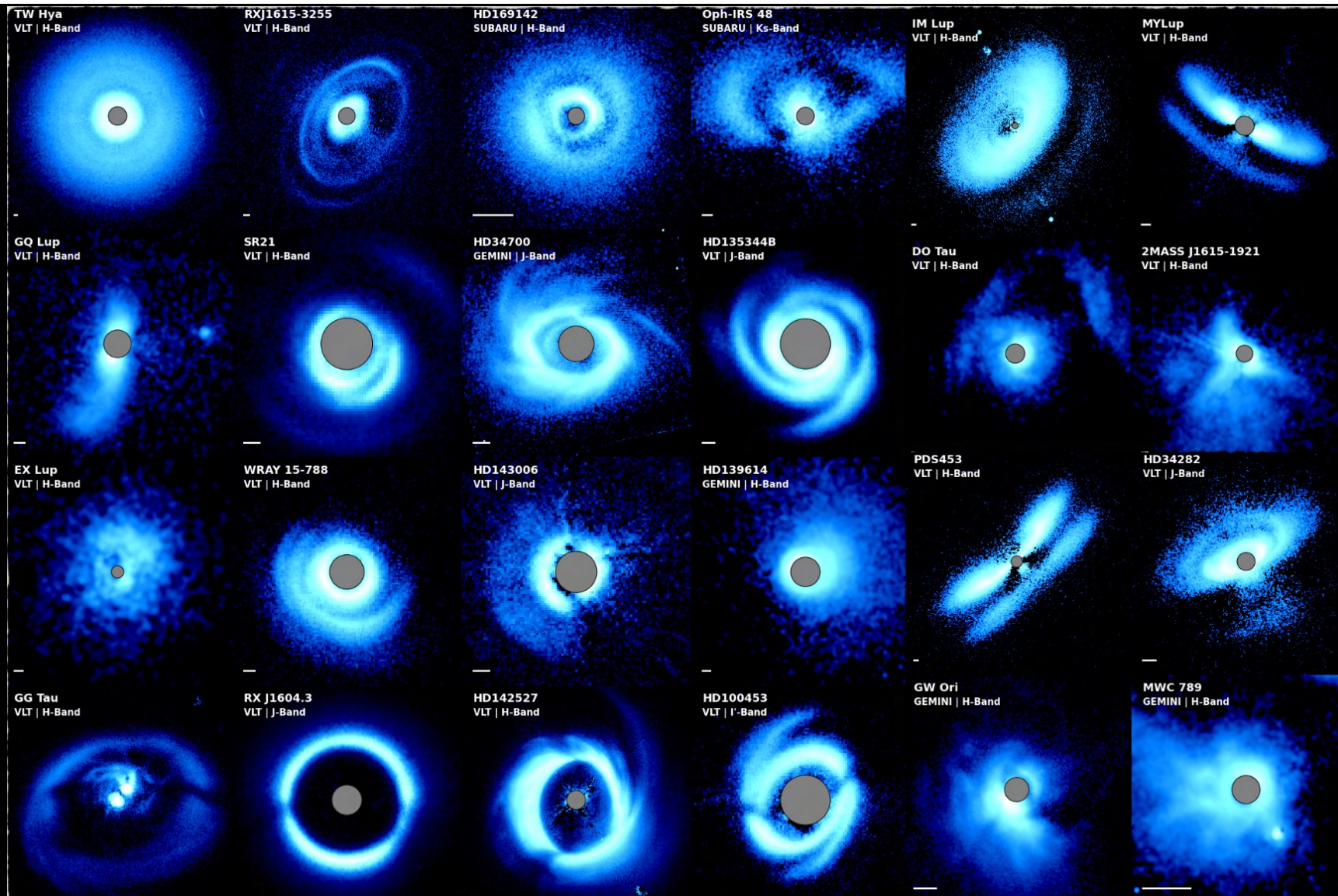
- Cavities
- Walls
- gaps
- rings
- Spiral
- Azimuthal asymmetries

reminiscent of the lopsided structures produced by dust trapping of large grains

Protoplanetary disks show enigmatic patterns!

DSHARP+2018

Transitional disks: evidence of planet formation?

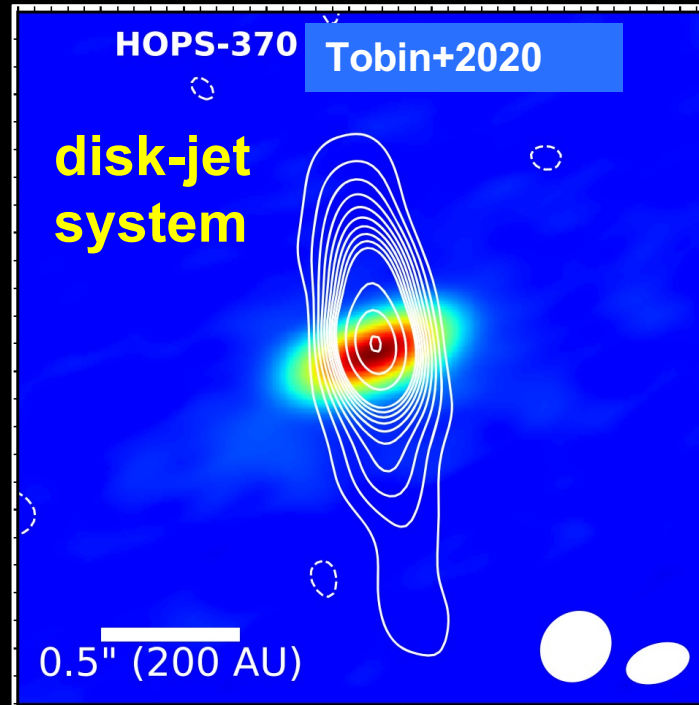
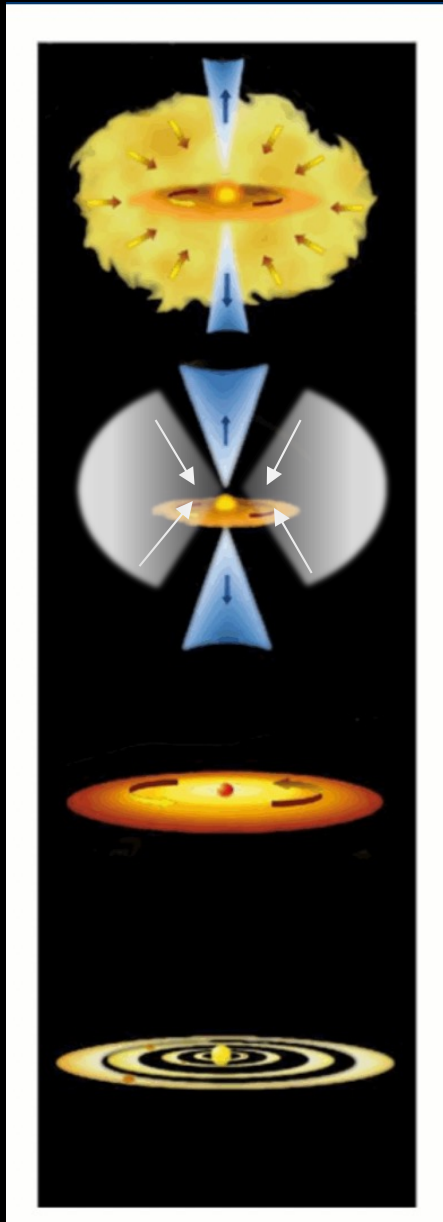


Protoplanetary disks show enigmatic patterns!
such as bright rings, empty gaps, spirals, etc

Garufi+ 2023 PPVII

Planet formation in extreme conditions

How do other scenarios fit into the standard star-formation paradigm?



Are disk/jet systems present in other star-formation scenarios?

Examples:

- Multiple star formation
- High-mass star formation

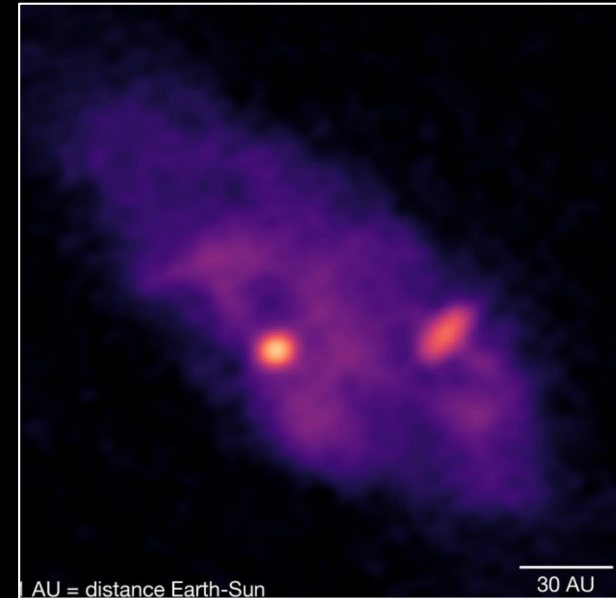
Can planetary systems be formed in these scenarios?
How these planetary systems are expected to be?

The study of these additional star/planet forming scenarios can provide the clues for a better understanding of the conspicuous diversity of all the known exoplanetary systems.

Formation of disks in multiple stellar systems

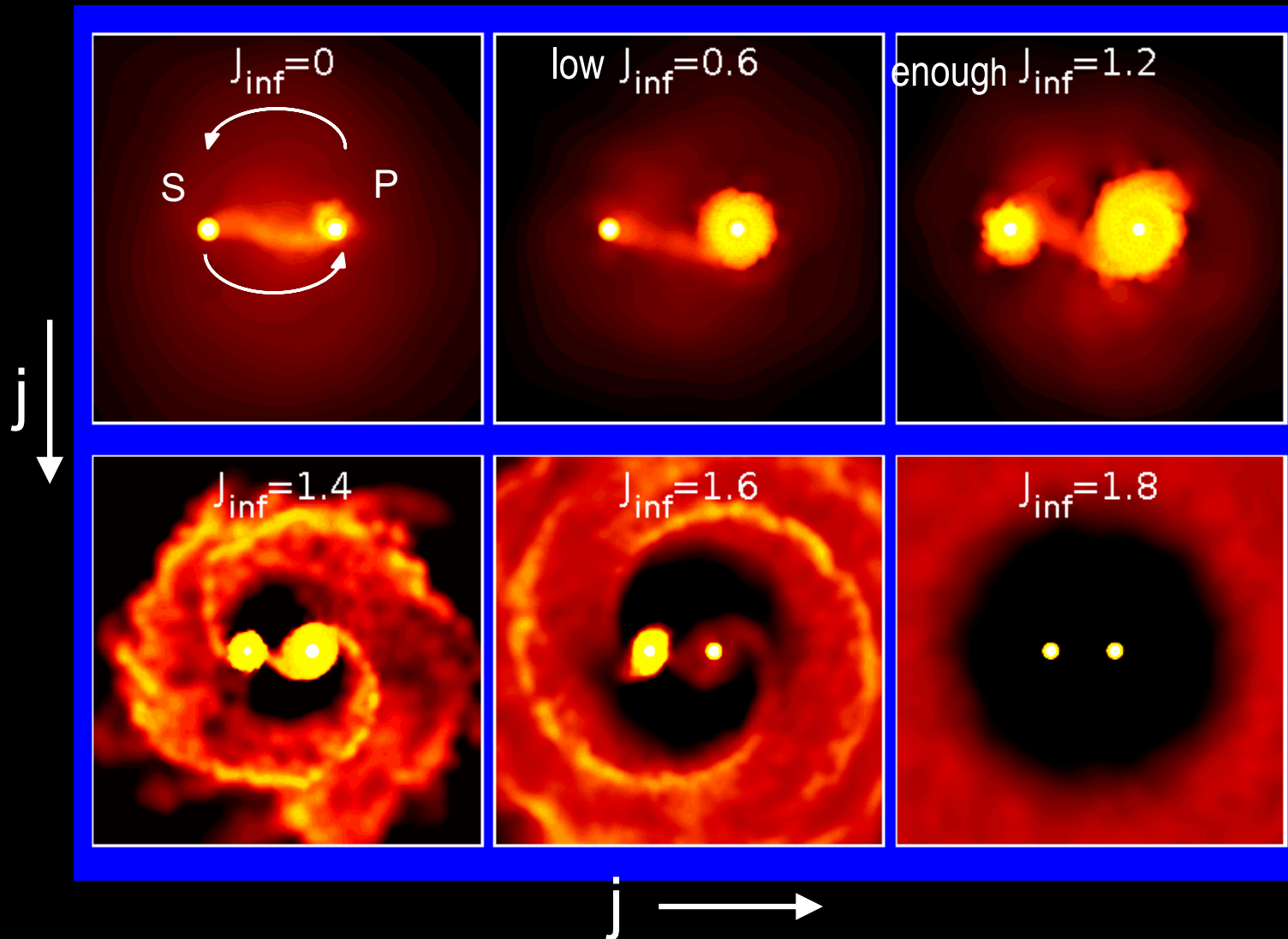
The standard star formation scheme is not directly applicable to binaries:

- Each star may develop its own disk and jet, but the whole system is **more complex**.
- There may be up to three disks: **circumstellar disks** and also a **circumbinary disk**.
- In close binaries there may be **interaction** between stars, even matter exchange.
- The evolution of one star may affect the other; one of the stars may even be formed from the **fragmentation** of its companion's disk.
- **High angular resolution** data ($<1''$) are needed to disentangle the emission of the two stars, especially for those separated by <100 au.
- **The standard paradigm of star formation may be used for the study of binary star formation, but it is more complex.**



Formation of disks in binary systems (hydrodynamic simulations)

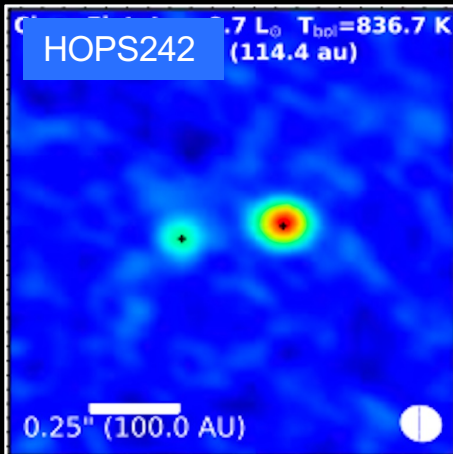
The relevant parameter for disk formation is the specific angular momentum of the infalling gas ($\mathbf{j} = \mathbf{r} \times \mathbf{v}$)



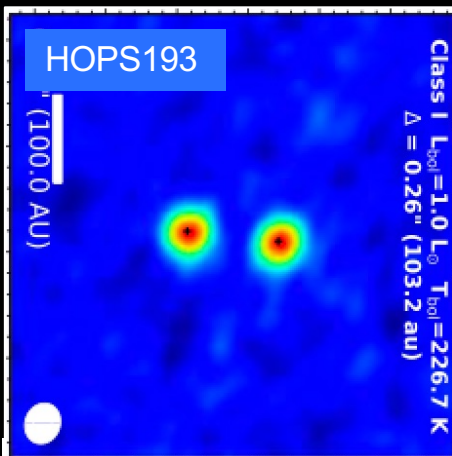
Bate 1997

In summary, a protobinary system might develop one, two or three protoplanetary disks.

Disks in binary systems (observations)

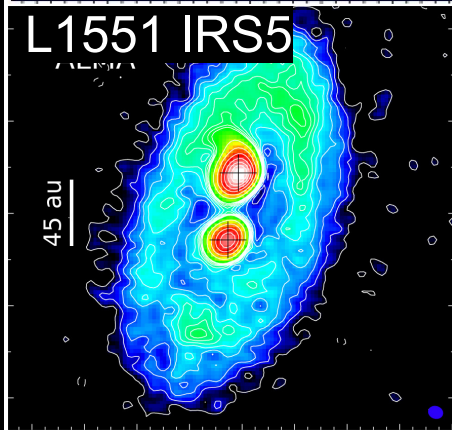


Tobin+2020

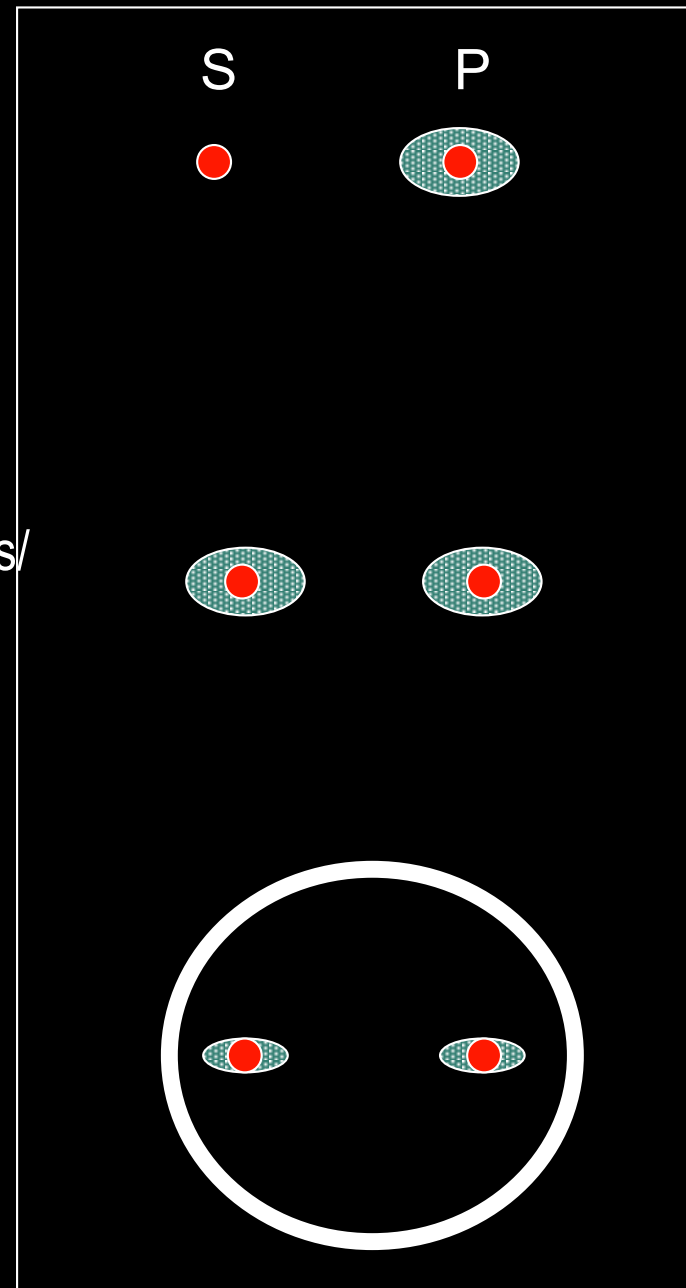


Tobin+2020

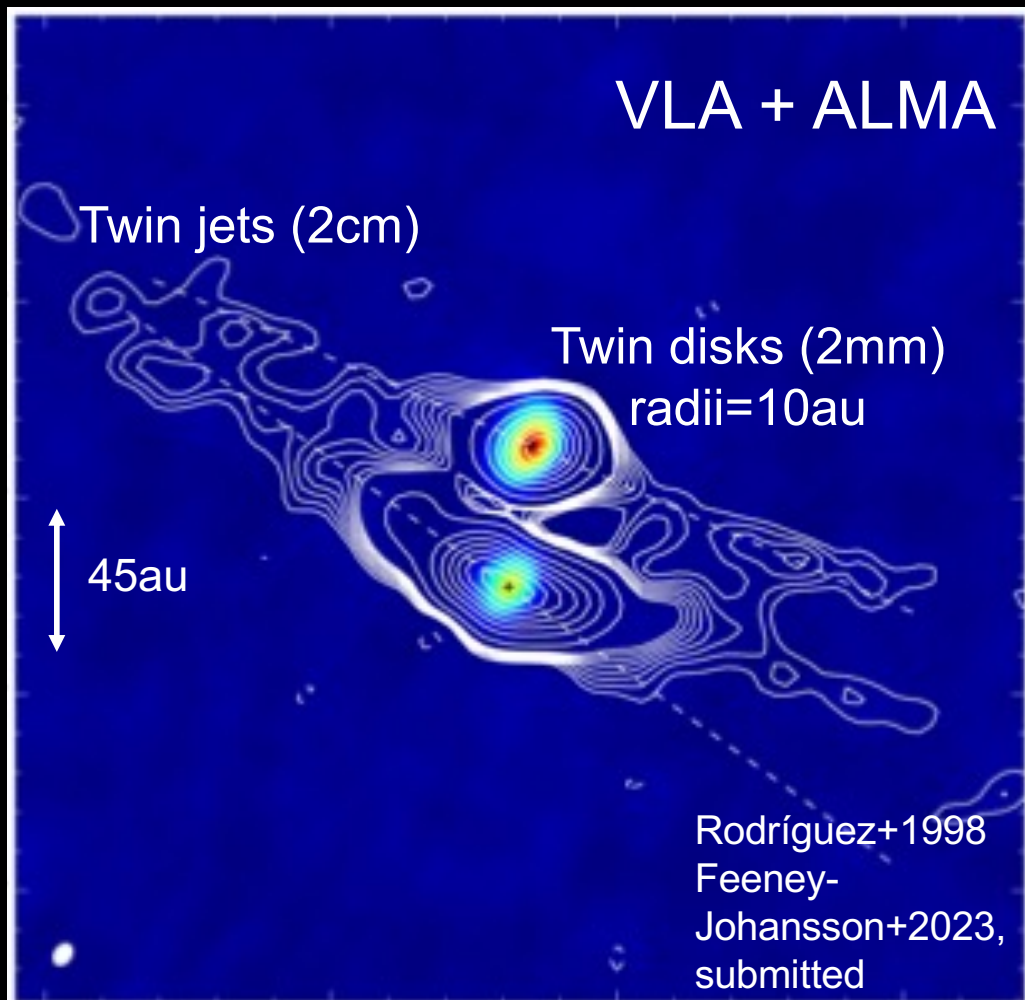
<https://spfe.es/en/catalogs/>



Takakuwa+2020



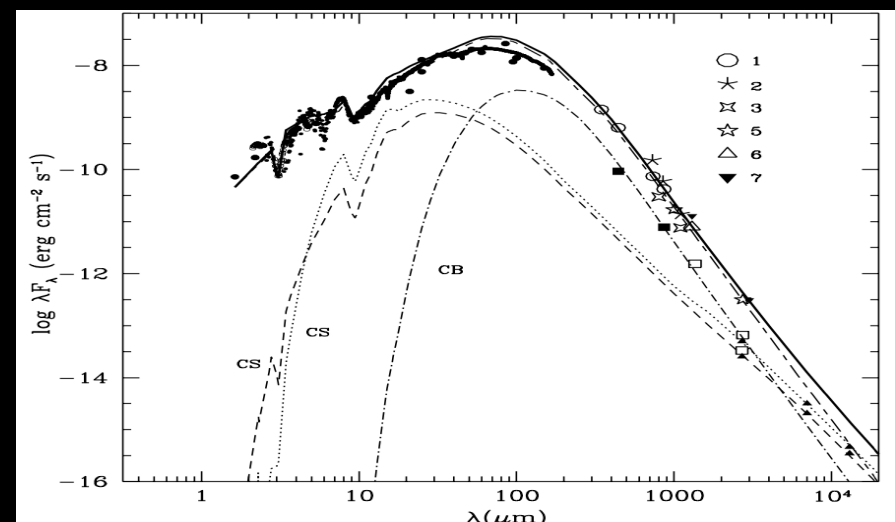
A comprehensive modeling of the protobinary system L1551 IRS 5 (Osorio et al. 2003)



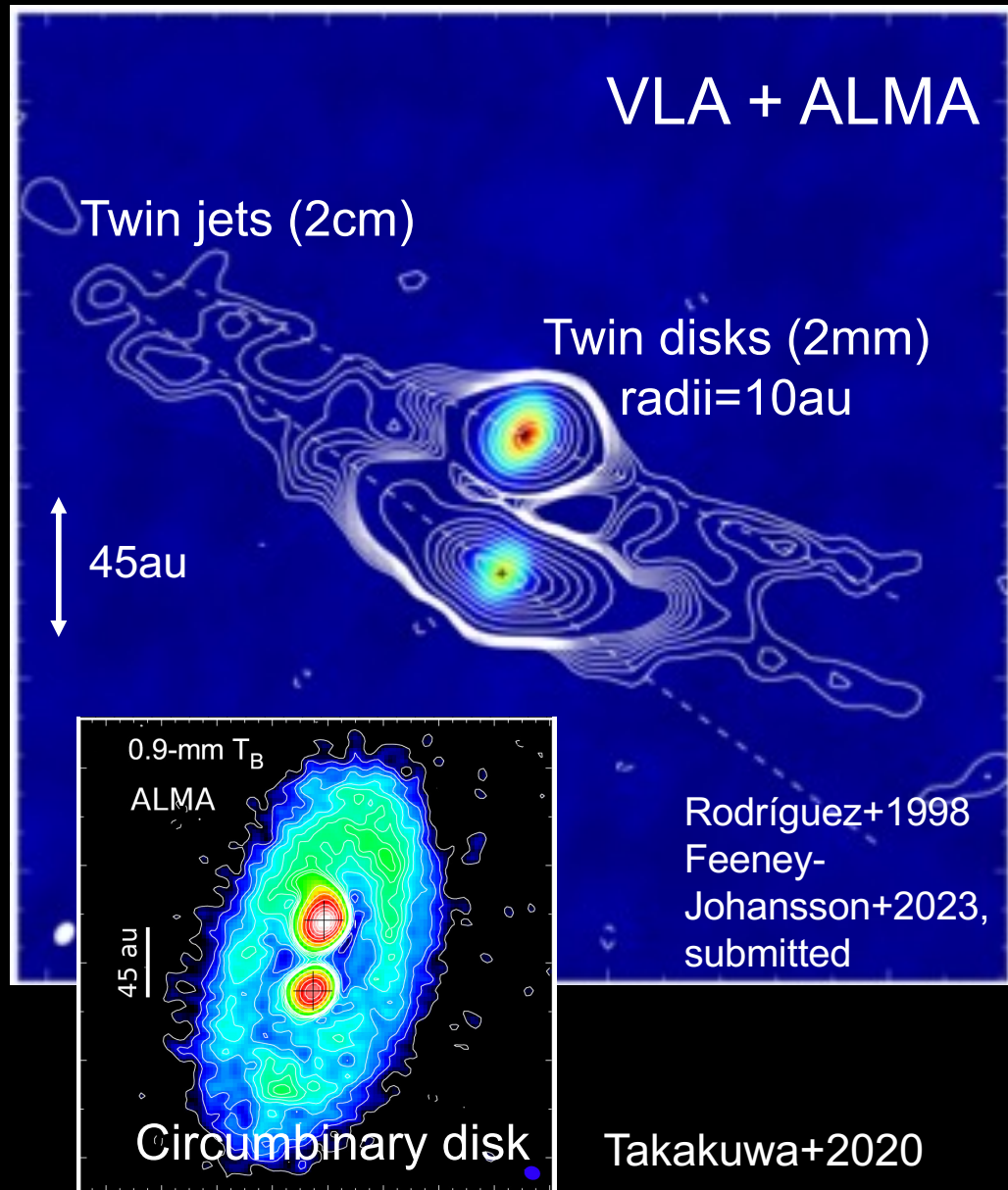
First binary where a disk/jet system in each protostar was identified. => formation similar to single stars.

Disks probably truncated by tidal interaction, they have small radii of 10 au.

Osorio+2003 modeling predicted a circumbinary disk ($R \sim 100-300$ au).



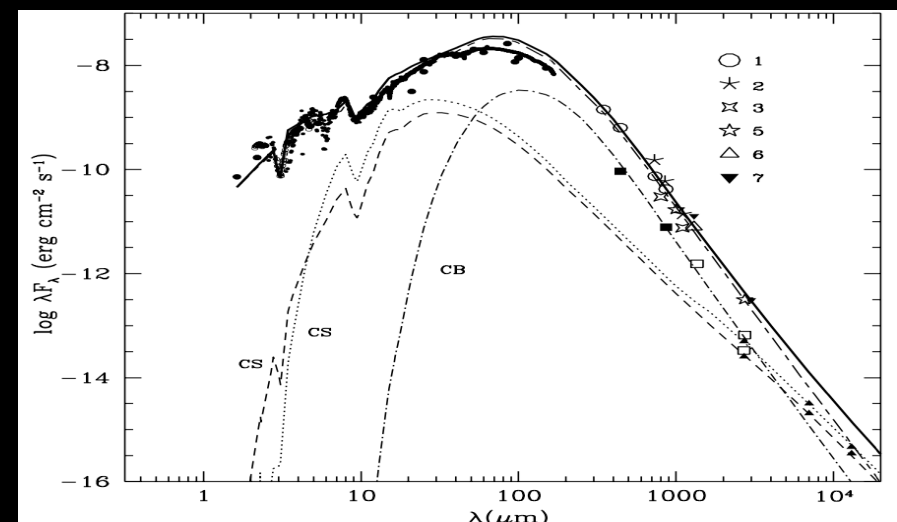
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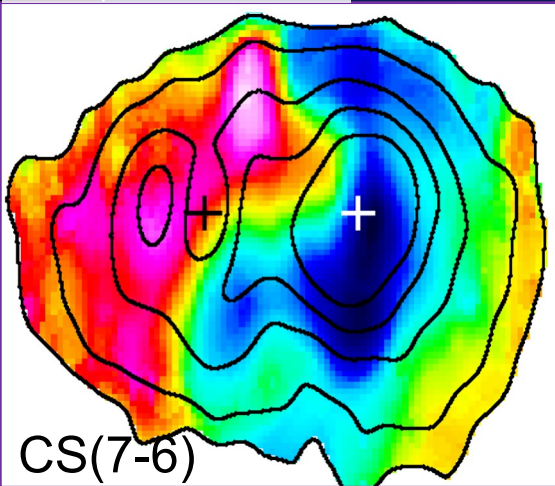
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A comprehensive observational study of the protobinary system SVS 13: a circumbinary disk in the making

(Díaz-Rodríguez+2022)

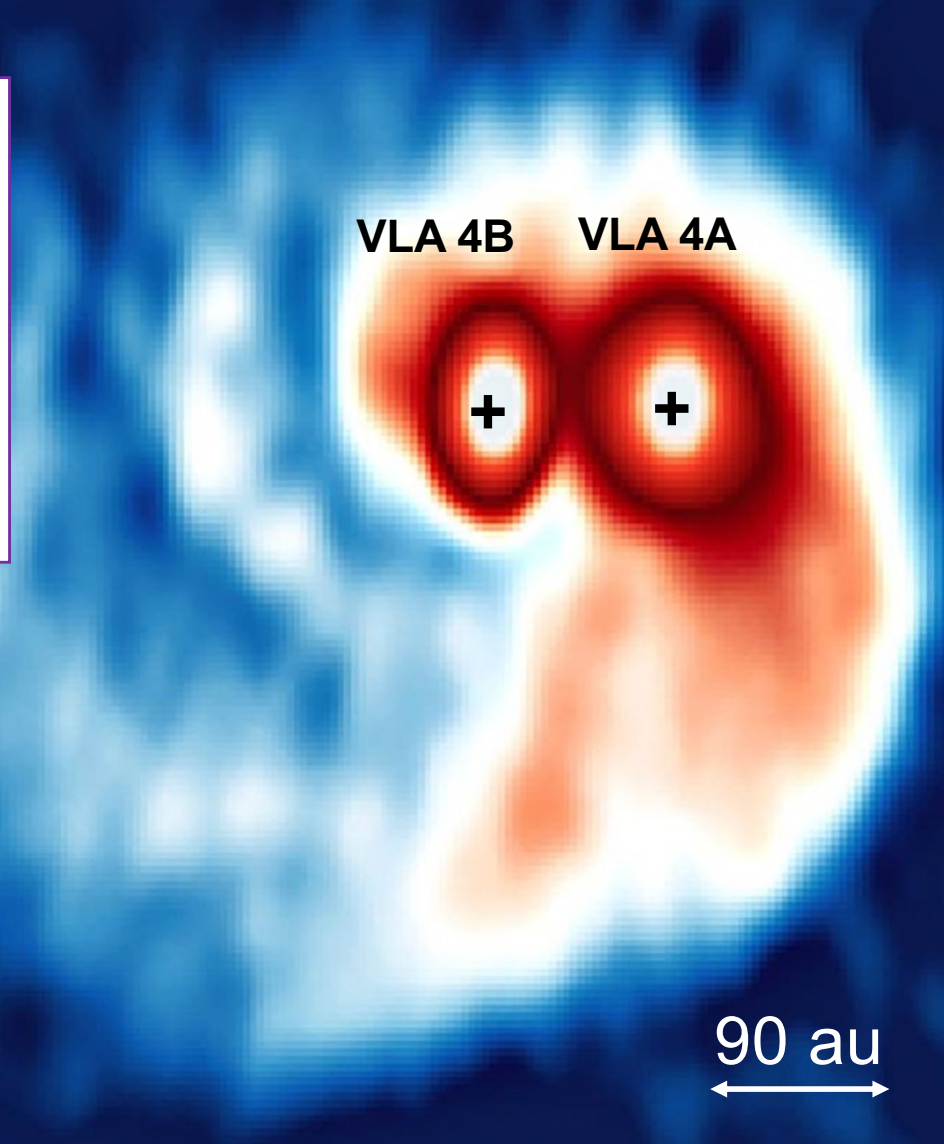
Circumbinary
rotation+infall



Color: mean velocity
Contours: integrated intensity



Continuum 0.9 mm



Circumstellar disks

R ~10 au (dust)
~30 au (gas)

VLA 4A

$M_* = 0.3 M_\odot$

$M_{\text{disk}} = 4-9 M_J$

$T_{\text{dust}} = 140-300 \text{ K}$

VLA 4B

$M_* = 0.6 M_\odot$

$M_{\text{disk}} = 9-30 M_J$

$T_{\text{dust}} = 140-400 \text{ K}$

Circumbinary disk

R ~500 au,

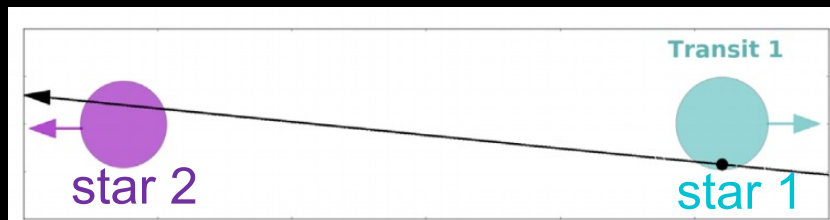
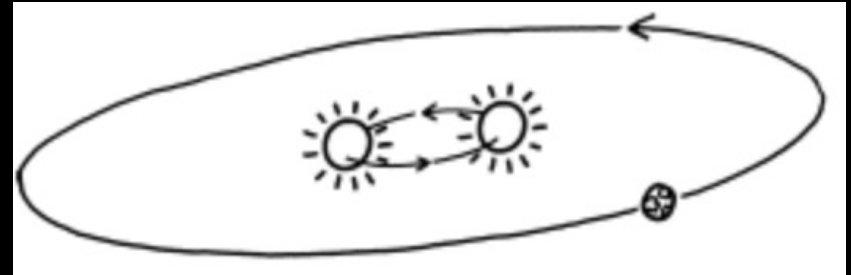
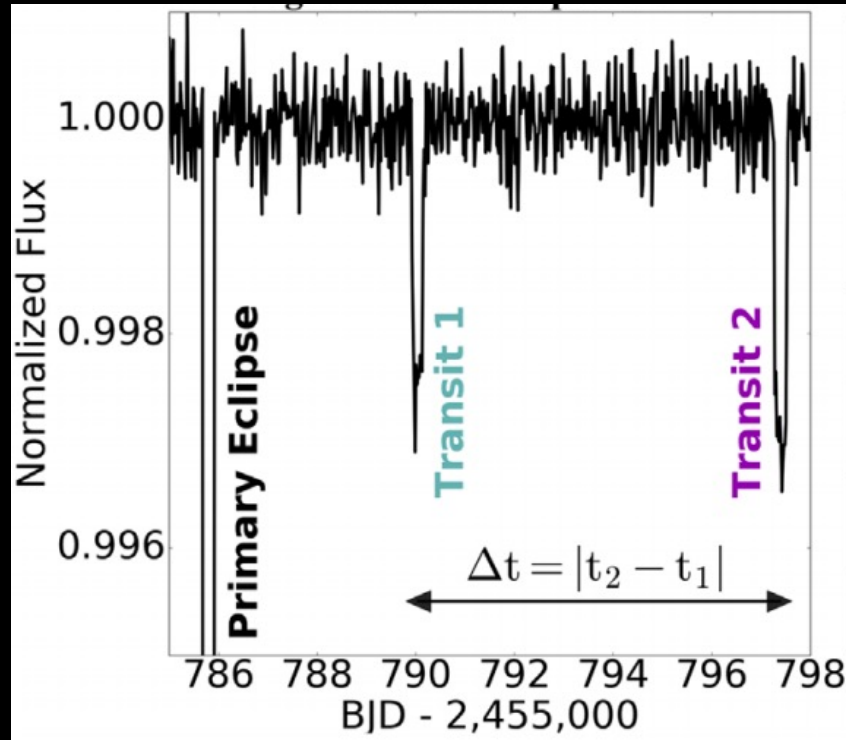
$M_{\text{disk}} = 50 M_J$

T ~140 K

counter-clockwise rotation

Detection of circumbinary exoplanets

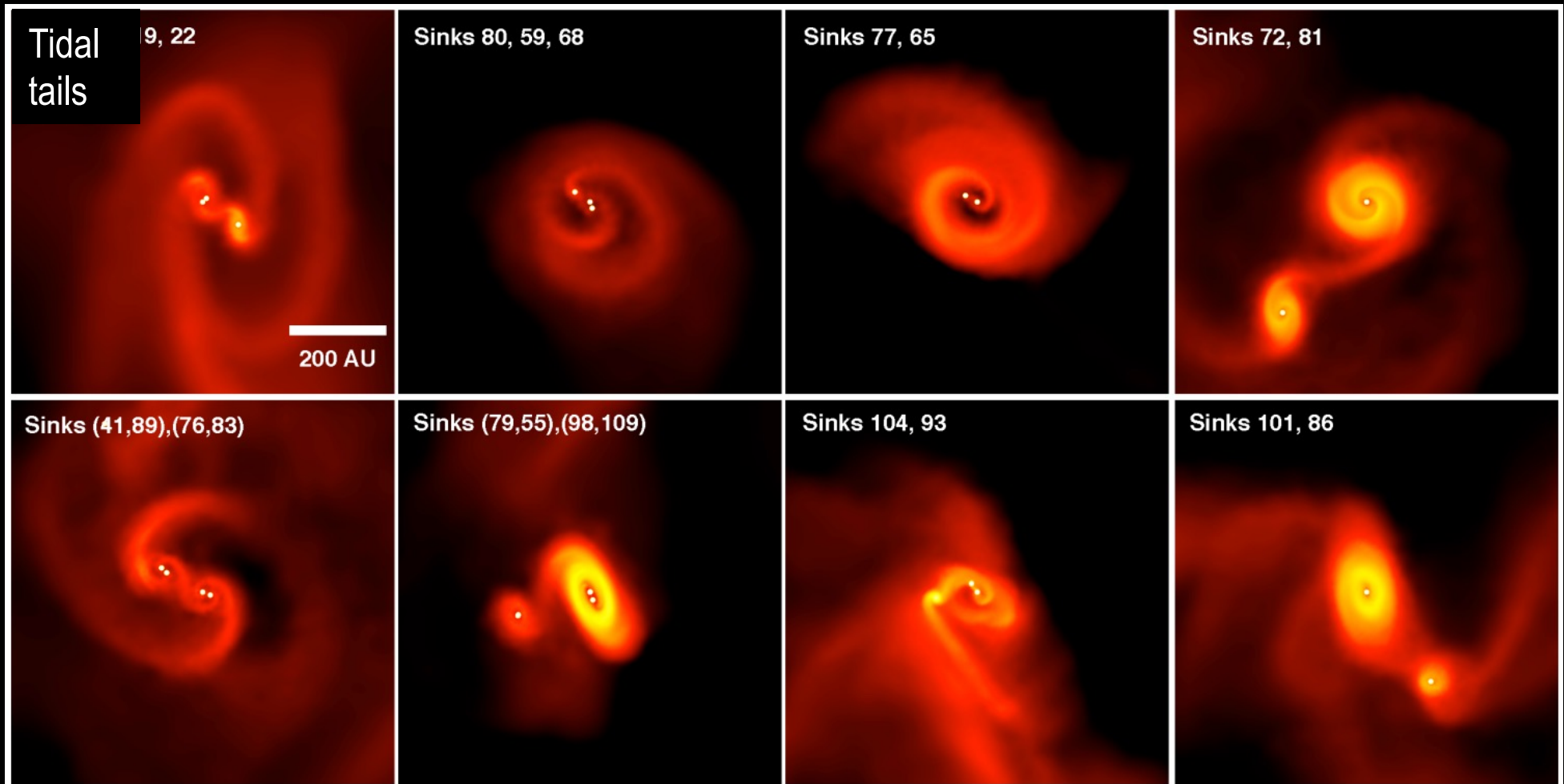
Kepler 34 (TESS data; Kostov+2021)



Eclipse of the primary star by the secondary star (big dip), transit of the planet across the primary (transit 1), and across the secondary star (transit 2)

Simulations of disk interaction in multiple systems: matter exchange, tidal interaction, spirals, fragmentation.

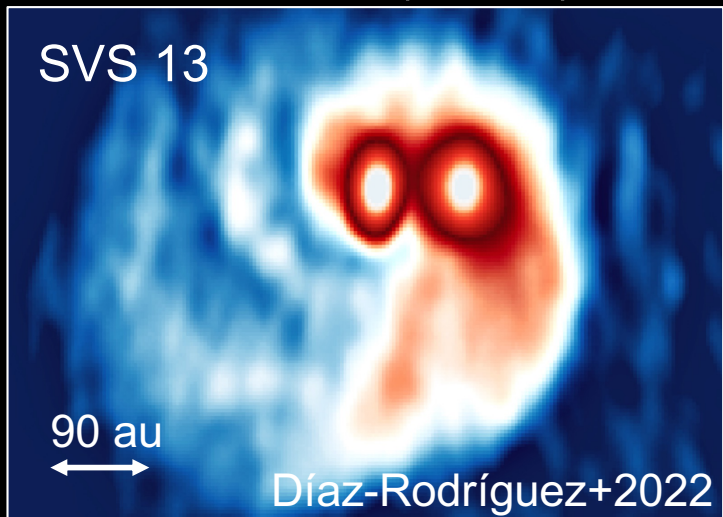
These panels are close-ups of several groups of stars, corresponding to a snapshot of the whole system at a given time



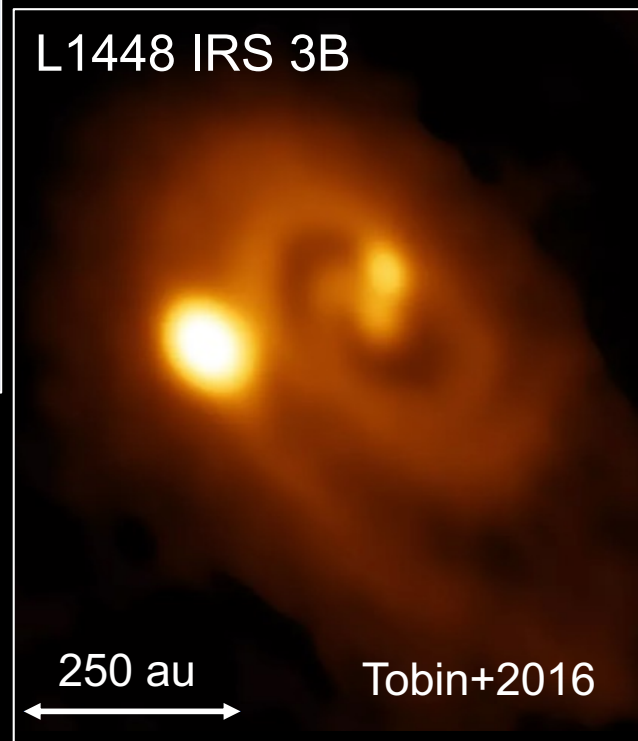
(Bate 2018)

Observations of disk interaction in multiple systems

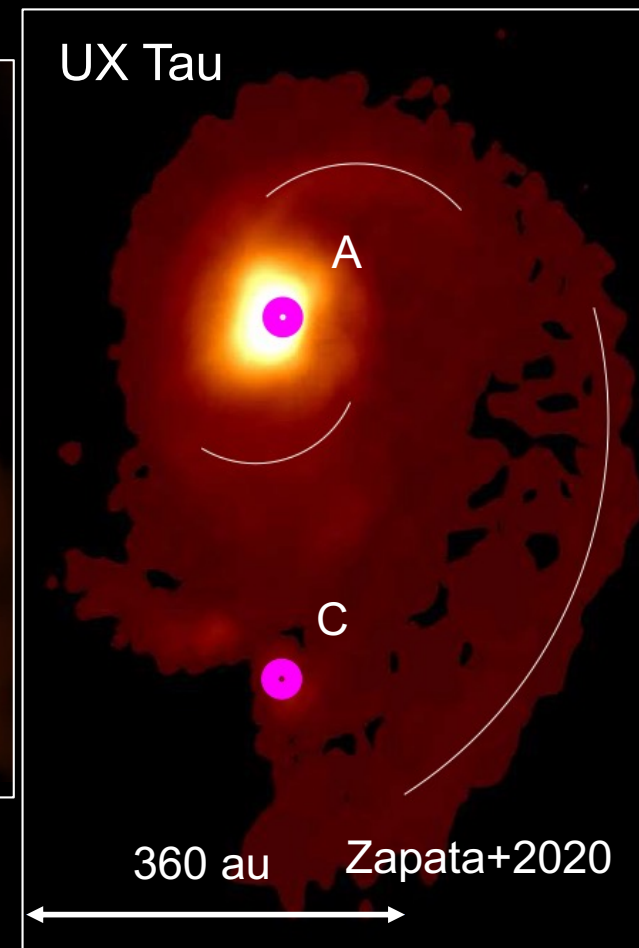
Two stars (spirals)



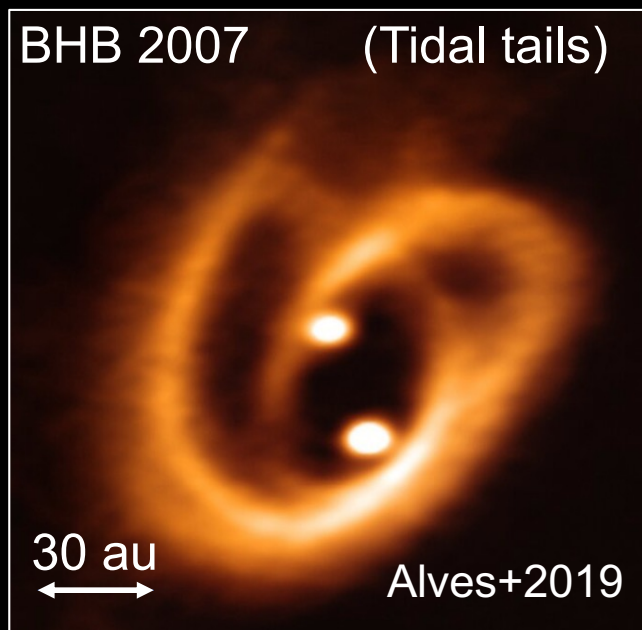
Three stars
(fragmentation)



High-order multiple
stellar systems (flybys)

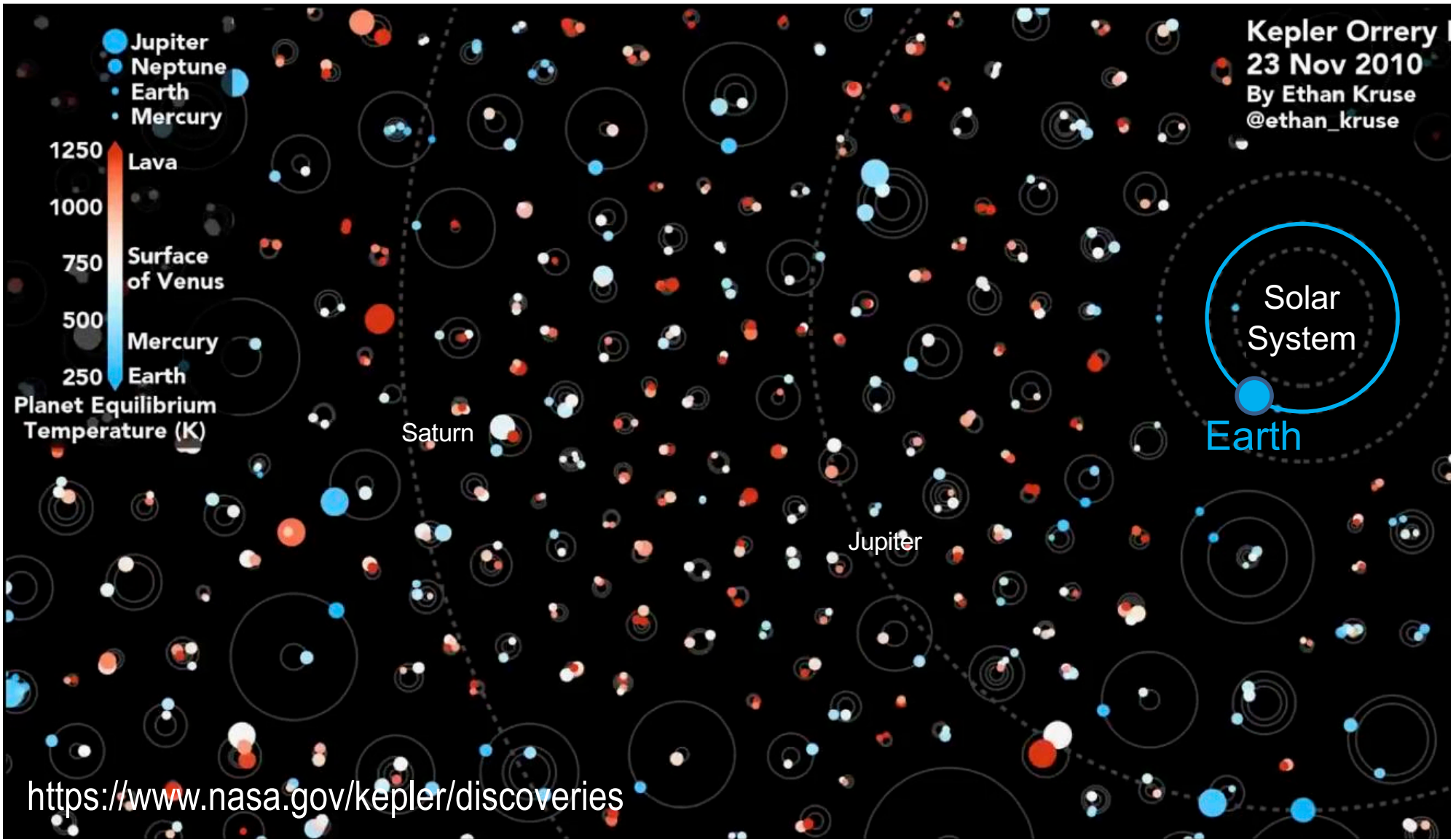


BHB 2007 (Tidal tails)



The interaction of a triple system usually results in the ejection of the less massive star.

Diversity of exoplanetary system architectures



Closely-packed compact exoplanetary systems

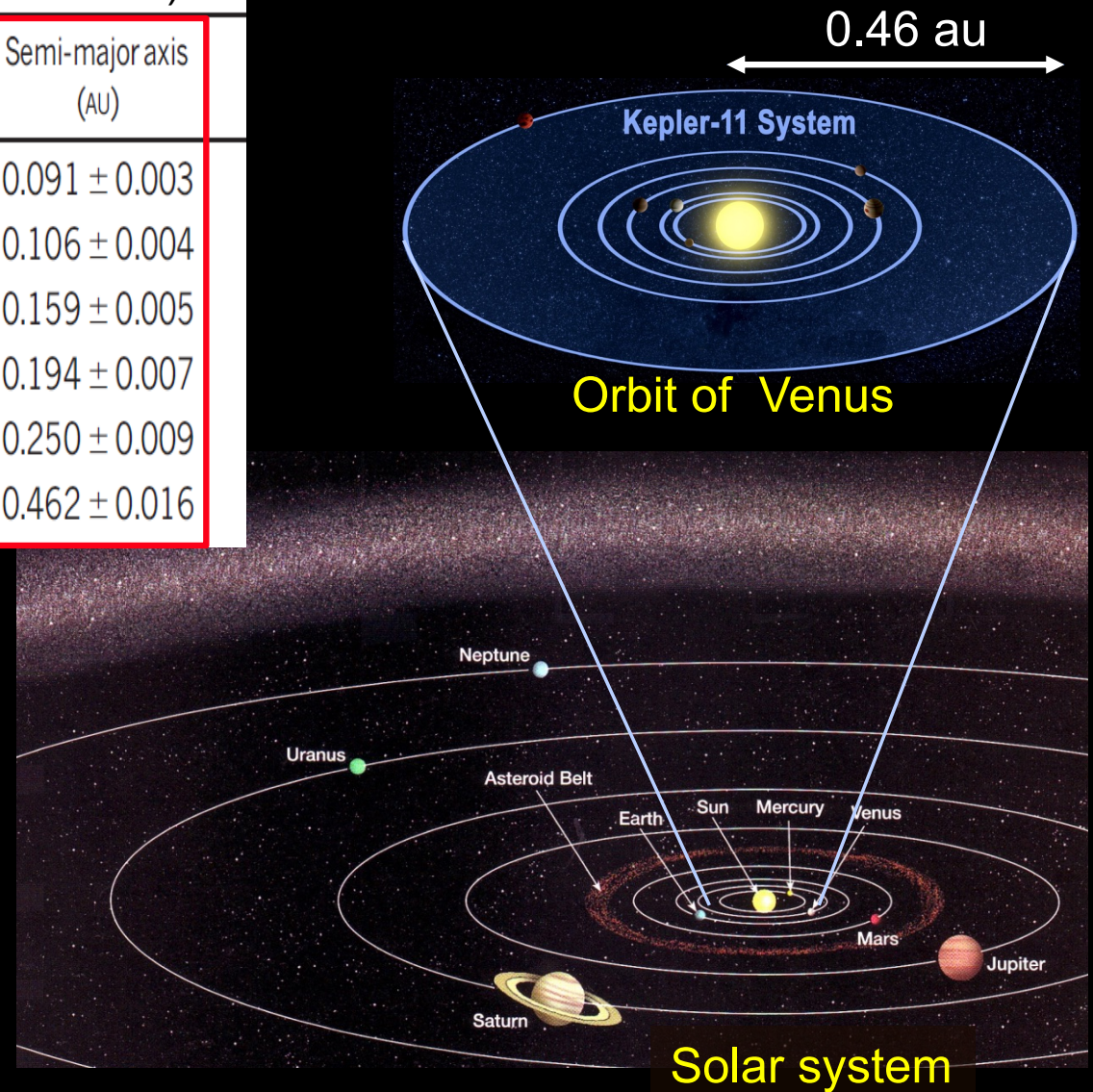
Kepler-11 Planet Properties (Lissauer et al 2011)

Planet	Period (days)	Epoch (BJD)	Semi-major axis (AU)
b	10.30375 ± 0.00016	$2,454,971.5052 \pm 0.0077$	0.091 ± 0.003
c	13.02502 ± 0.00008	$2,454,971.1748 \pm 0.0031$	0.106 ± 0.004
d	22.68719 ± 0.00021	$2,454,981.4550 \pm 0.0044$	0.159 ± 0.005
e	31.99590 ± 0.00028	$2,454,987.1590 \pm 0.0037$	0.194 ± 0.007
f	46.68876 ± 0.00074	$2,454,964.6487 \pm 0.0059$	0.250 ± 0.009
g	118.37774 ± 0.00112	$2,455,120.2901 \pm 0.0022$	0.462 ± 0.016

Kepler 11: 6 planets at $r < 0.46$ au

Trappist 1: 6 planets at $r < 0.7$ au

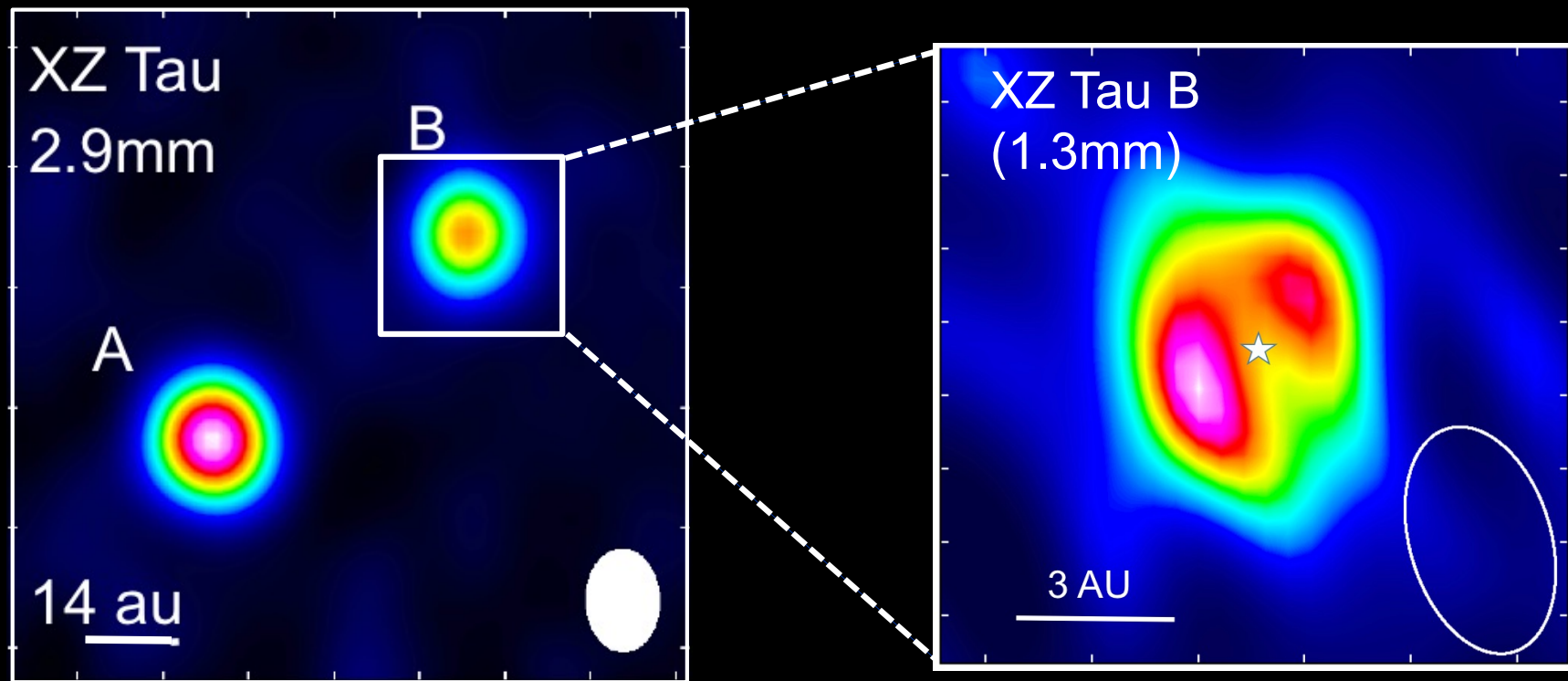
Do compact exoplanetary systems originate from dwarf protoplanetary disks?



A dwarf protoplanetary disk in XZ Tau B

(Osorio et al. 2016)

ALMA Long Base Science Verification observations



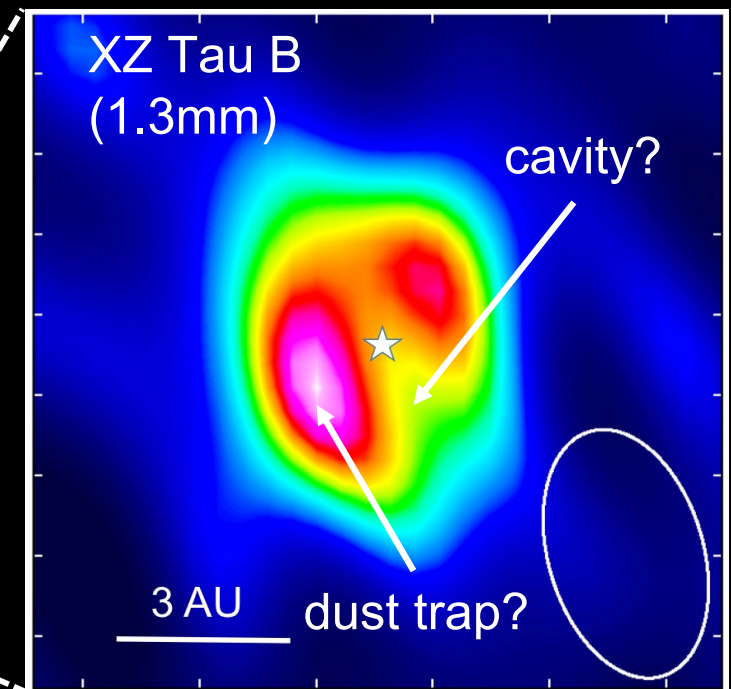
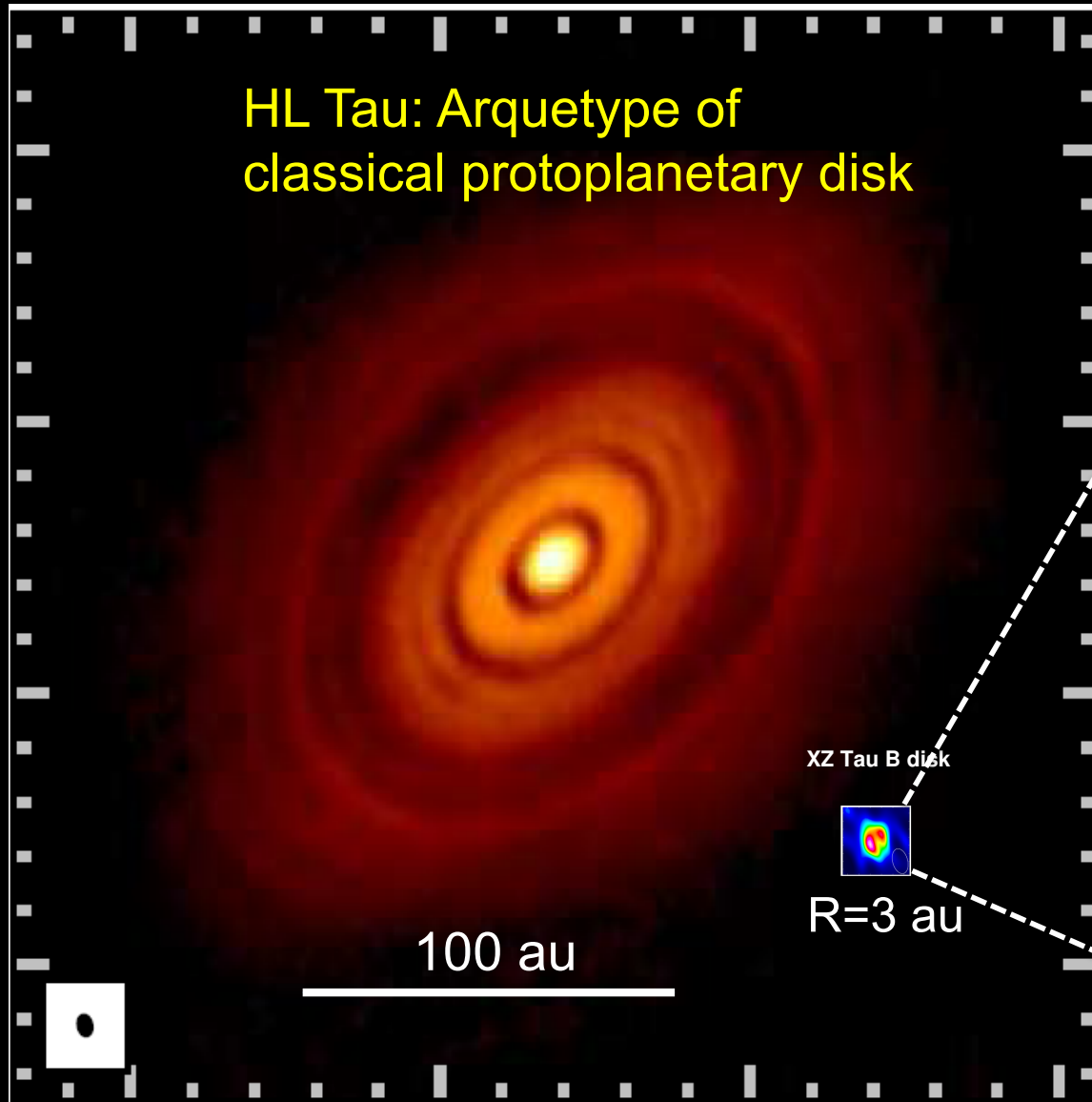
beam = $0.09'' \times 0.06'' = 13\text{au} \times 8\text{au}$ beam = $0.03'' \times 0.02'' = 4.2\text{au} \times 2.8\text{au}$

The large improvement in the angular resolution of ALMA allow us to probe scales of only 2au

Comparison between HL Tau and XZ Tau B disks

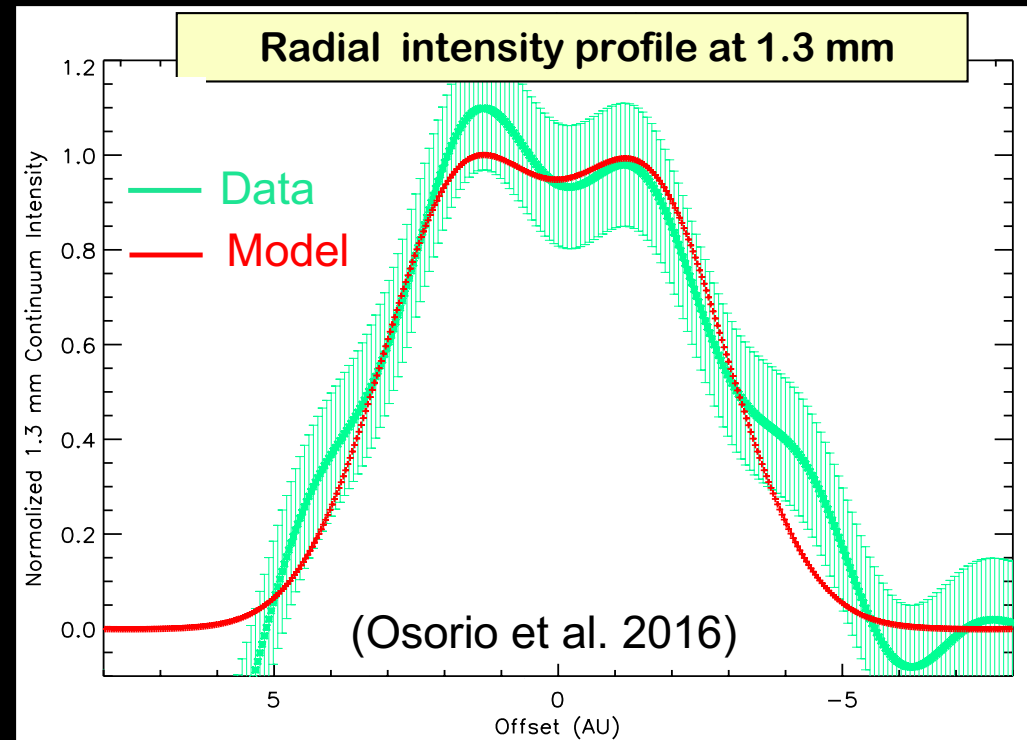
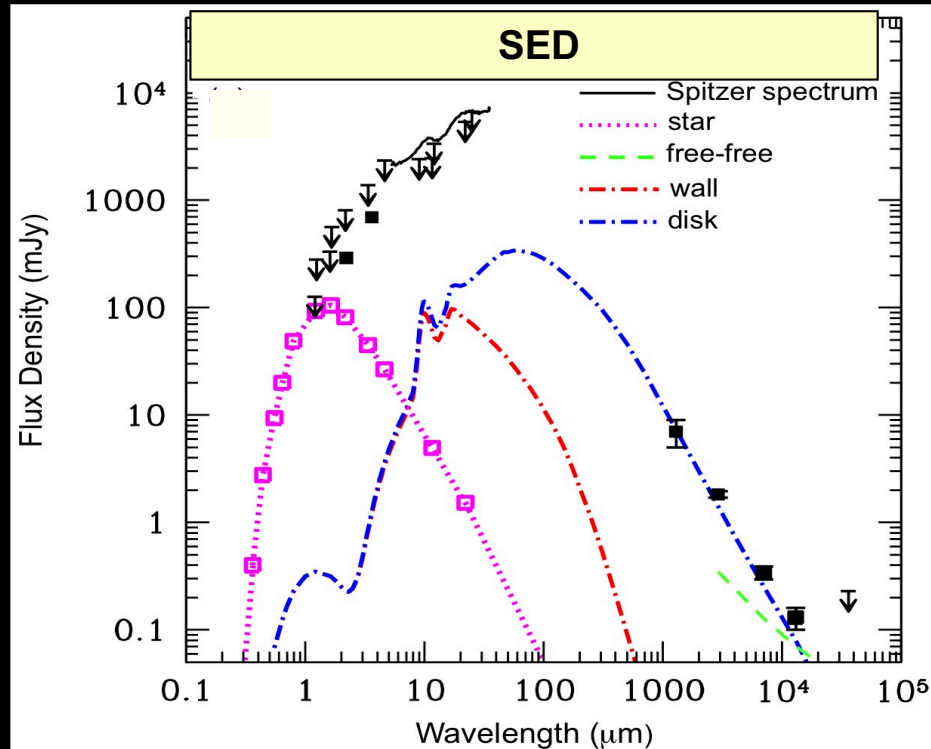
HL Tau: Arquetype of classical protoplanetary disk

XZ Tau B: Arquetype of dwarf disk



we think, that just as HL Tau is considered an archetype for classic disks, XZ Tau B may be a prototype of small disks.

Modeling of the dwarf disk in XZ Tau B



XZ Tau B Disk

$$R_{\text{disk}} = 3.4 \text{ au}$$

$$M_{\text{disk}} = 9 M_{\text{Jup}}$$

$$\Sigma(r=1 \text{ ua}) = 2700 \text{ g/cm}^2$$

Classical Disk

$$50\text{-}100 \text{ au}$$

$$50\text{-}500 M_{\text{Jup}}$$

$$100\text{-}1000 \text{ g/cm}^2$$

The model predicts a disk with enough mass to form planets ($9 M_{\text{Jup}}$) even though it is very small.

A large population of these small disks could exist but they have remained hidden, because, so far only few extremely high angular resolution observation have been performed.

Formation of massive stars

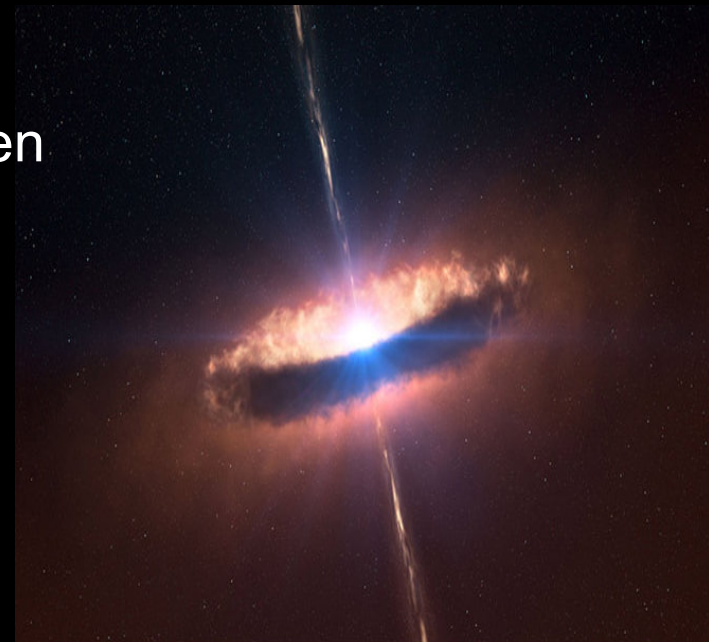
If massive stars ($>8M_{\odot}$) could be formed via accretion, with a disk/jet system, then they could form a planetary system, however:

-The contraction time that a massive star requires to reach the main sequence is very short (10^4 yr).

-Massive stars reach the main sequence when they are still accreting material.

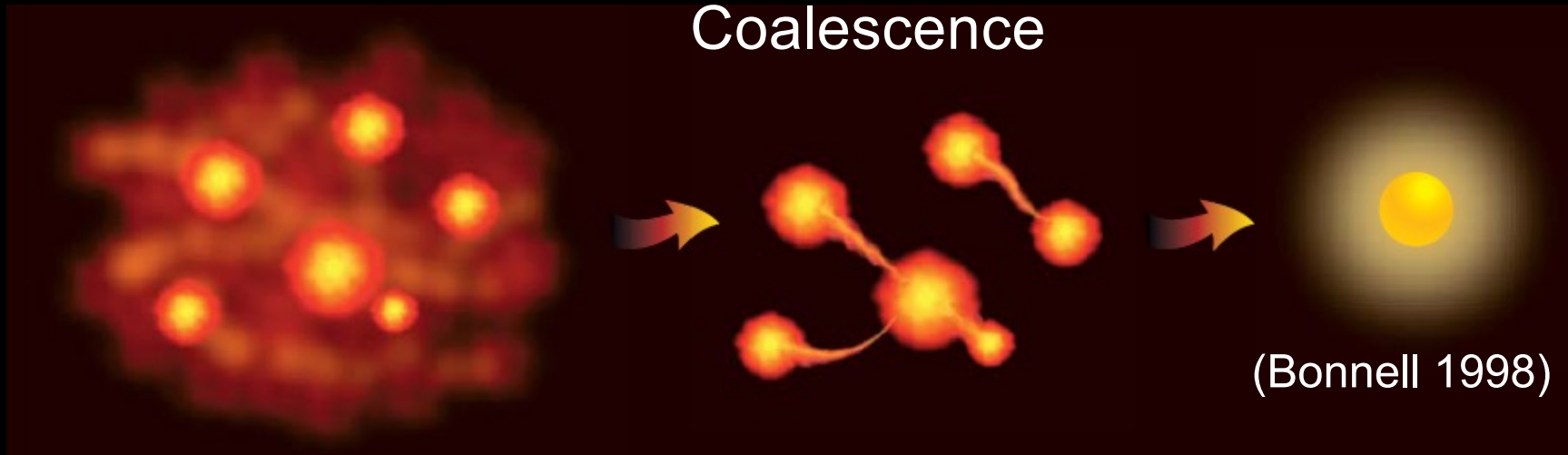
-Massive stars develop HII regions with a high UV photon rate. **Radiation pressure may halt the collapse, therefore the star cannot grow larger to reach a high mass.**

So, in recent years it has been debated if it is possible to form high-mass stars via accretion.

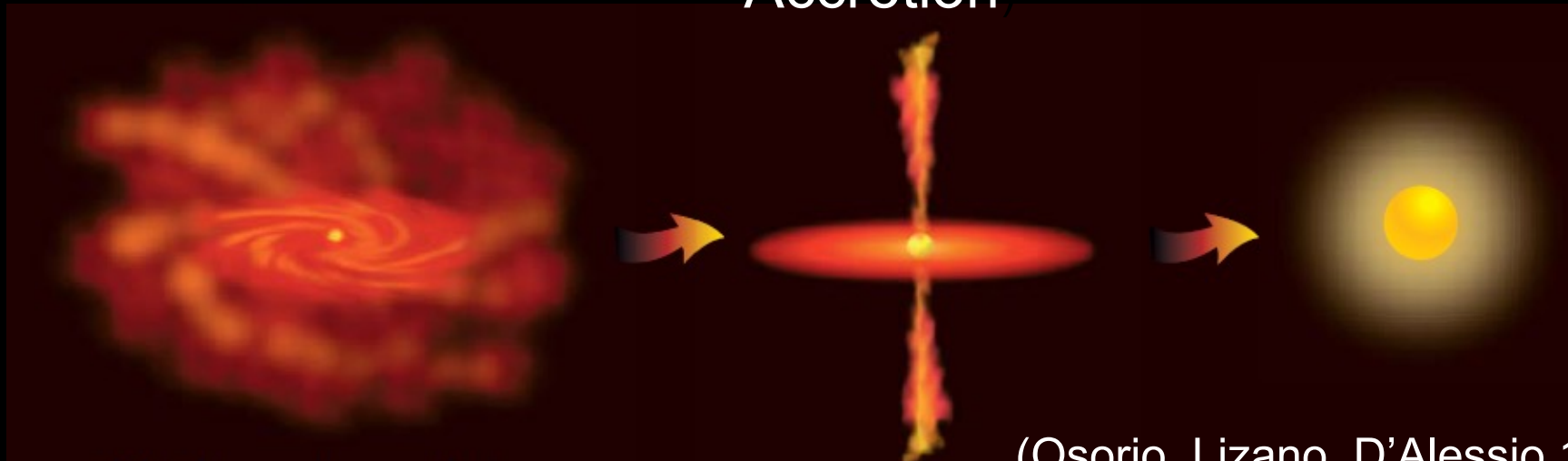


Star formation of high mass protostars: Accretion or coalescence?

Coalescence

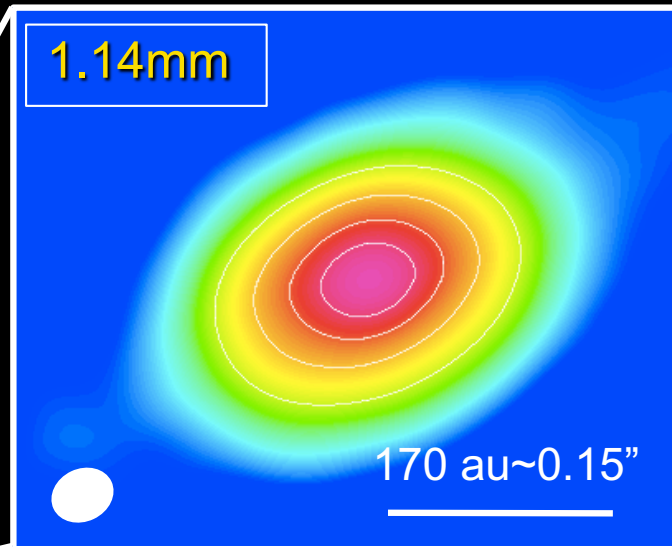
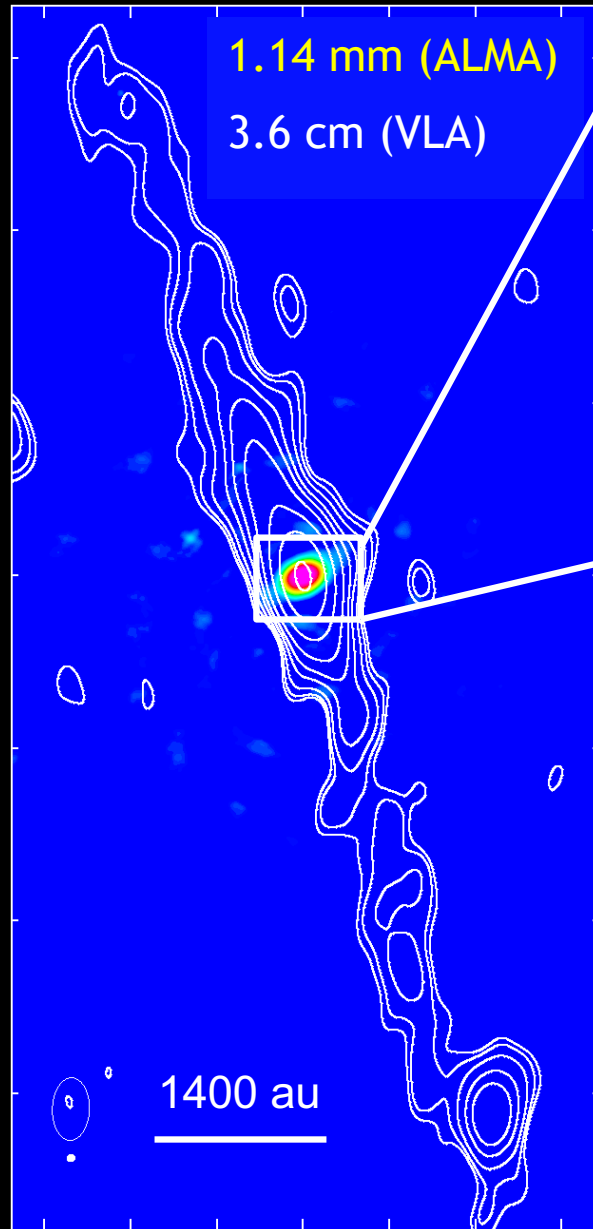


Accretion



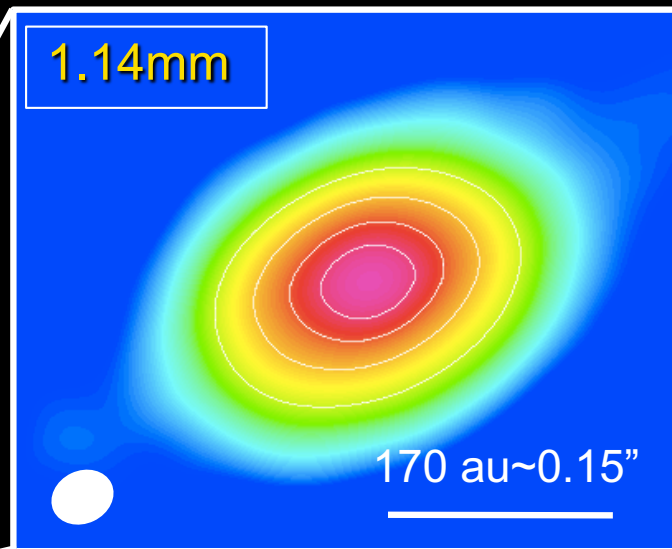
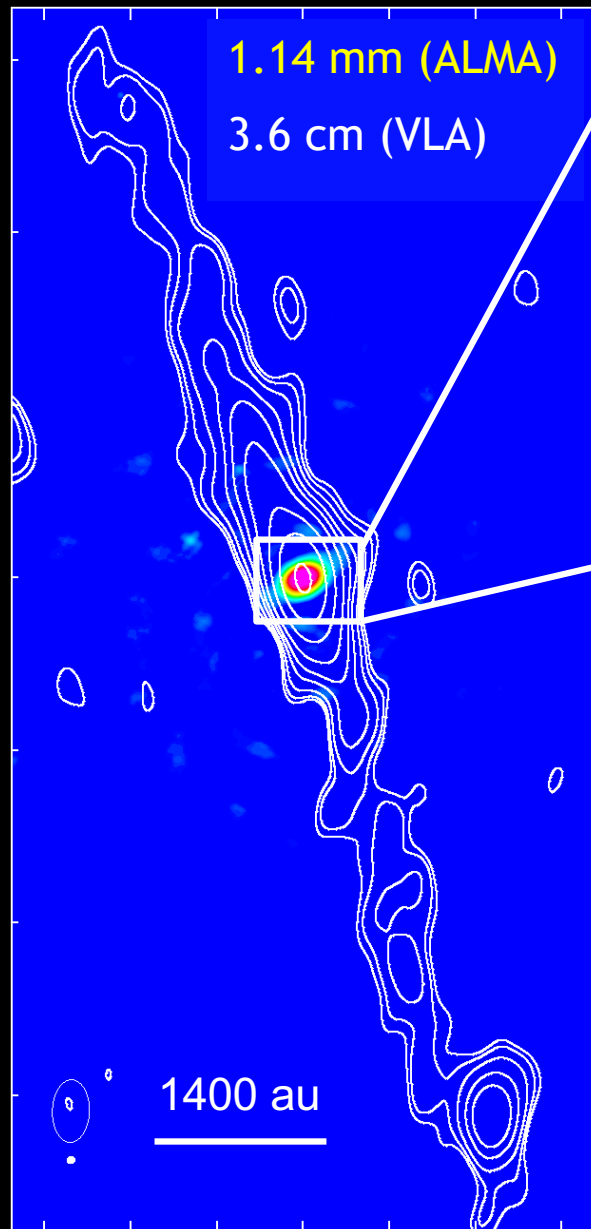
(Osorio, Lizano, D'Alessio 1999)

The disk/jet system of the HH80 massive protostar

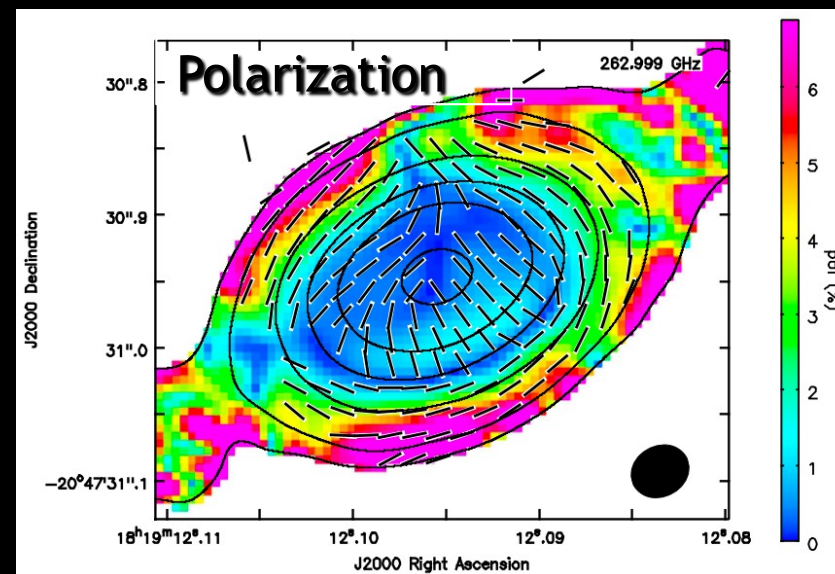


Resolved disk,
perpendicular to the
jet, $R \sim 200$ au
→ **bona-fide disk**

The disk/jet system of the HH80 massive protostar



Resolved disk,
perpendicular to the
jet, $R \sim 200$ au
→ **bona-fine disk**



The polarization
pattern indicates that
there is no settling
onto the disk
midplane, and that
the size of the dust
grains is
 ~ 50 - $500 \mu\text{m}$.
→ Inputs for the
models.

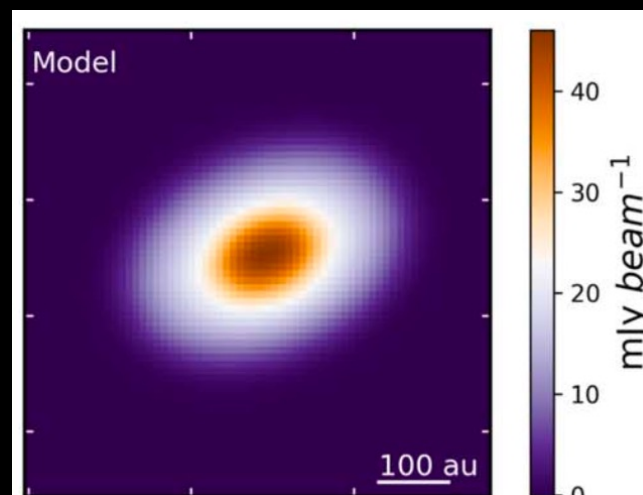
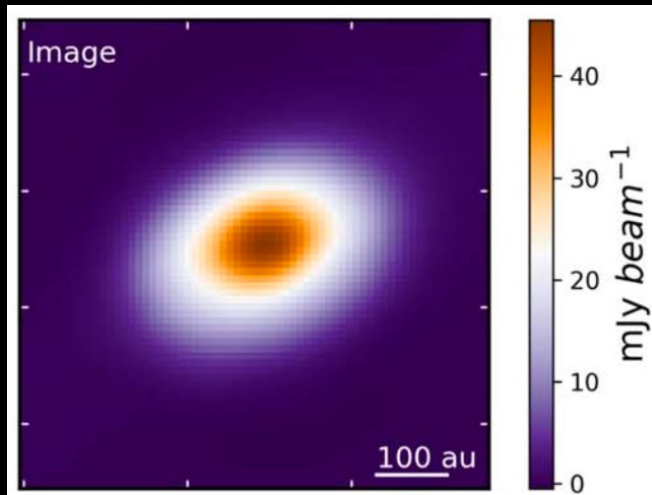
Girart et al. 2018; Añez-López, Osorio et al. 2020

Modeling of the HH80 disk (massive protostar)

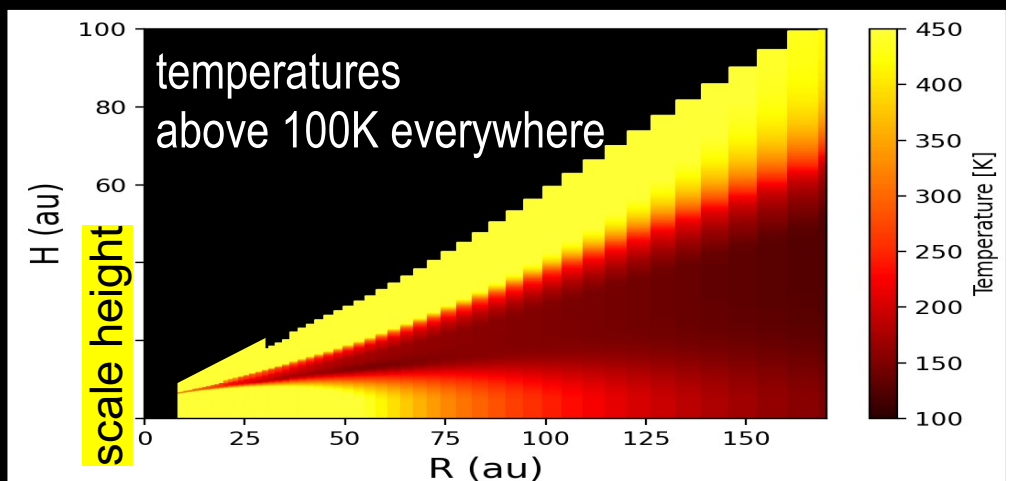
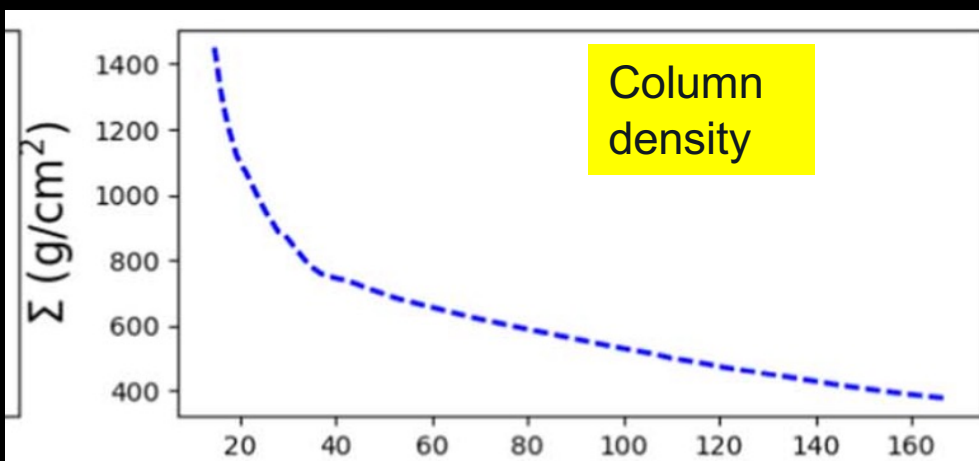
(Añez-López+2020)

Sources of heating:

- Irradiation by the stellar luminosity
- Irradiation by the accretion shock (disk surface)
- Viscous dissipation (disk mid-plane)



\dot{M} ($M_{\odot} \text{ yr}^{-1}$)	7×10^{-5}
M_{disk} (M_{\odot})	5
R_{disk} (au)	168
H_{100} (au)	7
R_{in} (au)	14
i ($^{\circ}$)	49
a_{max} (μm)	500
α	0.1

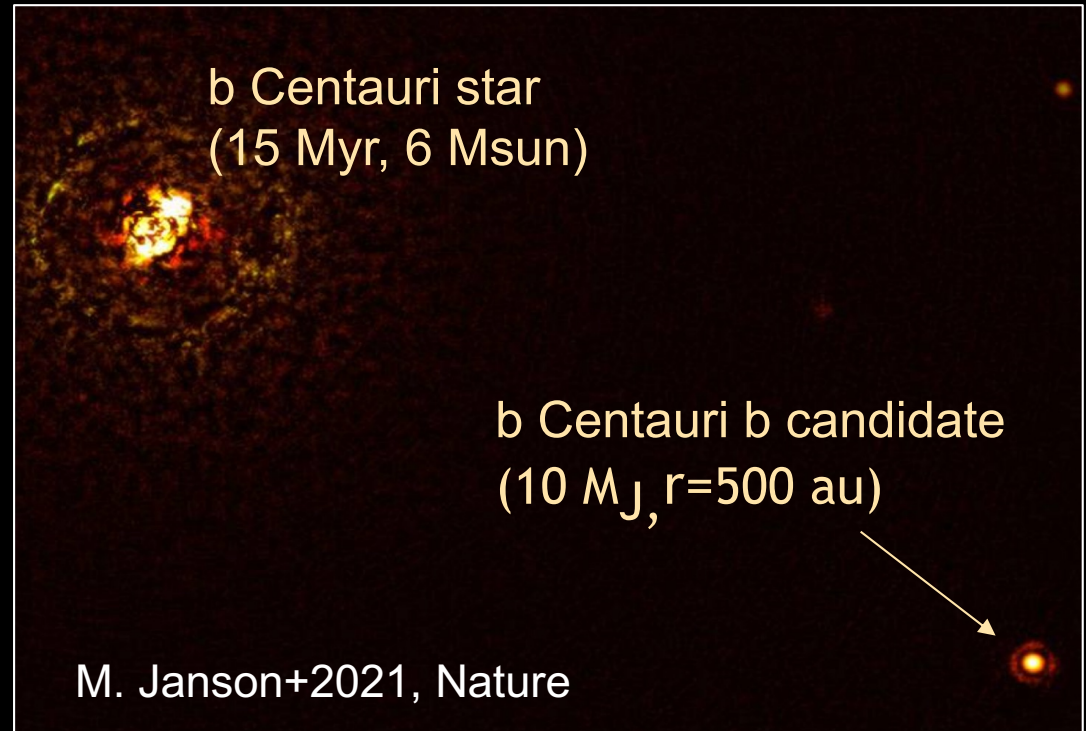
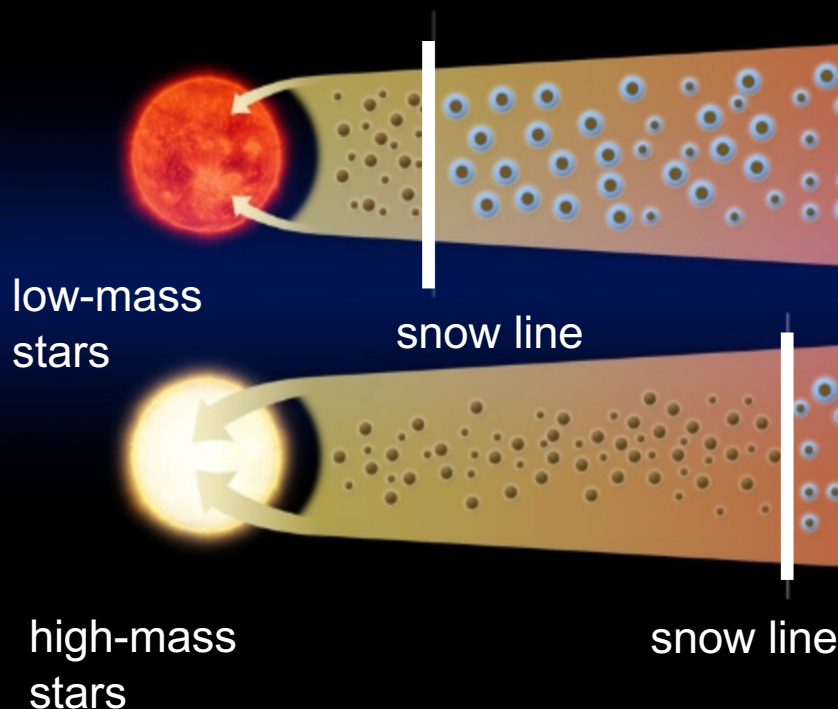


Implications of the HH80 disk modeling for planet formation

Very hot and dense disk: $T > 120 \text{ K}$ ($T < 30\text{-}40 \text{ K}$ for low-mass stars).

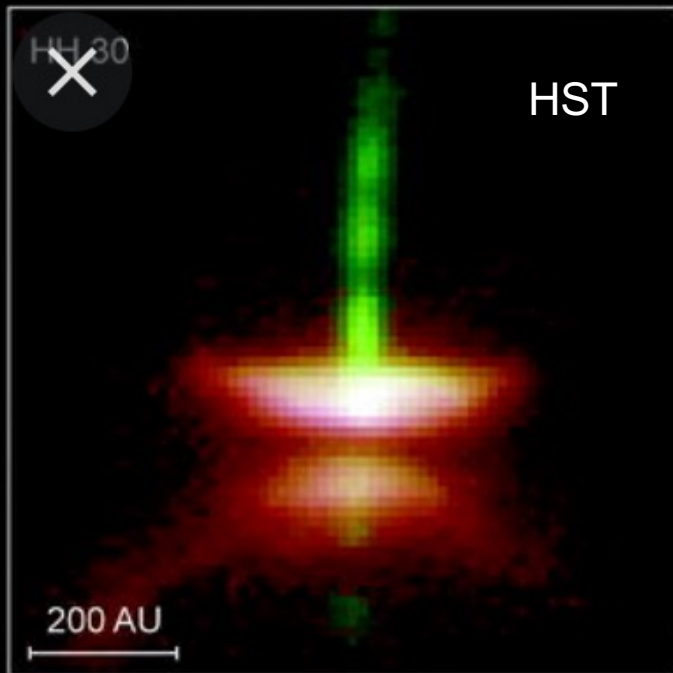
- The volatile molecules (CO , CH_3OH , H_2O) will be sublimated in the disk, due to the high temperatures (snow line outside/at the border of the disk). So, volatile molecules could only be found in the envelope surrounding the disk.
- The formation of giant gas planets would be difficult. Only the formation of rocky planets would be possible in the HH80 disk.

Discovery of a planetary system in a high-mass star

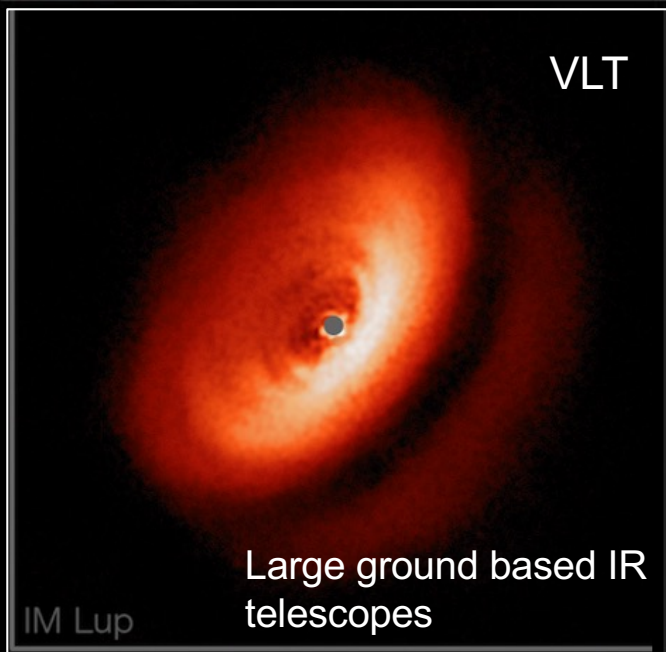
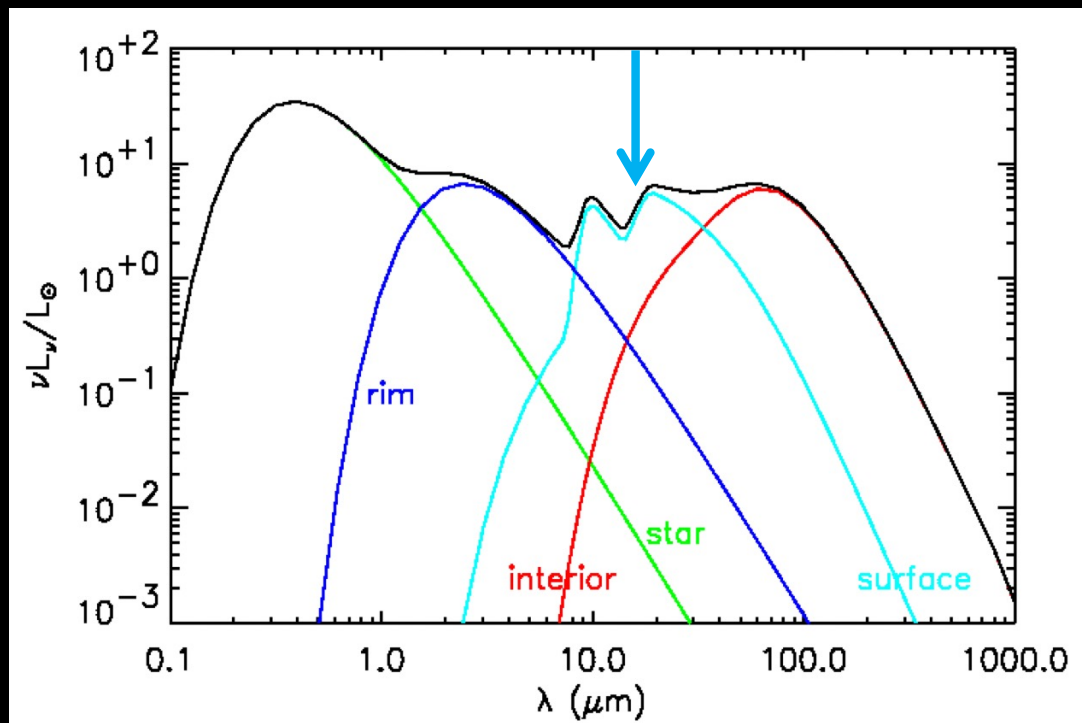


Fundamental properties of accretion disks

Flared disks



HST



VLT

Large ground based IR telescopes

H: **pressure scale height** (ratio of the sound speed $c_s = \sqrt{kT/\mu m_H}$ and the Keplerian angular speed Ω).

$$H = \frac{c_s}{\Omega} = \left(\frac{kT}{\mu m_H} \frac{r^3}{GM_*} \right)^{1/2}$$

For a typical temperature profile of $T \propto r^{-1/2}$
 The ratio $H/r \propto r^{1/4}$ increases with r .
 Disks thus have a **flaring geometric shape**



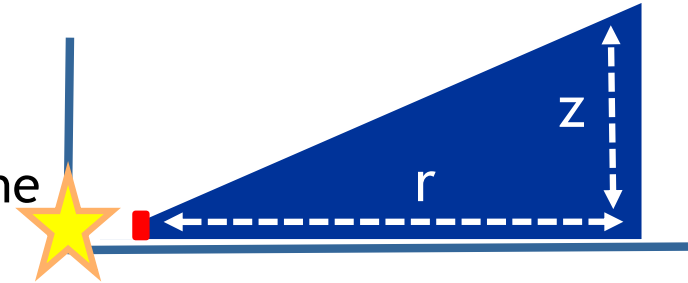
Hartmann 1987

Physical Structure (disk density)

Volumetric density: assuming that disk is in vertical hydrostatic equilibrium:

$$\frac{\partial P}{\partial z} = -\rho g_z,$$

Assuming: $P = \rho k T$ and that the star dominates the gravitational potential. The latter equation becomes:



$$\frac{\partial \ln \rho}{\partial z} = - \left[\frac{\mu m_H}{kT} \frac{GM_* z}{(r^2 + z^2)^{3/2}} + \frac{\partial \ln T}{\partial z} \right],$$

For a geometrical thin disk ($z \ll r$) and with a small Temp gradient ($\partial T / \partial z = 0$), the solution of this equation is a simple Gaussian distribution:

$$\rho = \rho_0 \exp \left[-\frac{1}{2} \left(\frac{z}{H} \right)^2 \right],$$

with

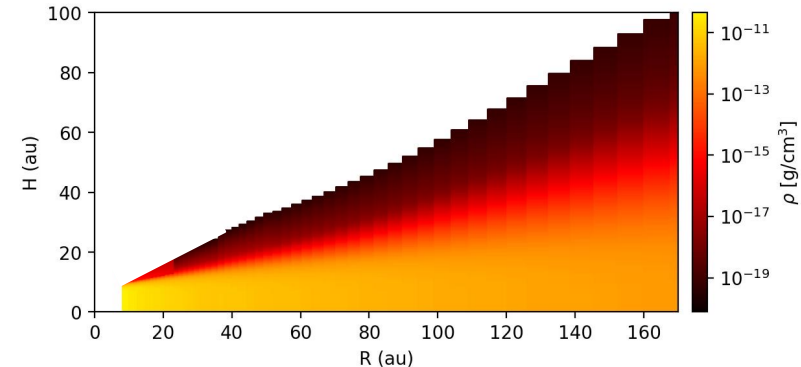
$$H = \frac{c_s}{\Omega} = \left(\frac{kT}{\mu m_H} \frac{r^3}{GM_*} \right)^{1/2} \quad \text{H: pressure scale height}$$

ρ decreases exponentially with height H .

Physical Structure (disk surface density)

The disk surface density Σ is defined as the vertical integral of the $\rho(r,z)$

$$\Sigma_{\infty} = \int_{-z_{\infty}}^{z_{\infty}} \rho dz$$



If $\rho(r,z)$ is unknown, Σ can be estimated from the equation for conservation of angular momentum in a steady α -disk:

$$\Sigma_{\infty} = \frac{\dot{M}}{3\pi \langle v_t \rangle} \left[1 - \left(\frac{R_*}{R} \right)^{1/2} \right]$$

ν is the viscosity, given in terms of the characteristic length ($\ll H$) and α turbulent velocity ($\ll c_s$). So it can be parametrized as $\nu = \alpha c_s H$, where $\alpha < 1$ (Shakura & Sunyaev 1973) thus:

$$\Sigma = (\dot{M}_{\text{dot}} / dt / 3 \pi \alpha c_s H) \left[1 - (R_*/R)^{1/2} \right]$$

Equation of Σ in term of dM_{dot}/dt , α and c_s that involved T.

Σ increases if dM_{dot}/dt increases and α and T decrease.

Disk mass, simple assumptions

$$S_{\nu} = B_{\nu}(T_d) \kappa_{\nu} M_{\text{disk}} / D^2$$

\uparrow ALMA flux
 Band 6
 \uparrow 20 K
 \uparrow 0.001 g/cm²
 \uparrow Distance (Gaia)

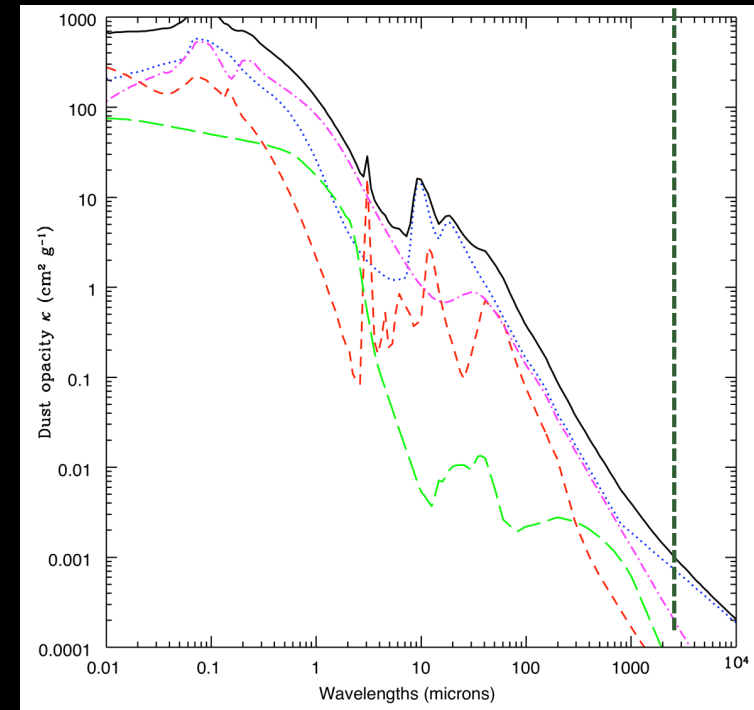
Assumptions:

- an averaged disk temperature
- an averaged dust opacity
- $\tau \ll 1$, optically thin regime
- $\kappa T \ll h\nu$; Rayleigh-Jeans regime

HD169142 disk $M_{\text{disk (dust)}} = 20 M_{\text{earth}}!$

$$M_{\text{disk(dust)}} = 400 M_{\text{earth}}$$

$$M_{\text{disk}} = 0.12 M_{\odot}$$



Minimum-Mass Solar Nebula
 MMNS $\sim 40 M_{\text{earth}}$. (Andrews 2020)

Disk mass, simple assumptions

$$S_v = B_v (T_d) \kappa_v M_{\text{disk}} / D^2$$

↑ ↑ ↑ ↑

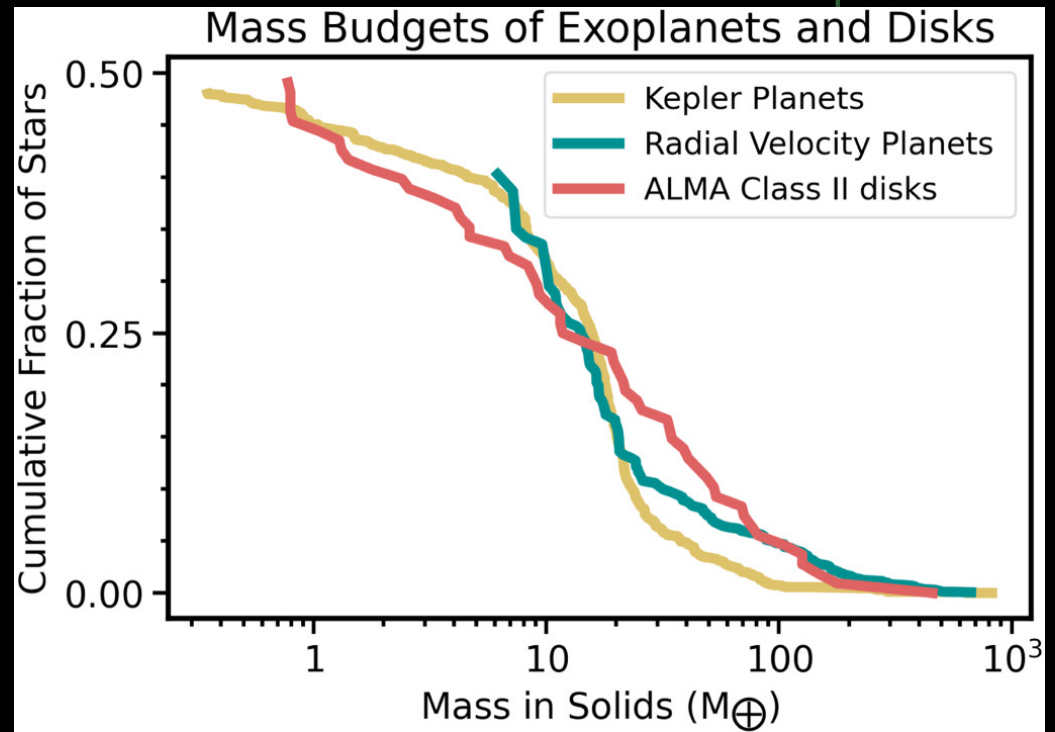
ALMA flux 20 K 0.001 g/cm² Distance (Gaia)

Band 6

Assumptions:

- an averaged disk temperature
- an averaged dust opacity
- $\tau \ll 1$, optically thin regime
- $\kappa T \ll h\nu$; Rayleigh-Jeans regime

careful with the masses derived from this approximation!



Minimum-Mass Solar Nebula
MMNS $\sim 40 M_{\text{earth}}$

Disk dust opacities

-At short wavelengths κ_ν is almost flat.

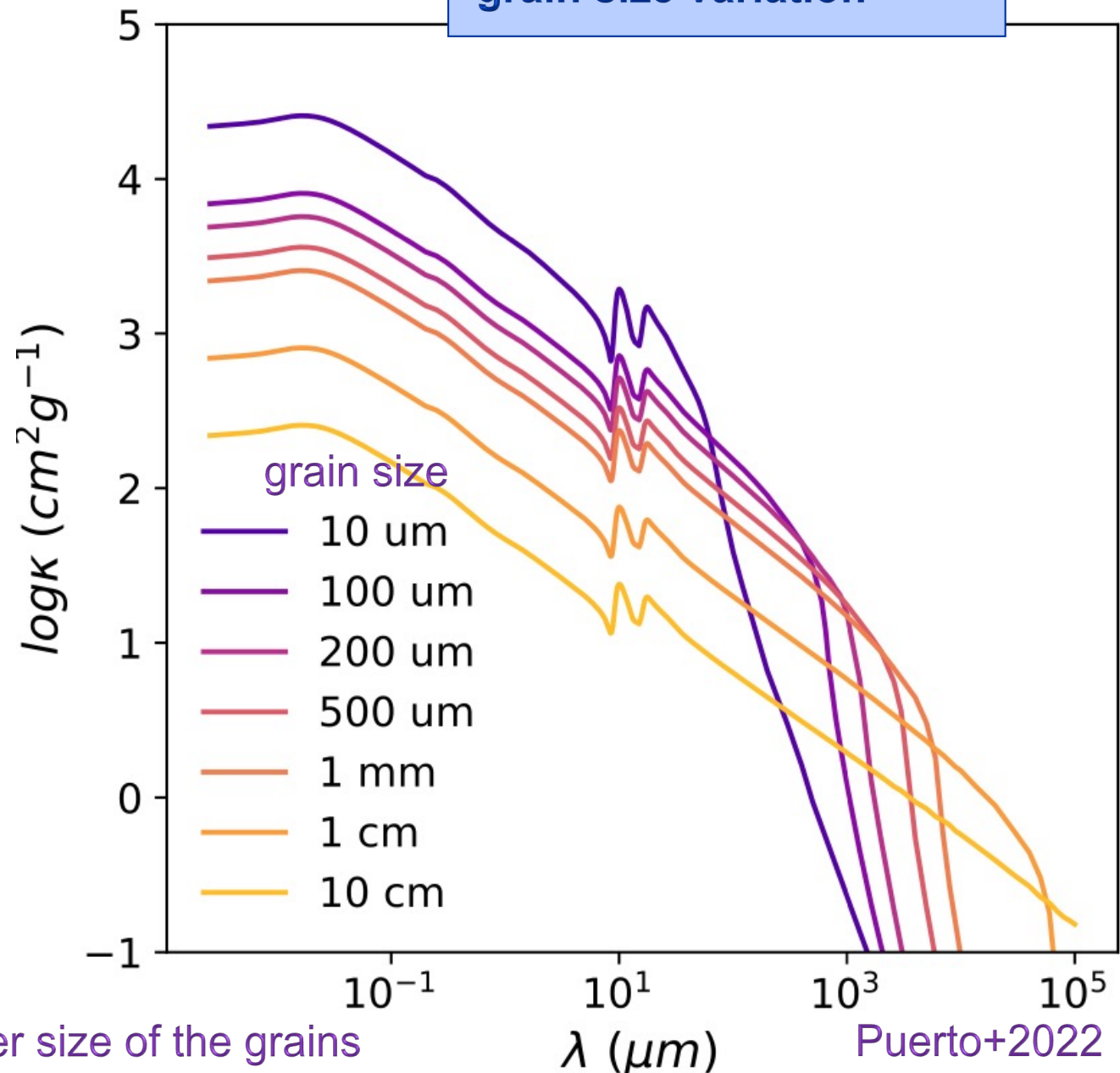
-At long wavelengths a power-law is assumed:

$$\kappa = \kappa_0 (\nu/\nu_0)^\beta = \kappa'_0 (\lambda/\lambda_0)^{-\beta}$$

small grains
(steep slope
 $\beta \sim 2$)

big grains
(shallow slope
 $\beta \sim 0$)

grain size variation



β decreases with larger size of the grains

Disk dust opacities

The dust opacity depends on:

***grain-size distribution:**

Standard grain-size
distribution

$n(a) \propto a^{-3.5}$, with

$a=0.05\mu\text{m} - 1\text{mm}$

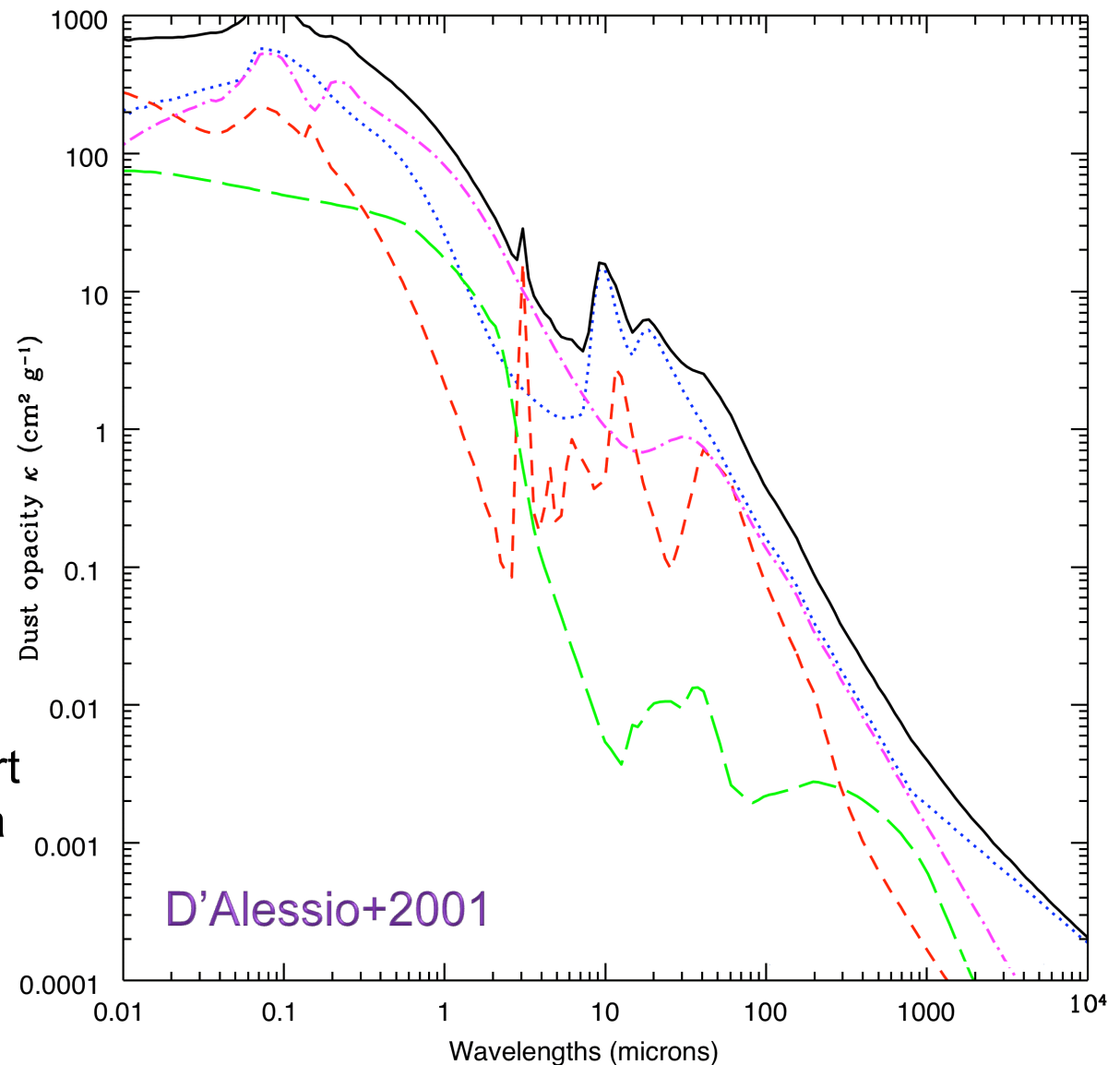
***Shapes of the grains** (Mie
Spheres)

***Composition**

***Abundance**

Since the gas is the dominant
component we need to convert
the dust into gas. To do that, a
standard dust-to-gas mass
ratio= 1/100 is assumed.

Absorption coefficient for a mixture of
grains of different **compounds**



How could the dust opacity exponent (β) be inferred from the observations?

(Sub) millimeters Spectral index

$$F_{\nu_1} = F_{\nu_2} (\nu_1/\nu_2)^\alpha$$

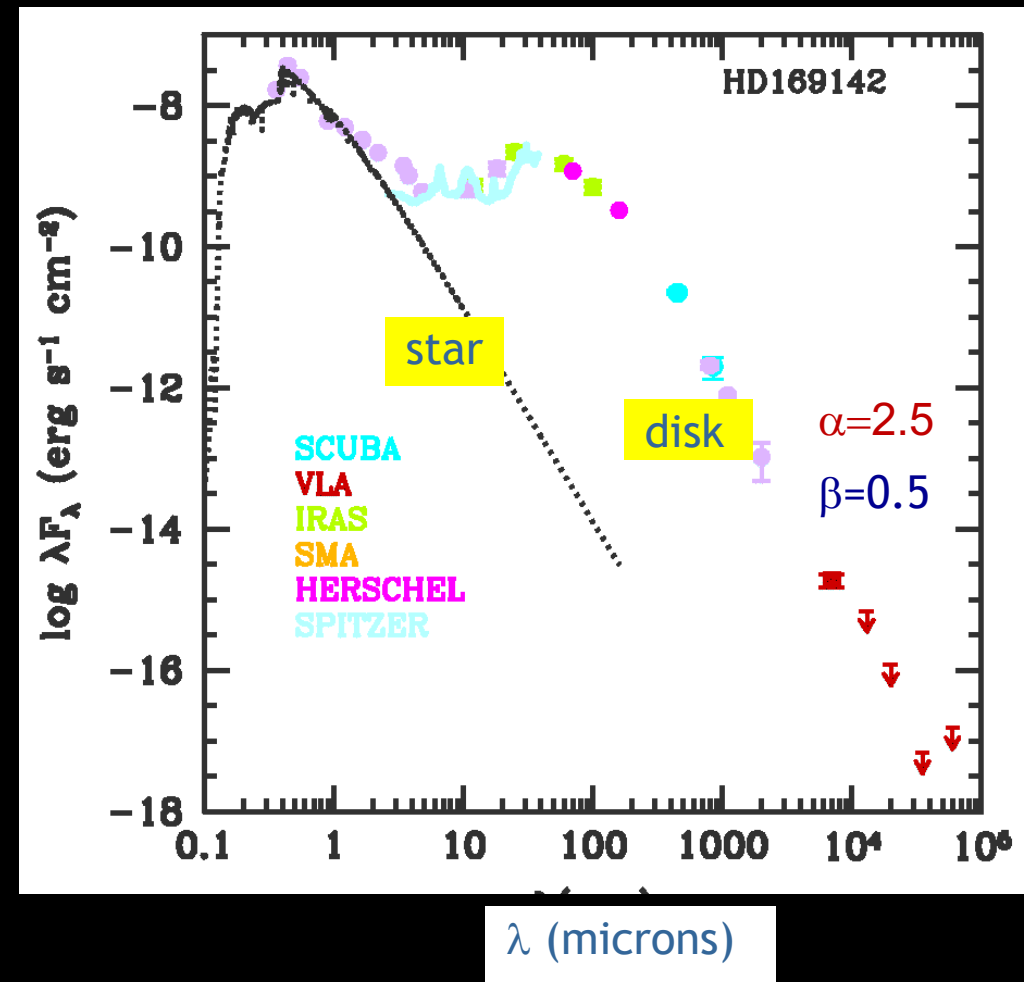
$$\alpha = \log(S_{\nu_1} - S_{\nu_2}) / \log(\nu_1 - \nu_2)$$

If optically thin emission
 $\alpha = \beta + 2$ so that
 $\alpha \sim 2, \beta \sim 0$ big grains
 $\alpha \sim 4, \beta \sim 2$ small grains

HD169142 disk:

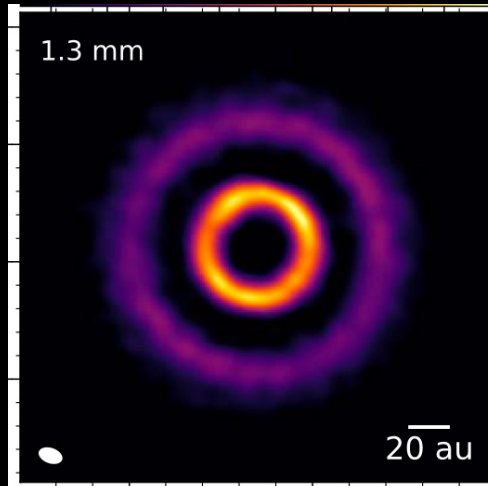
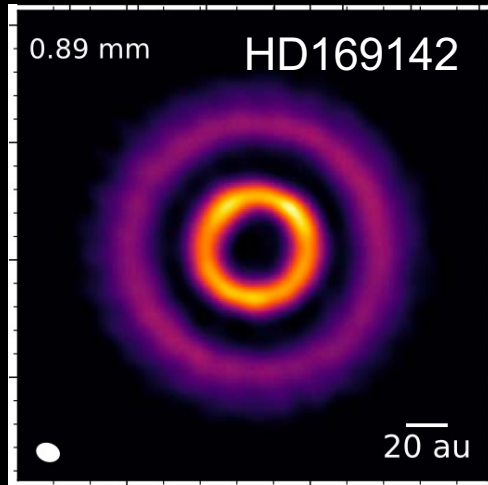
The spectrum has a spectral index, $\alpha=2.5 \Rightarrow \beta=0.5$ indicative of a shallow dependence of the opacity with the frequency which in turn implies dust grain growth.

The presence of big grains is a prerequisite for planet formation.

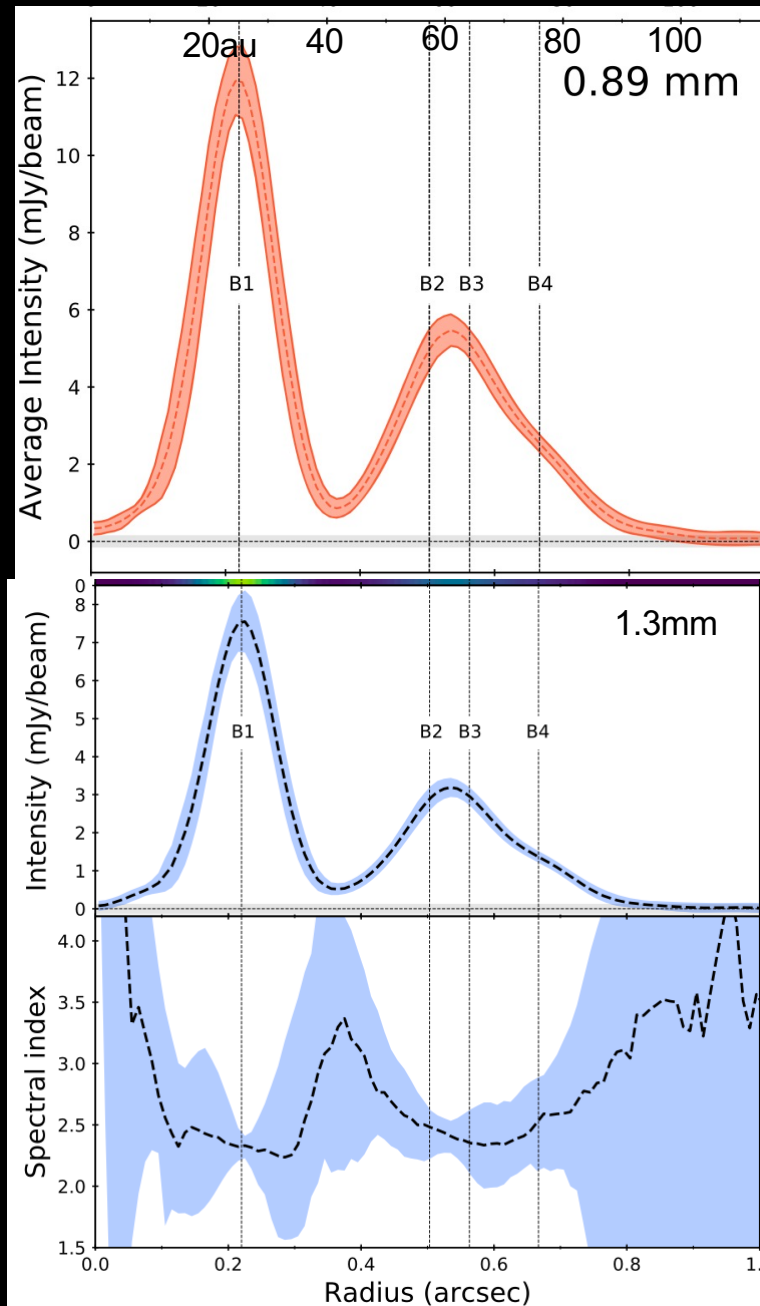


Dent + 2006, Osorio + 2014

The resolved images of the disks allow us to infer radial variations of the grain size



Macías+2019



*From ~ 20 to 100 au α increases from ~ 2.4 to 3.5 , as derived from two ALMA bands. This indicates a decrease of the grain size with radius, consistently with radial migration

*There might be an accumulation of big grains in the rings.

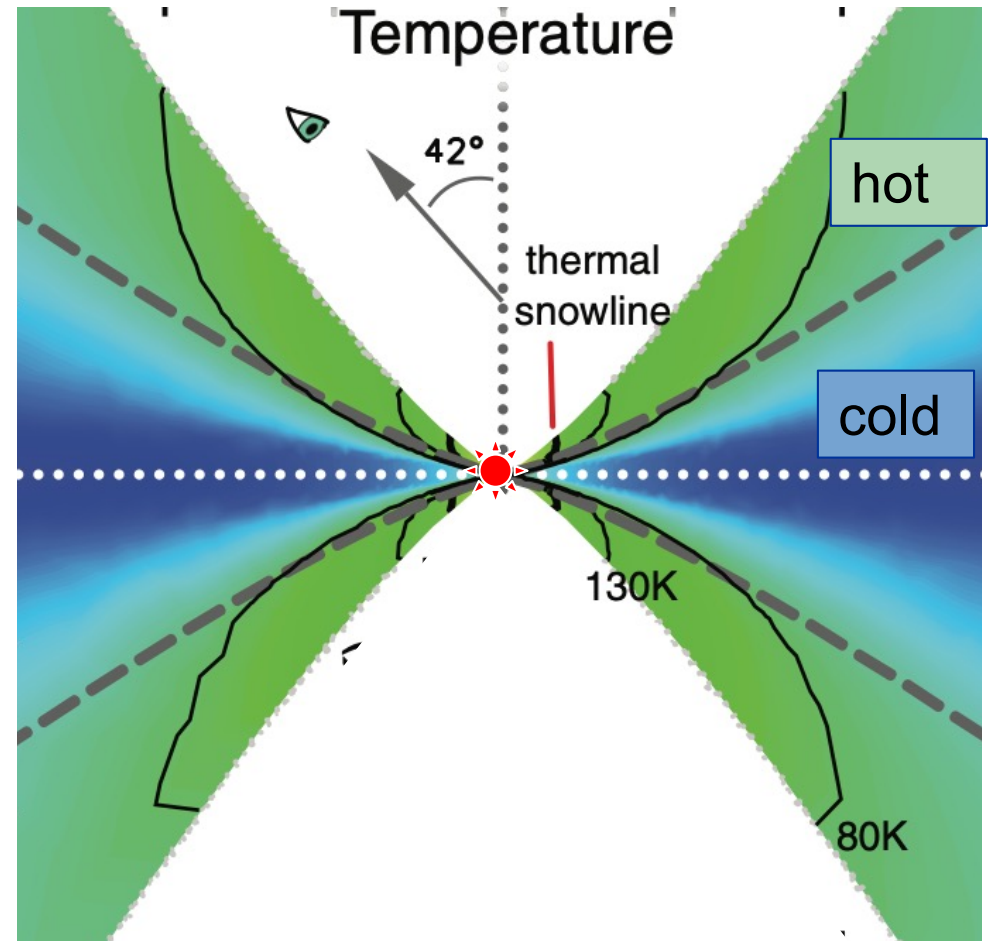
Physical Structure (disk temperature distribution (r,z))

Main heating mechanisms:

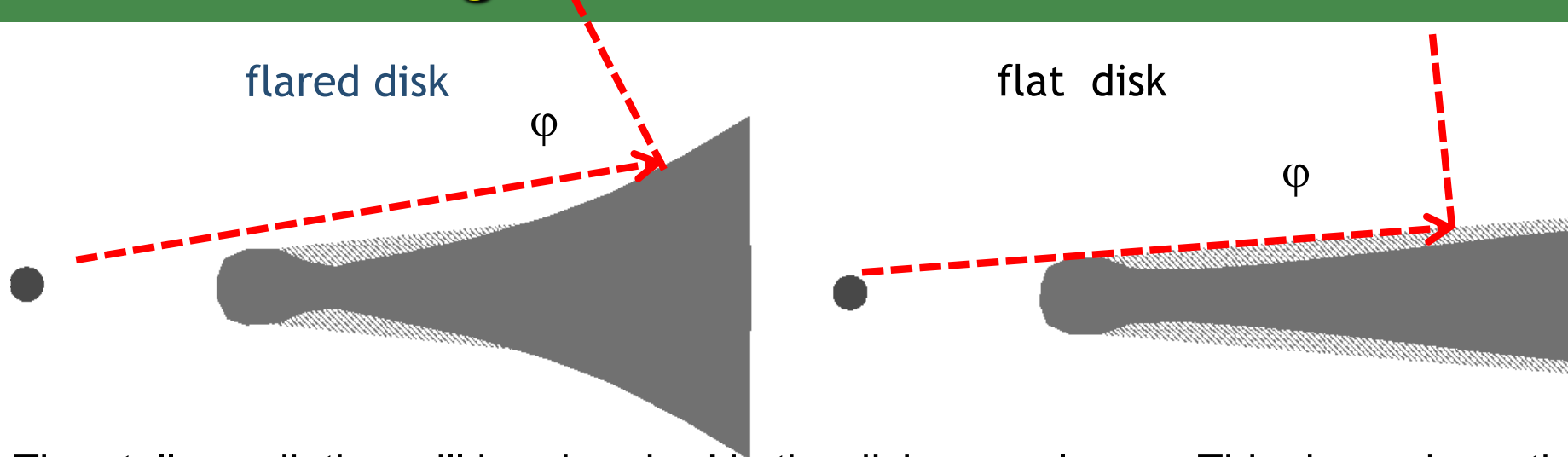
- Stellar irradiation
- Viscous dissipation

the idea behind this is that the friction (viscosity) between two neighboring disk layers (at different velocities) produces heat.

The density and temperature are crucial to determine the snowline position(s).



Heating due to the stellar irradiation



The stellar radiation will be absorbed in the disk upper layers. This depends on the geometric shape of the disk. A **flared disk** will absorb more stellar radiation than a **flat disk**. The flux of stellar radiation at distance r from the star is $L_*/2 \pi r^2$. however, the projection of this flux onto the disk surface is $\cos(\varphi) L_*/2 \pi r^2$, where φ is the grazing angle.

heating rate: $Q_+ = \varphi L_*/2 \pi r^2$

Assuming the disk's surfaces can radiate as a Planck functions with temperature T_{eff} :

cooling rate: $Q_- = 2\sigma_{\text{SB}}T_{\text{eff}}^4$

equating these two equations, T_{eff} can be obtained by:

$$T_{\text{eff}} \propto L_*^{1/4} r^{-1/2}$$

Viscous heating due to the accretion process

The heating per gram of gas is proportional to the viscosity coefficient ν and the square of shear, $\nu (r d\Omega / dr)^2$.

The total heating per cm^2 is then:

heating rate:
$$Q_+^{\text{accr}} = \Sigma \nu \left(r \frac{d\Omega_K}{dr} \right)^2 = \frac{9}{4} \Sigma \nu \Omega_K^2$$

If we insert $dM_{\text{acc}}/dt = 3 \pi \Sigma \nu$, then we get

$$Q_+^{\text{accr}} = \frac{3}{4\pi} \dot{M} \Omega_K^2$$

The release of energy is proportional to the amount of matters that accrete.

Assuming the disk's two surfaces can radiate as a Planck functions with temperature T_{eff} :

cooling rate:
$$Q_- = 2\sigma_{\text{SB}} T_{\text{eff}}^4$$

equating these two equations, T_{eff} can be obtained by:

$$T_{\text{eff}} \propto \dot{M}^{1/4} r^{-3/4}$$

Physical Structure (Temperature distribution (r,z))

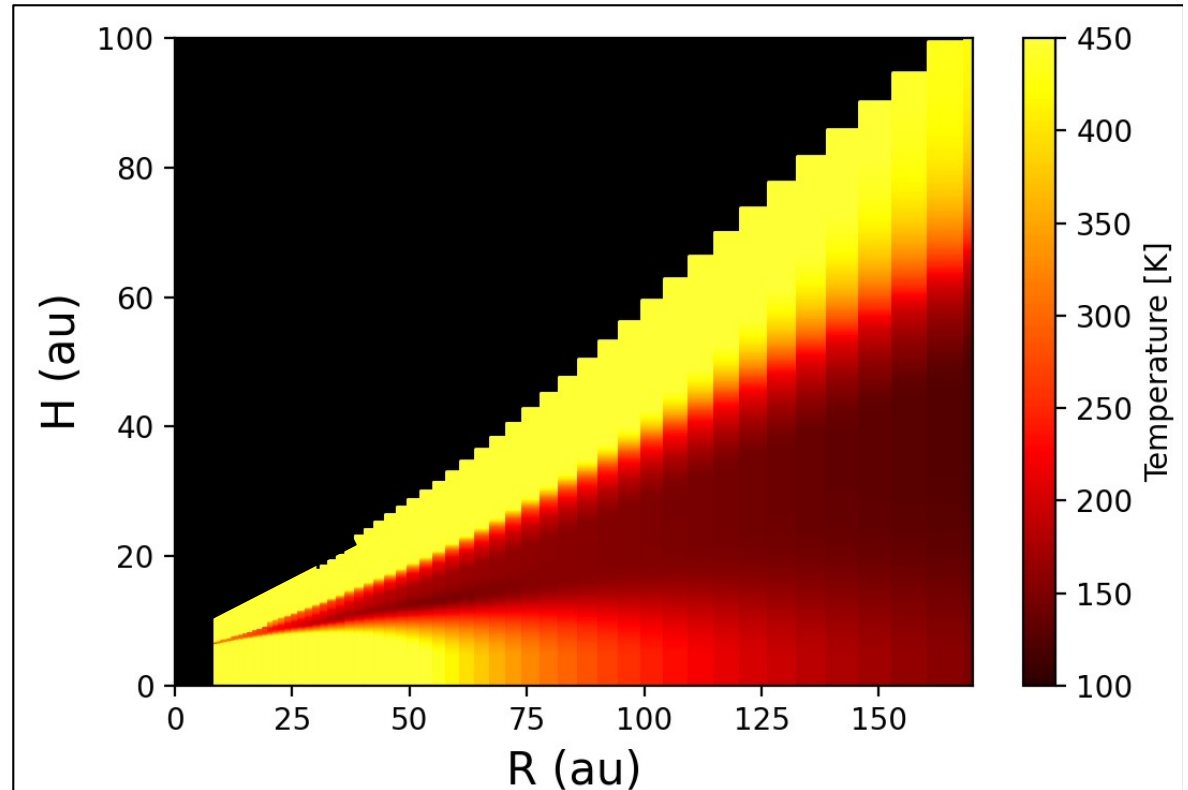
Main heating mechanisms:

- Viscous dissipation (accretion)
- Stellar irradiation

$$T_{\text{eff}} = \left\{ \frac{3}{8\pi\sigma_{\text{SB}}} \dot{M} \Omega_K^2 + \varphi \frac{L_*}{4\pi\sigma_{\text{SB}} r^2} \right\}^{1/4}$$

$$T_{\text{eff}} \propto (r^{-3} + r^{-2})^{1/4}$$

Since for viscous dissipation the T_{eff} drops steeper than for irradiation, T_{eff} is dominated by irradiation for large r , while by accretion for small r .

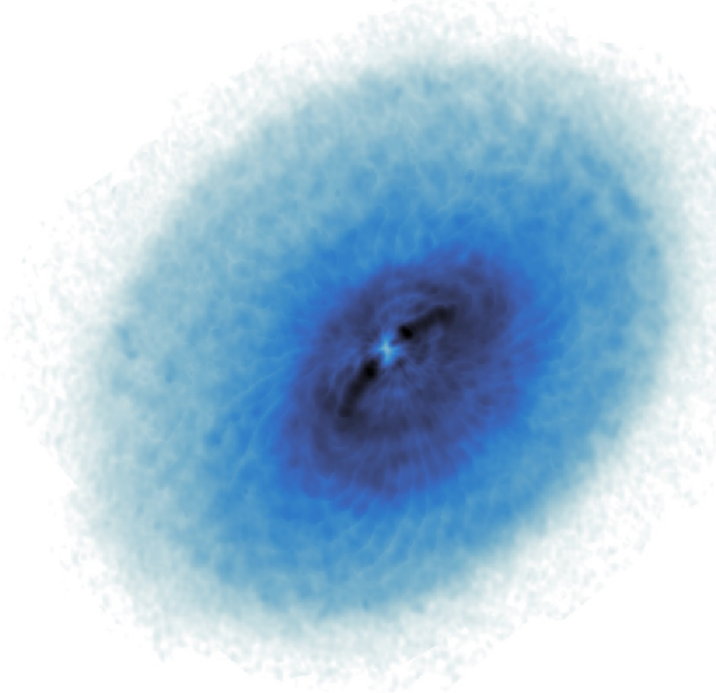


Others:

- Envelope irradiation, (class 0, MYSO)
- Accretion shock (eruptive variables, MYSO)
- Cosmic rays
- Resistive dissipation by magnetic fields (Lizano+2016)

Physical Structure (disk velocity distribution)

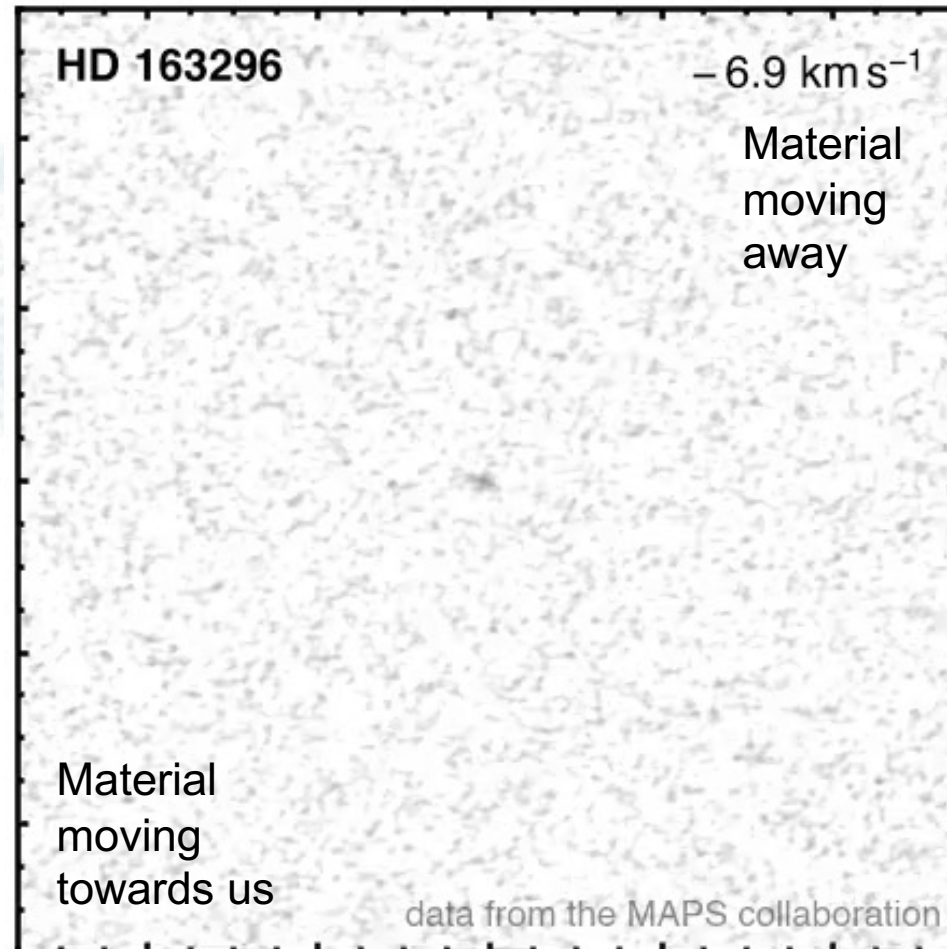
HD 163296



CO (2-1) emission
Teague+2021

Additional information can be obtained from the molecular line emission.

Channel map as a function of the V_{los}

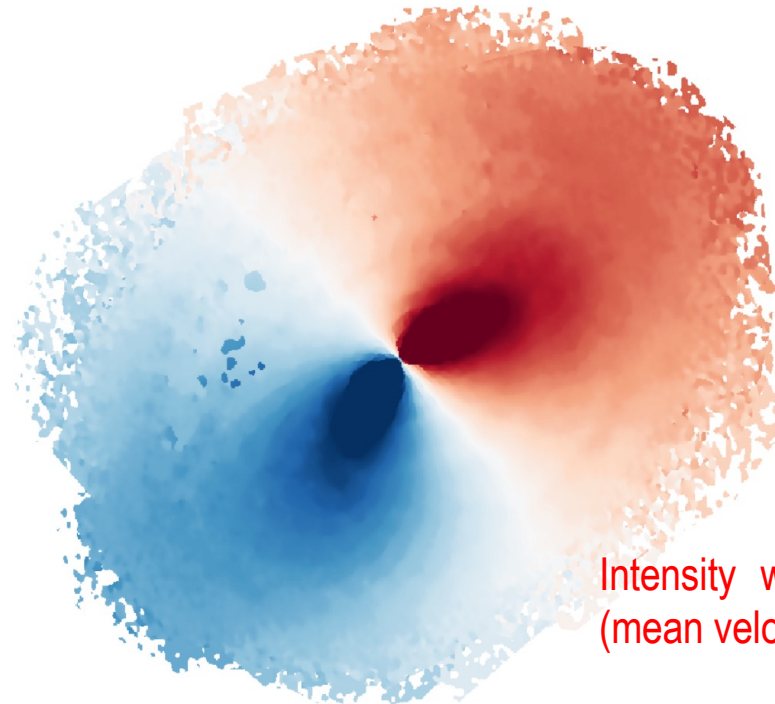
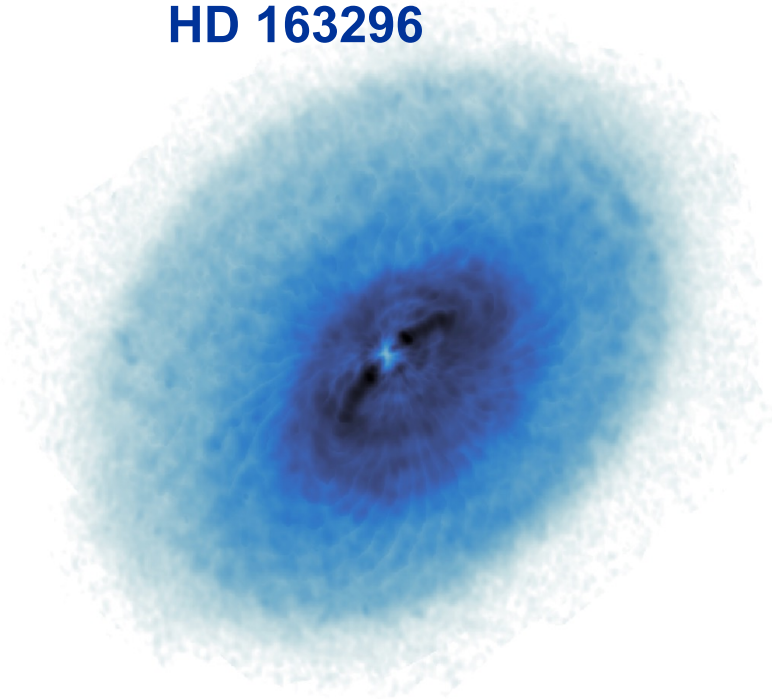


Observed butterfly pattern

This is better defined when disks have an intermediate inclination

Physical Structure (velocity distribution)

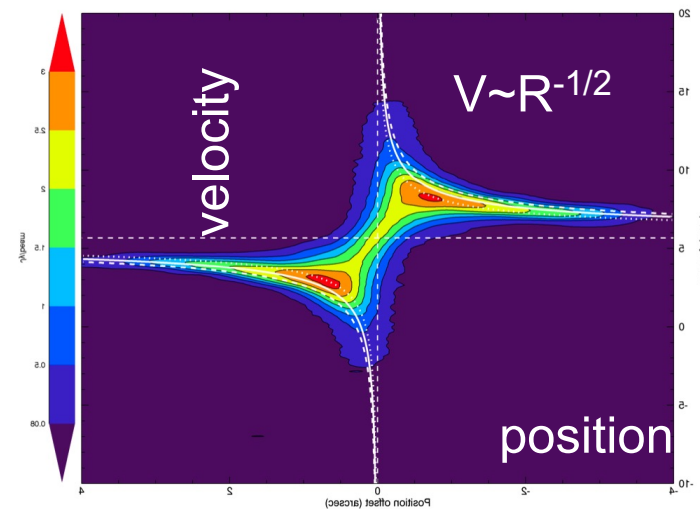
HD 163296



Intensity weighted velocity
(mean velocity)

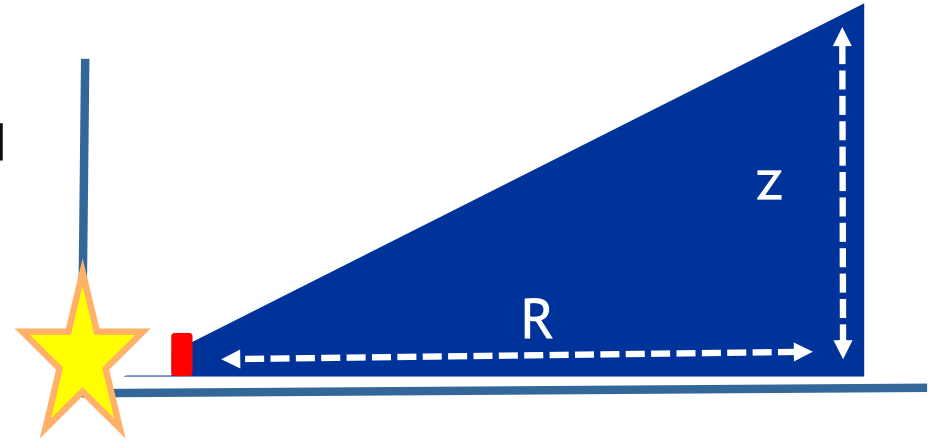
CO (2-1) emission
Teague+2021

PV diagram shows a perfect Keplerian pattern. This is the most reliable method used to derive M_*



D'Alessio Irradiated Accretion Disk models, DIAD

- Geometrically thin and axisymmetric
- Steady mass accretion rate dM_{acc}/dt and viscosity through the disk (α -prescription)
- Hydrostatic equilibrium
- Keplerian rotation $\Omega = \sqrt{GM_*/R^3}$
- Viscous and irradiated disks



The dust mixture includes two grain populations:

- $a_{\text{min}} = 0.005 \mu\text{m}$, $a_{\text{max}} \sim 1 \mu\text{m}$ (disk upper layers)
- $a_{\text{min}} = 0.005 \mu\text{m}$, $a_{\text{max}} \sim 1 \text{ mm} - 1 \text{ cm}$ (disk midplane).

Disk model catalog (D'Alessio+1998, 1999, 2001, 2006)
<https://dept.astro.lsa.umich.edu/datasets/accretion/catalog.html>

Paola's calculations have defined the state of the art in models of protoplanetary disks





END

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