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Dimensionally reduced EFTs for cosmological phase transitions

16/10/24 · LISA Spain Meeting 2024 · Institute of Space Sciences (ICE, CSIC and IEEC)

Luis Gil

Based on [2406.02667], by:

M. Chala, J. C. Criado, LG and J. López Miras

FTAE
High Energy Theory



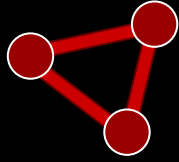
UNIVERSIDAD
DE GRANADA

A Junta
de Andalucía

The search for *new physics* in the early universe

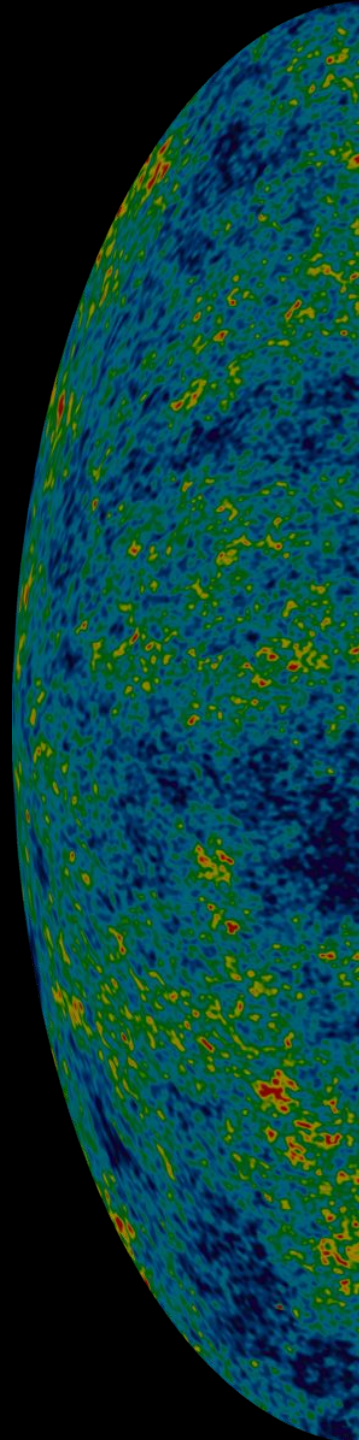


Colliders



Cosmic rays

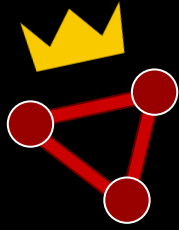
**Cosmic
Microwave
Background**



The search for *new physics* in the early universe

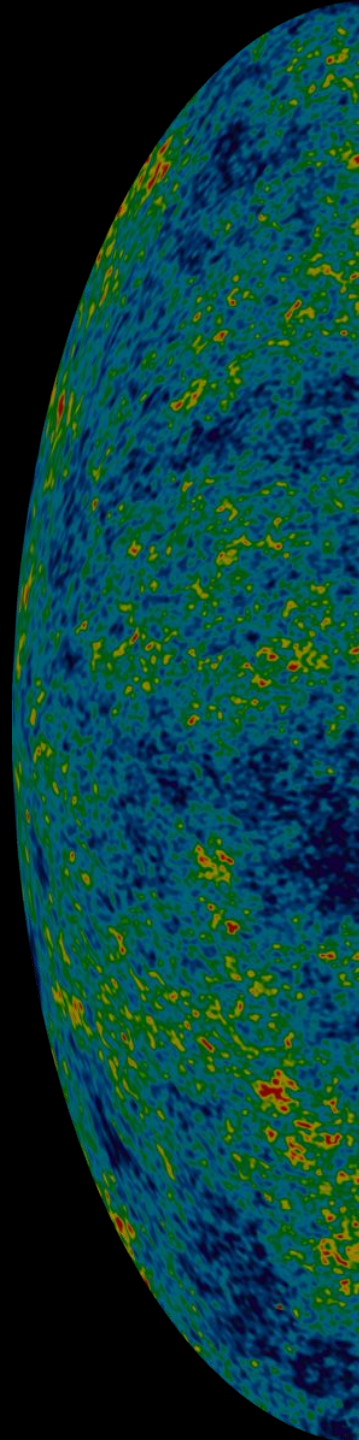


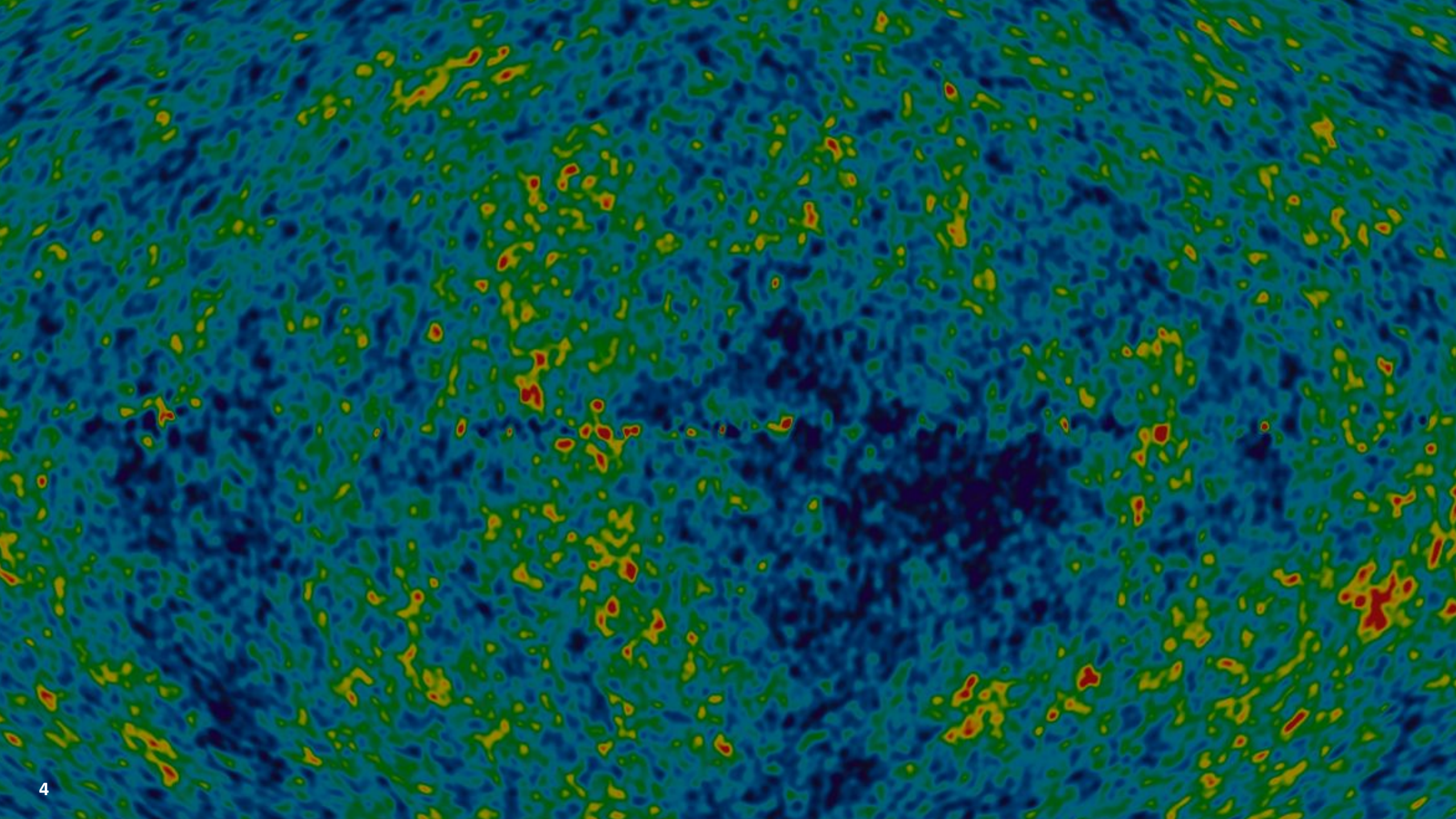
Colliders

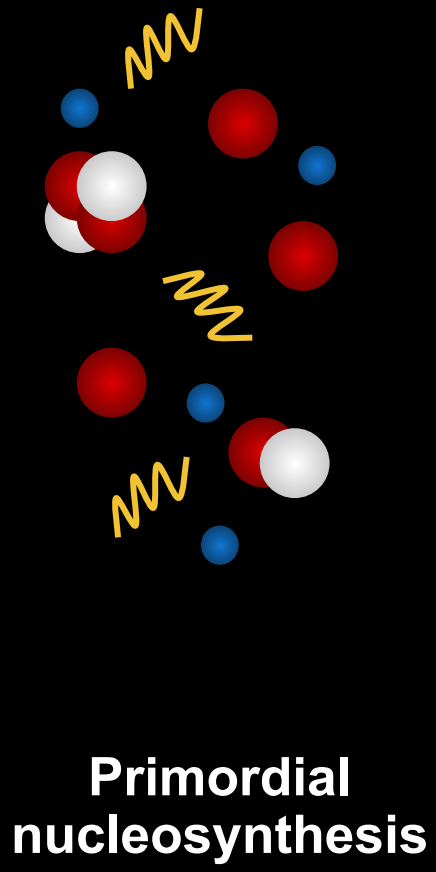
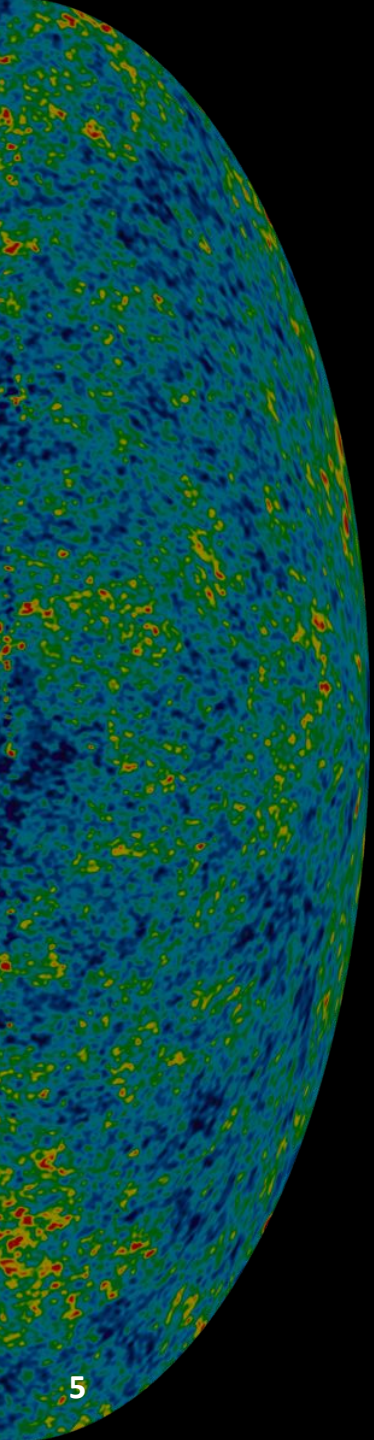


Cosmic rays

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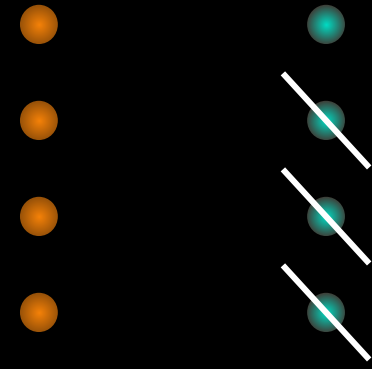




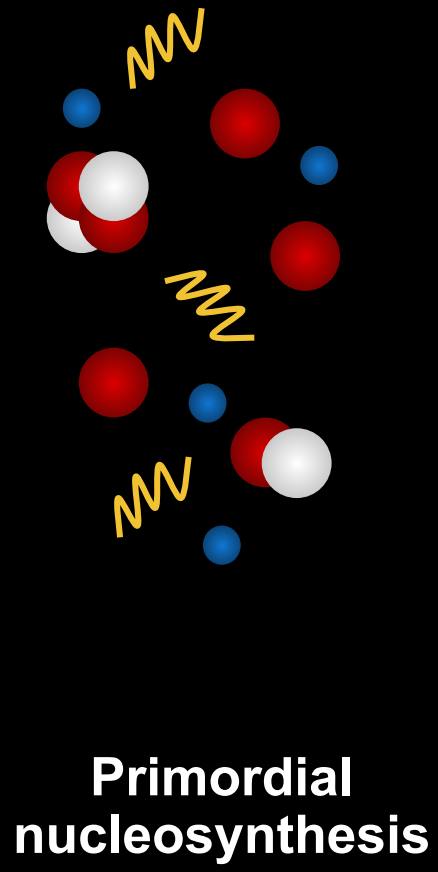
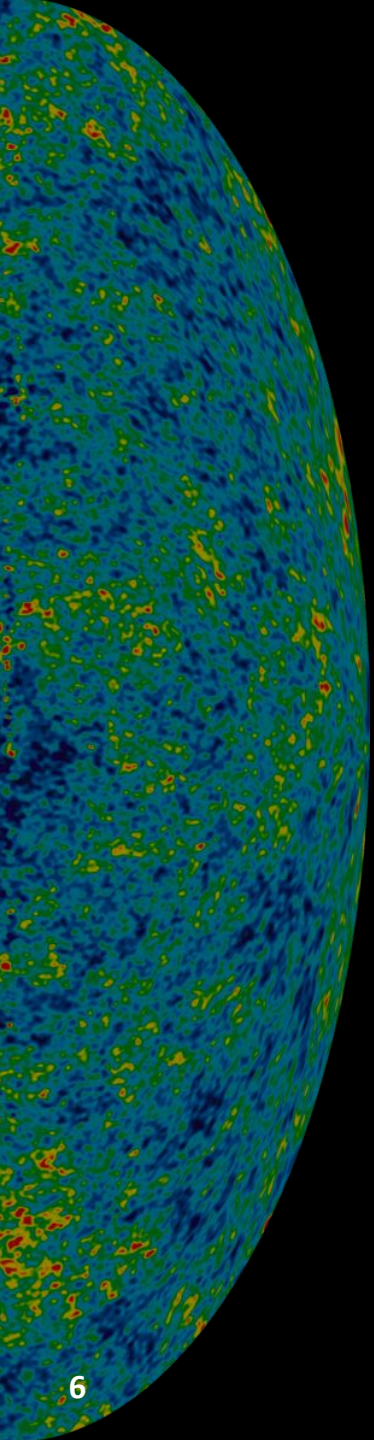


Matter

Antimatter

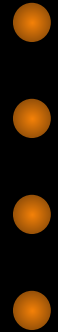


Inflation



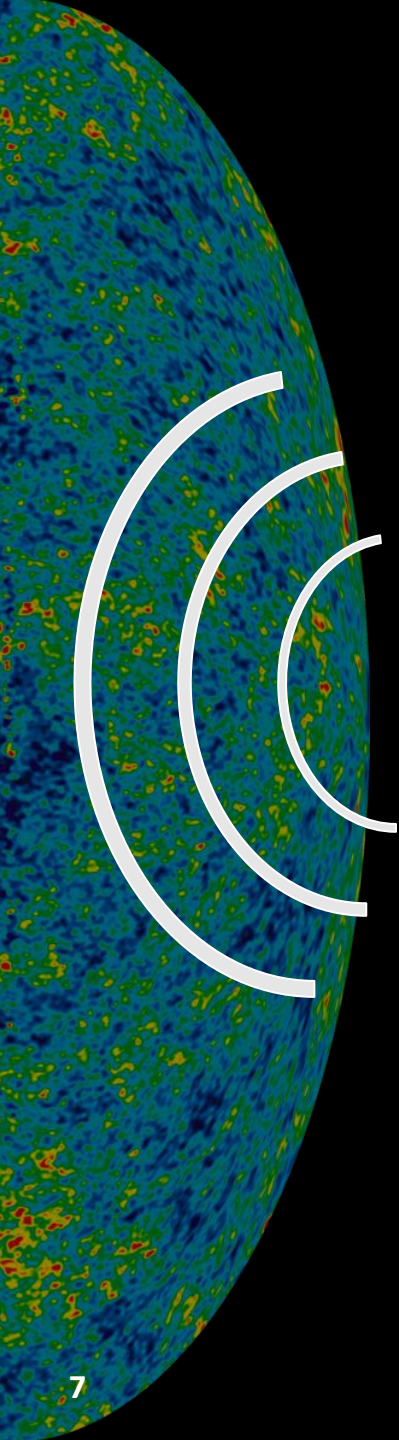
Matter

Antimatter



Electroweak (?)
Baryogenesis

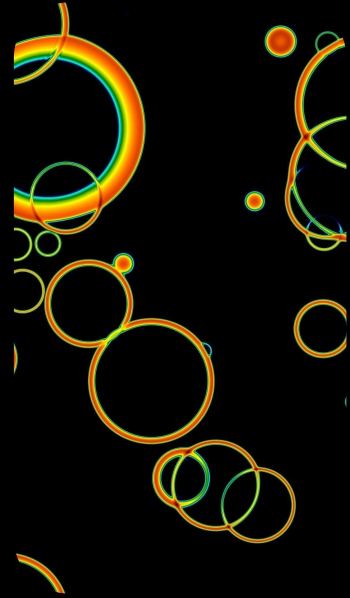
Inflation



Cosmological phase transition

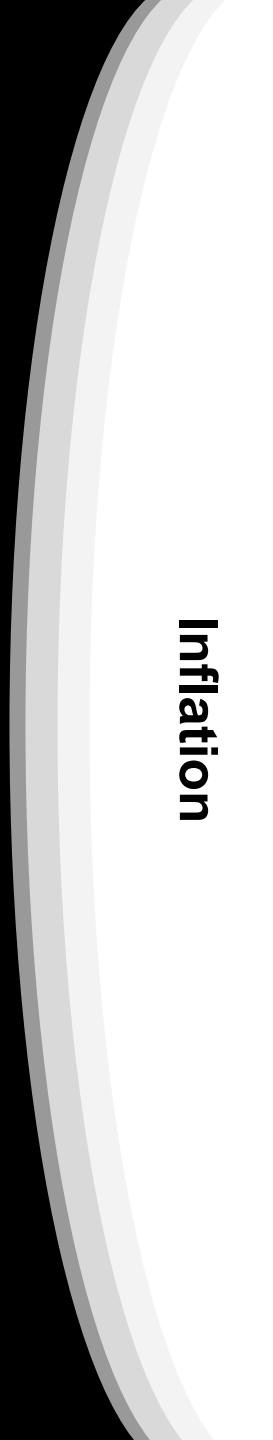
Stochastic GW background
 $h\Omega_{\text{GW}} \simeq h\Omega_{\phi} + h\Omega_{\text{sw}} + h\Omega_{\text{turb}}$

Finite temperature QFT
 $T_*, \alpha, \beta/H_*, v_w$



We need precise theoretical predictions!

Electroweak (?)
Baryogenesis



Inflation



Thermal field theory

Outline of the Matsubara formalism

- **Generating functional** ($J=0$) in QFT:

$$\mathcal{Z}[0] = \langle q' t | q 0 \rangle = \langle q' 0 | e^{-i\mathcal{H}t} | q 0 \rangle = \mathcal{N} \int \mathcal{D}q \exp(iS)$$

- **Partition function** in quantum statistical mechanics:

$$\mathcal{Z}_{\text{th}} = \text{Tr} (e^{-\beta\mathcal{H}}) = \sum_q \langle q 0 | e^{-\beta\mathcal{H}} | q 0 \rangle = \mathcal{N} \int_{q(0)=q(-i\beta)} \mathcal{D}q \exp(-S_E)$$

Field correlators

QFT at finite temperature

=

Euclidean QFT at zero temperature
with periodic time



Thermal field theory

Outline of the Matsubara formalism

- In the path integral formalism, at thermal equilibrium, fields decompose in **Matsubara modes** that live in 3D Euclidean space:

$$\phi(\tau, \mathbf{x}) \equiv T \sum_{n=-\infty}^{\infty} \phi_n(\mathbf{x}) e^{i\omega_n \tau}$$

$$\omega_n = 2\pi n T$$

$$\psi(\tau, \mathbf{x}) \equiv T \sum_{n=-\infty}^{\infty} \psi_n(\mathbf{x}) e^{i\omega'_n \tau}$$

$$\omega'_n = 2\pi \left(n + \frac{1}{2} \right) T$$

- Each mode acquires a **thermal mass** given by its **Matsubara frequency**.

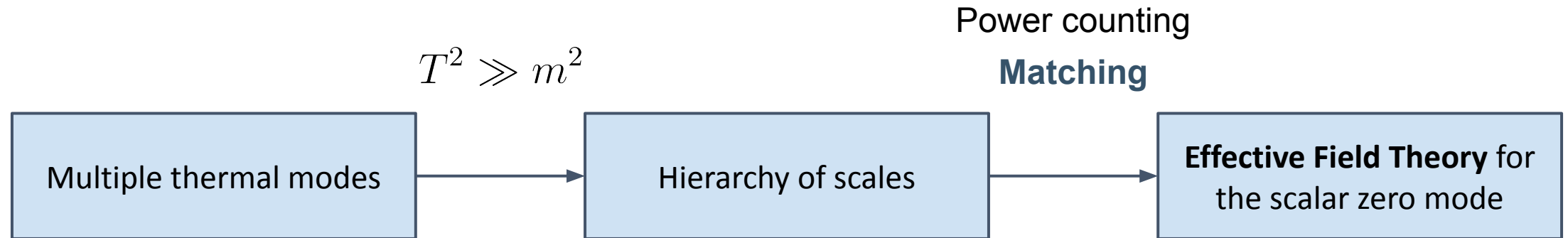
This introduces a **hierarchy of scales**.

$$M \sim \pi T^2$$



Thermal field theory

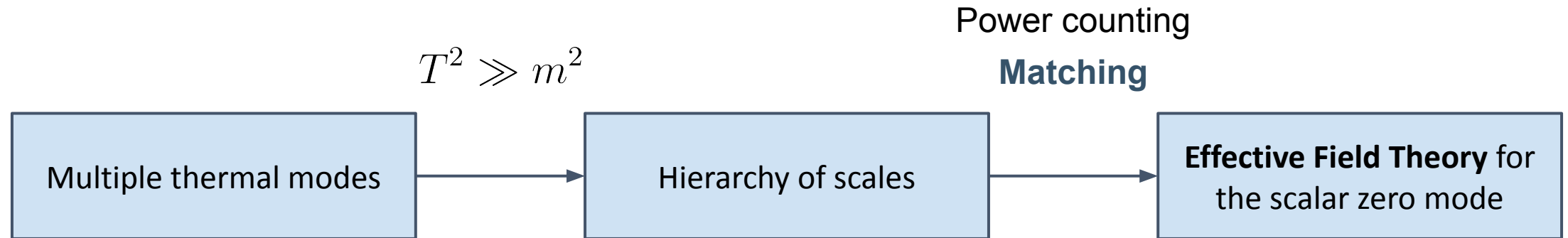
The 3D EFT approach



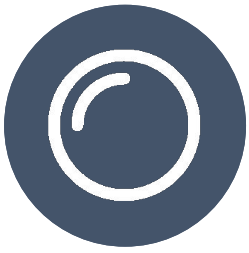


Thermal field theory

The 3D EFT approach



We have a simplified theory to study cosmological phase transitions in the high temperature limit



Bubbles and gravitational waves

Thermally induced field phase transitions

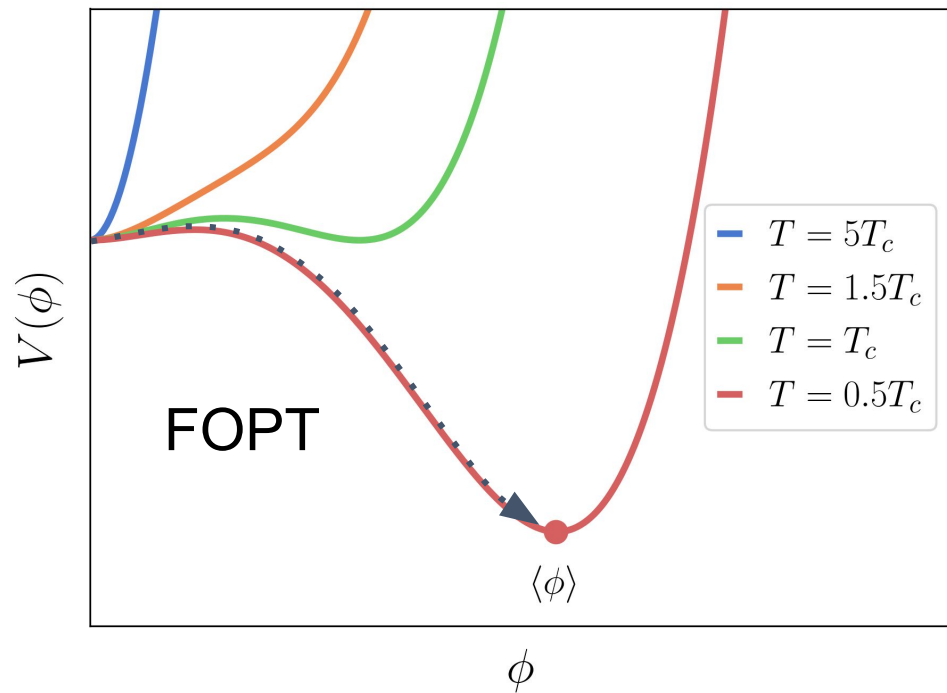


Fig. 1: Temperature evolution of scalar potential

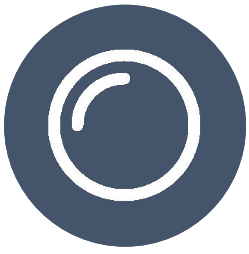
Transition rate: $\Gamma = A(T)e^{-S_E[\varphi_b](T)}$

Depends on static, non-homogeneous solutions to the (Euclidean) EoMs:

$$\left. \frac{\delta S_E}{\delta \varphi} \right|_{\varphi = \varphi_b} = 0$$

These are the so-called **bounce solutions**.

[Coleman - PhysRevD.15.2929]



Bubbles and gravitational waves

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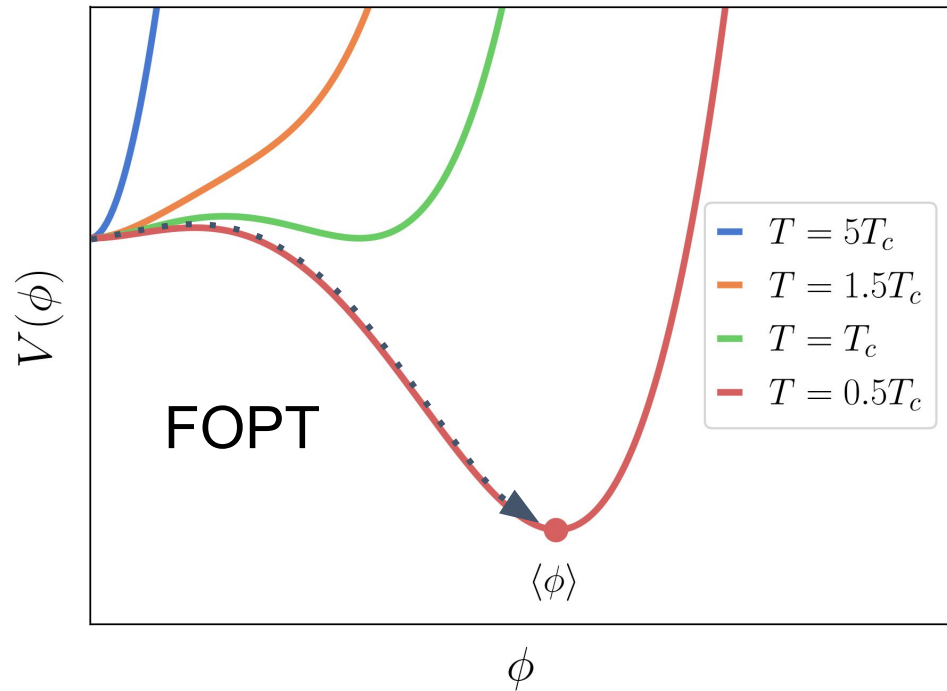


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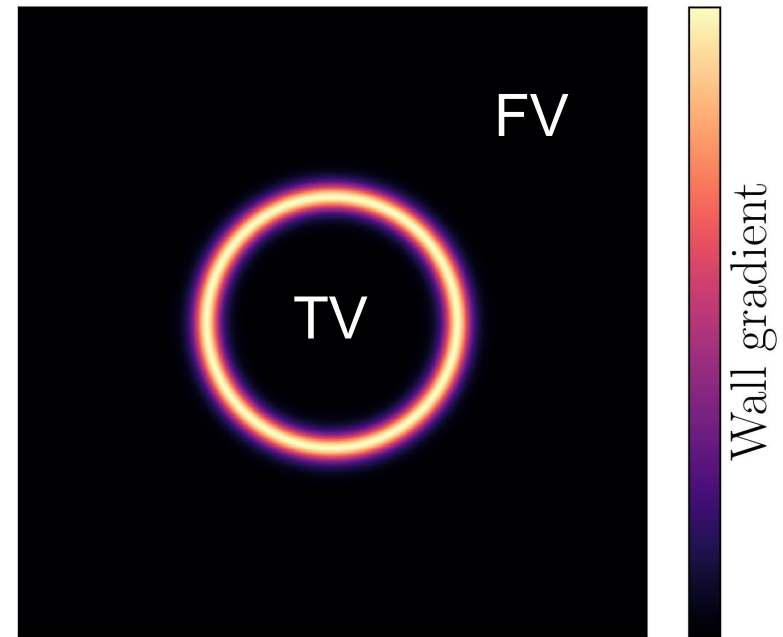
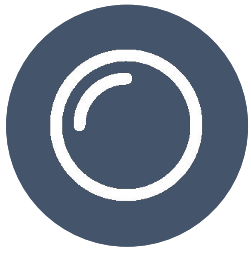


Fig. 2: $O(3)$ symmetric bounce solution



Bubbles and gravitational waves

From PTs to GWs

[Caprini *et al.* - 1512.06239]

As the transition rate grows, in a **first-order PT (FOPT)**:

- 1) Bubbles of true vacuum nucleate and expand in a hot plasma
- 2) Bubble fronts collide
- 3) Sound waves
- 4) Turbulence



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Probe for NP

How do we connect a QFT model to these GW spectra?



Bubbles and gravitational waves

From PTs to GWs

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Bubbles and gravitational waves

From PTs to GWs

How do we connect a QFT model to these GW spectra?

Nucleation temperature

$$S_3[\varphi_c] \sim 100 - 4 \log \frac{T_*}{100 \text{ GeV}}$$

Strength parameter

$$\alpha = \frac{\Delta \left(V_3(\varphi) - \frac{T}{4} \frac{d}{dT} V_3(\varphi) \right) \Big|_{T_*}}{\rho_r(T_*)}$$

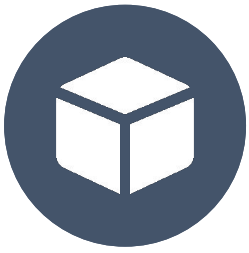
Inverse duration of PT

$$\frac{\beta}{H_*} = T_* \frac{dS_3[\varphi_c]}{dT} \Big|_{T_*}$$

Terminal bubble wall velocity

Easy to compute only in
some cases

[Lewicki *et al.* - 2111.02393]



3D EFT approach

Building our high-T EFT

UV theory in 4D Minkowski:

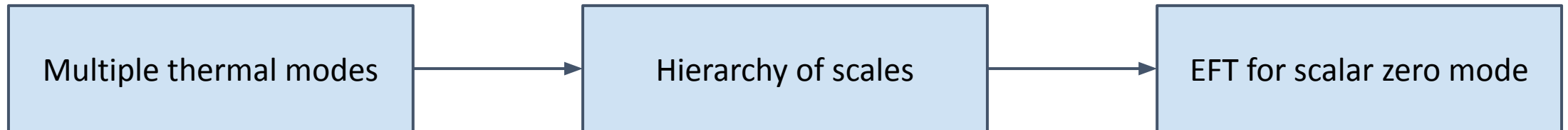
$$\mathcal{L}_4 = \frac{1}{2}(\partial_\mu\phi)^2 - \frac{1}{2}m^2\phi^2 - \kappa\phi^3 - \lambda\phi^4 + \bar{\Psi}i\not{\partial}\Psi - g\phi\bar{\Psi}\Psi \quad [\text{Gould, Xie - 2310.02308}]$$

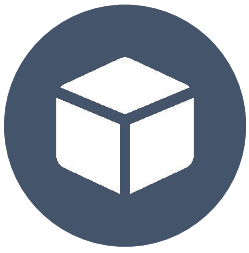
At temperature T , in Euclidean space:

$$S_0 = (-i)\frac{1}{T} \int d^3\mathbf{x} \sum_{n=-\infty}^{\infty} \left[\frac{1}{2}(\partial_i\phi_n)^2 + \frac{1}{2} \left(m^2 + (2\pi nT)^2 \right) \phi_n^2 + \bar{\psi}_n \not{\partial} \psi_n + \left(2\pi \left(n + \frac{1}{2} \right) T \right)^2 \bar{\psi}_n \psi_n \right]$$

$$T^2 \gg m^2$$

Power counting





3D EFT approach

Building our high-T EFT

[Kajantie *et al.* - 9508379]

[Criado - 1901.03501]

We determine our 3D EFT by matching correlators to the full theory in the **high-temperature limit**:

$$\mathcal{L}_3 = \frac{K_3}{2} (\partial\varphi)^2 + \frac{1}{2} m_3^2 \varphi^2 + \kappa_3 \varphi^3 + \lambda_3 \varphi^4 + \underbrace{\text{higher orders in } \left(\frac{\mathbf{m}}{\mathbf{T}}\right)}$$

Includes:

- New effective operators (**EO**)
- Corrections to operator coefficients



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How important are they?

Includes:

- New effective operators (**EO**)
- Corrections to operator coefficients



Results

PT magnitudes and effective interactions

[Chala, Criado, LG, López Miras -
2406.02667]

■ $(m^2, \kappa, \lambda)_A = (20\,000 \text{ GeV}^2, -40 \text{ GeV}, 0.01)$

■ $(m^2, \kappa, \lambda)_B = (31\,643.5 \text{ GeV}^2, -71.1 \text{ GeV}, 0.045)$

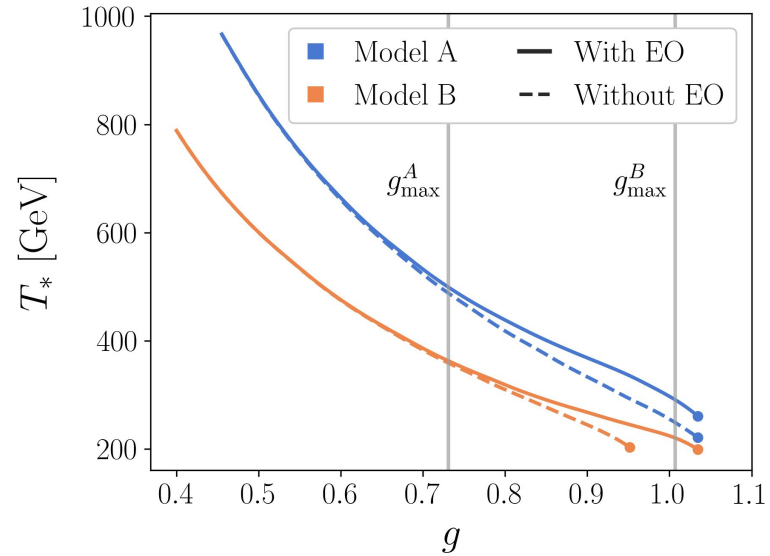


Fig. 5: PT magnitudes in two models, with and w/o effective operators



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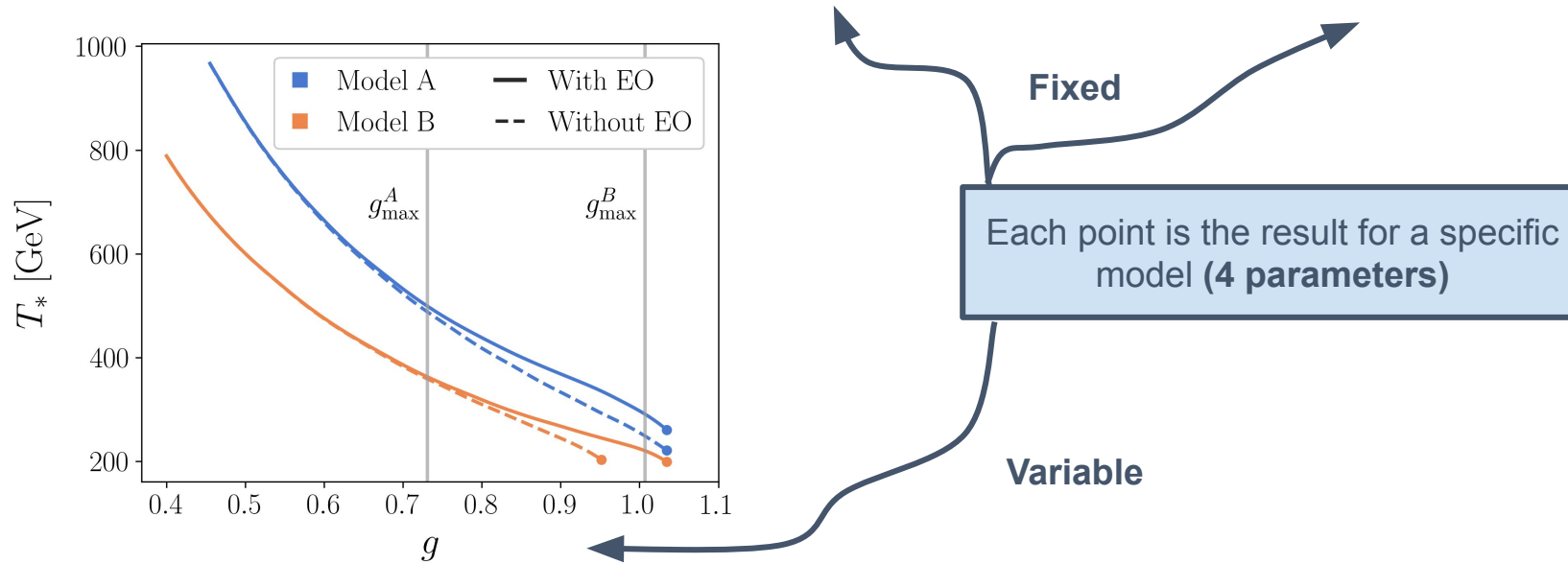


Fig. 5: PT magnitudes in two models, with and w/o effective operators



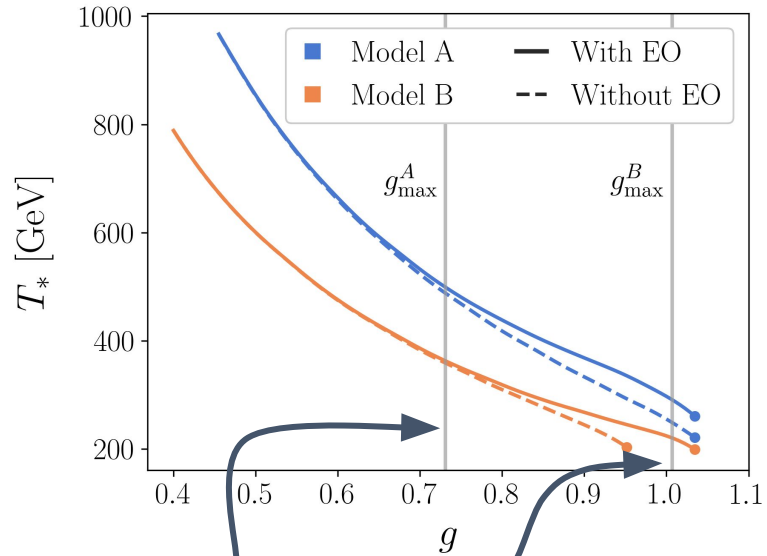
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3dEFT not valid beyond the corresponding line!

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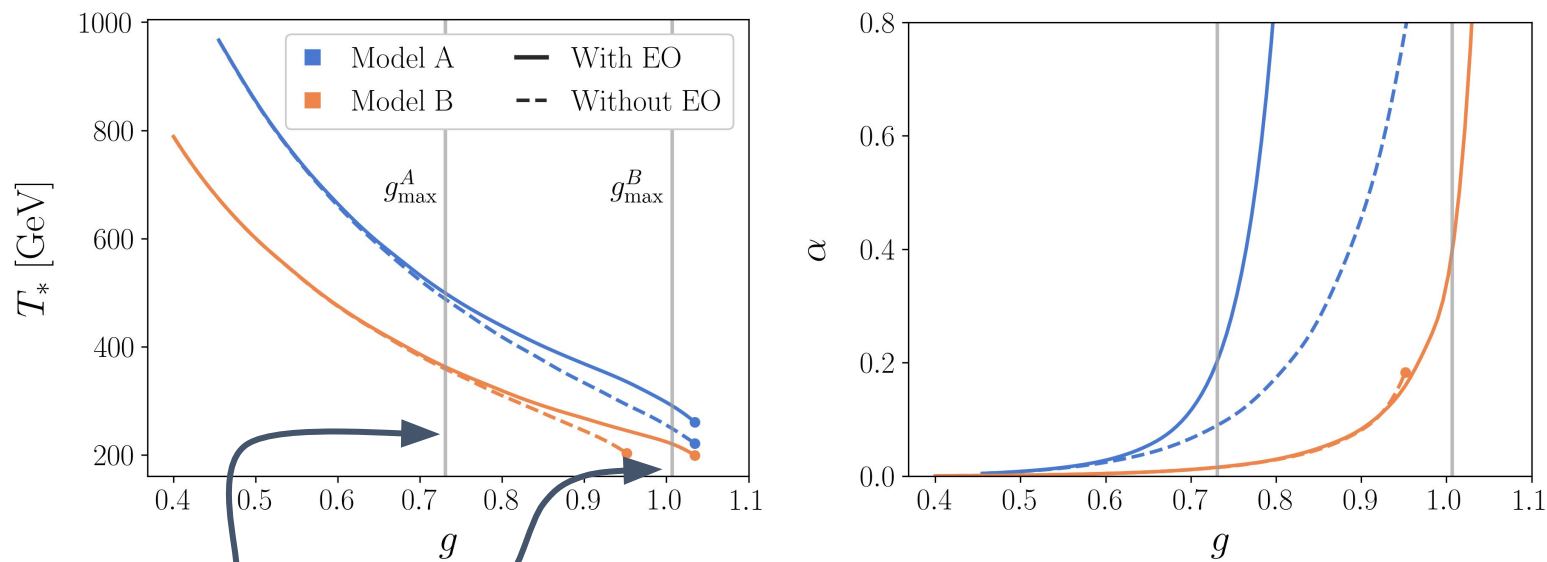
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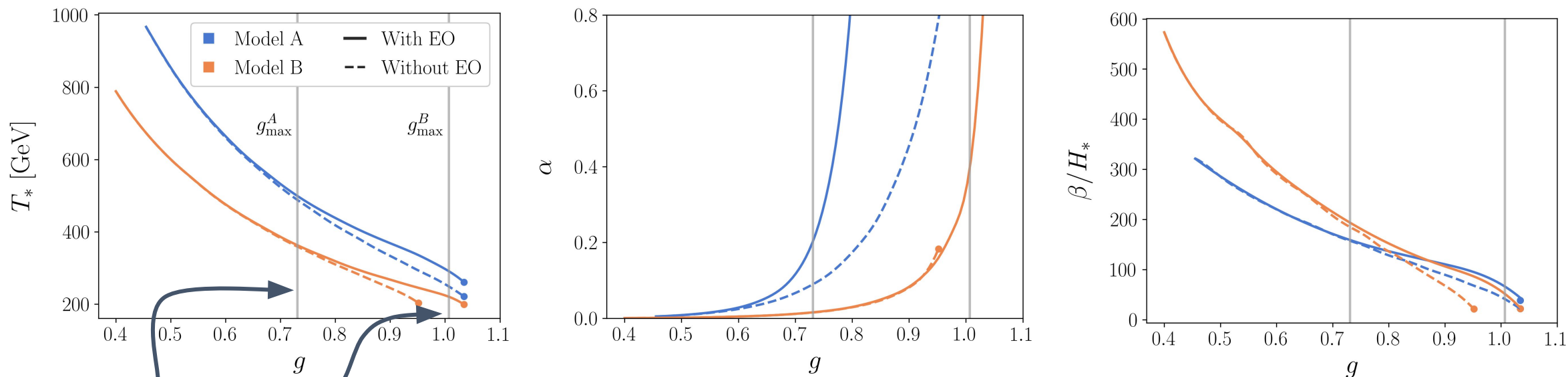
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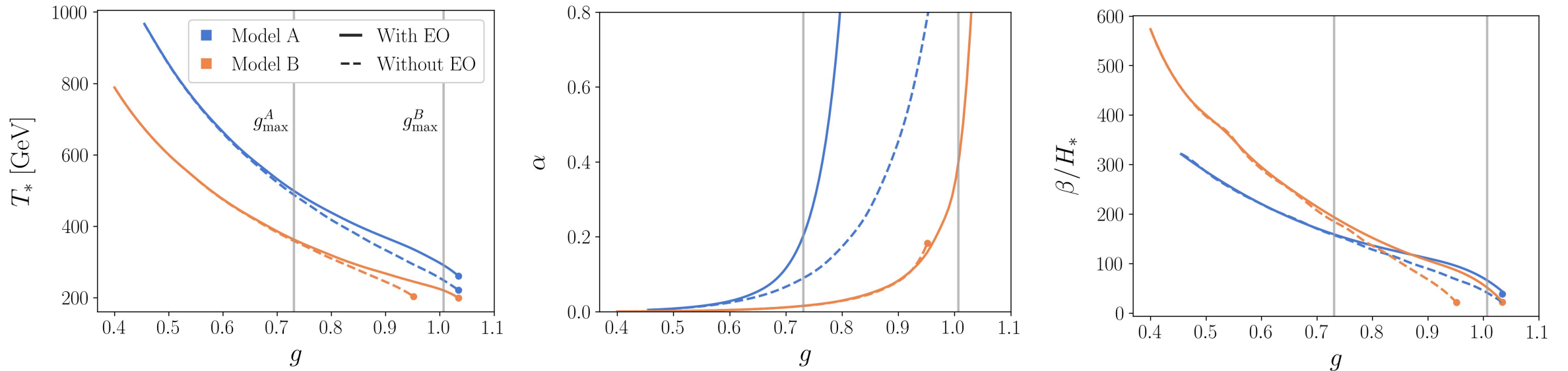


Fig. 5: PT magnitudes in two models, with and w/o effective operators

Two observations

- Eff. ops. **allow for PTs in a wider range** of values of the Yukawa
- Including eff. ops. yields **very different estimations** at large Yukawas



Results

[Chala, Criado, LG, López Miras - 2406.02667]

GW power spectra and effective interactions

■ $(m^2, \kappa, \lambda)_A = (20\,000 \text{ GeV}^2, -40 \text{ GeV}, 0.01)$

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$g \lesssim g_{\text{max}}$

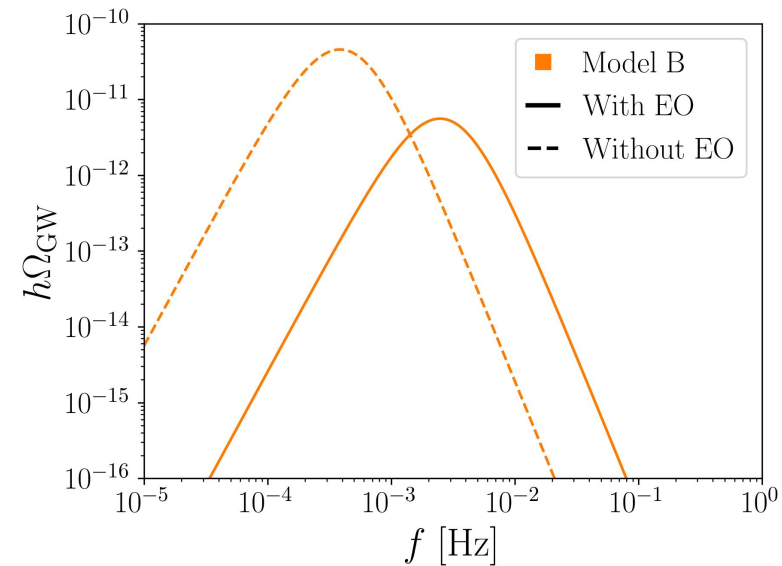
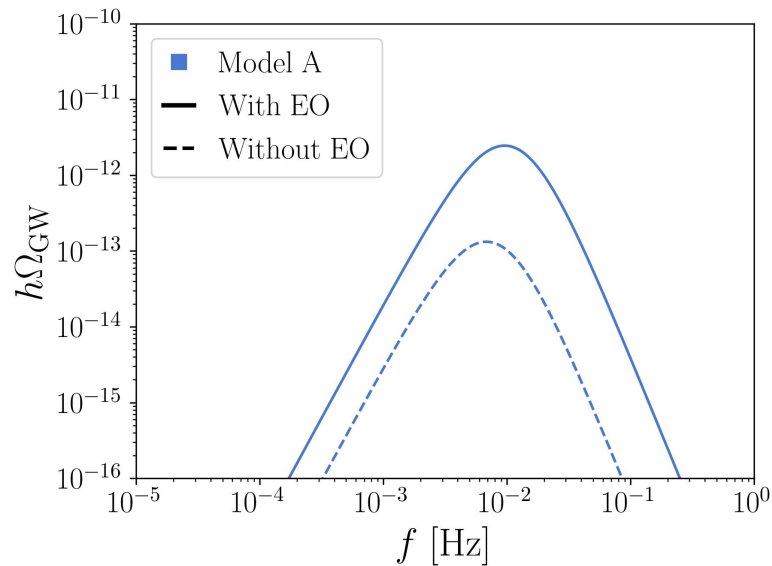


Fig. 6: GW power spectra in two models, with and w/o effective operators

(*) Generated with **PTPlot**
[Caprini *et al.* - 1910.13125]



Results

[Chala, Criado, LG, López Miras - 2406.02667]

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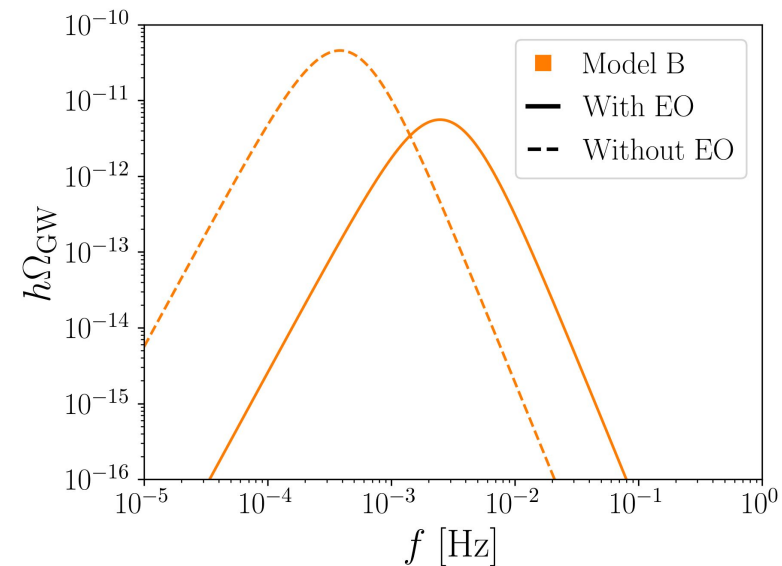
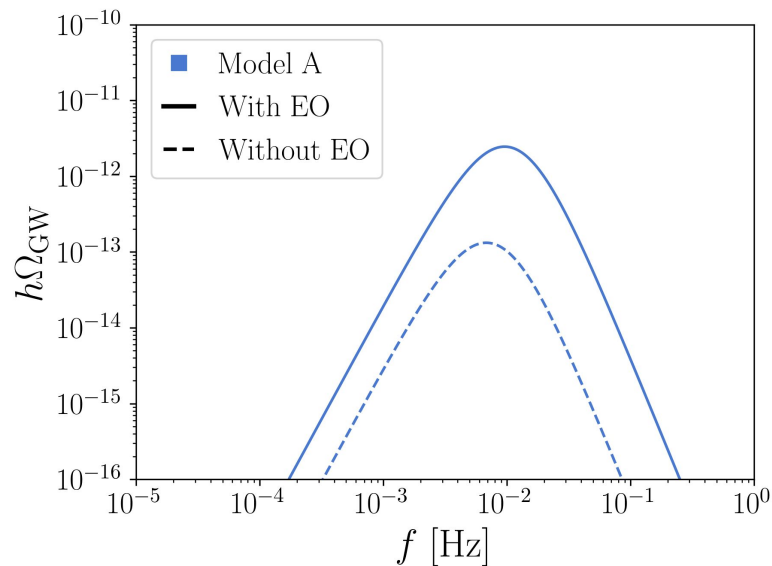


Fig. 6: GW power spectra in two models, with and w/o effective operators

Including eff. ops. can change the peak amplitude and frequency by **one order of magnitude**



Take home messages

We have learned

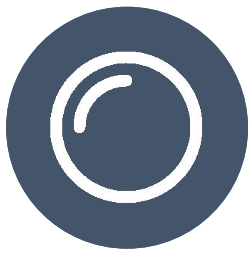
- Strong FOPTs occur in regions of parameter space **close the limit of validity** of the 3D EFT
- Higher-order-operator corrections in the 3dEFT **are thus relevant in the estimation of PT-related magnitudes** (order of magnitude differences in GW spectra)

Future work

- Application to well-motivated BSM models
- Full order corrections to assess dependence on unphysical parameters

Thank you for your attention!

¡Gracias por vuestra atención!



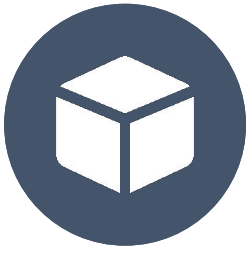
Bubbles and gravitational waves

What is the state of the art?

**This list is NOT
exhaustive!**

- Several studies of FOPTs in BSM extensions and in the SMEFT:
 - N⁴LO Higgs potential at finite T [Ekstedt *et al.* - 2405.18349]
 - EWPT in dim-6 SMEFT [Camargo-Molina *et al.* - 2103.14022]
 - EWPT in dim-6 XSM [Oikonomou *et al.* - 2403.01591]
 - EWPT in Σ SM [Niemi *et al.* - 1802.10500]
- Tools for PT-related computations:
 - *CosmoTransitions* [Wainwright - 1109.4189]
 - *FindBounce* [Guada *et al.* - 2002.00881]
 - *DRalgo* [Ekstedt *et al.* - 2205.08815]
 - *BubbleDet* [Ekstedt *et al.* - 2308.15652]
- Theoretical uncertainties persist from the particle physics side:
 - LISA [Caprini *et al.* - 1910.13125]

Effective operators with
derivatives **not included**



3D EFT approach

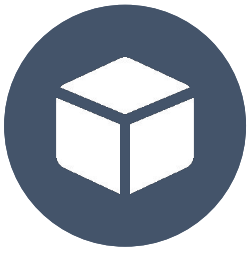
Building our high-T EFT

[Kajantie *et al.* - 9508379]
[Criado - 1901.03501]

We find an off-shell basis of operators up to (4D) dim-8 with the help of *BasisGen* and match at 1-loop (fermion only loops)¹:

$$\begin{aligned}\mathcal{L}_3 = & \frac{1}{2}K_3(\partial\varphi)^2 + \frac{1}{2}m_3^2\varphi^2 + \kappa_3\varphi^3 + \lambda_3\varphi^4 \\ & + \alpha_{61}\varphi^6 + \beta_{61}\partial^2\varphi\partial^2\varphi + \beta_{62}\varphi^3\partial^2\varphi \\ & + \alpha_{81}\varphi^8 + \alpha_{82}\varphi^2\partial_\mu\partial_\nu\varphi\partial^\mu\partial^\nu\varphi + \beta_{81}\varphi\partial^6\varphi + \beta_{82}\varphi^3\partial^4\varphi + \beta_{83}\varphi^2\partial^2\varphi\partial^2\varphi + \beta_{84}\varphi^5\partial^2\varphi\end{aligned}$$

¹ Scalar loop contributions to effective operators are subleading (see Appendix A in [2406.02667]).



3D EFT approach

Building our high-T EFT

[Kajantie *et al.* - 9508379]

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$$+ \alpha_{61}\varphi^6 + \beta_{61}\partial^2\varphi\partial^2\varphi + \beta_{62}\varphi^3\partial^2\varphi$$

$$+ \alpha_{81}\varphi^8 + \alpha_{82}\varphi^2\partial_\mu\partial_\nu\varphi\partial^\mu\partial^\nu\varphi + \beta_{81}\varphi\partial^6\varphi + \beta_{82}\varphi^3\partial^4\varphi + \beta_{83}\varphi^2\partial^2\varphi\partial^2\varphi + \beta_{84}\varphi^5\partial^2\varphi$$

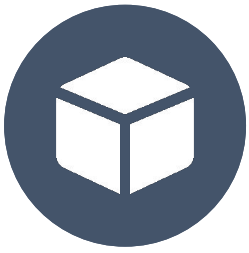
dim-6

Measure eff. ops. relevance

dim-8

Control EFT validity

¹ Scalar loop contributions to effective operators are subleading (see Appendix A in [2406.02667]).



3D EFT approach

Computing relevant PT magnitudes

Eff. ops. depend solely on the Yukawa (g)



Fix all other UV parameters in scan

[Chala *et al.* - 1905.00911]

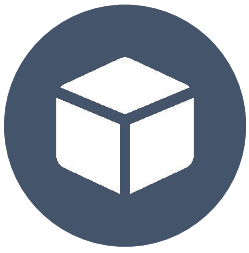
&

Study all PT magnitudes as functions of the **Yukawa**

m^2, κ, λ

$T_*, \alpha, \beta / H_*, v_w$

Bounce solution must be computed perturbatively!



3D EFT approach

Matching

[Chala, Criado, LG, López Miras -
2406.02667]

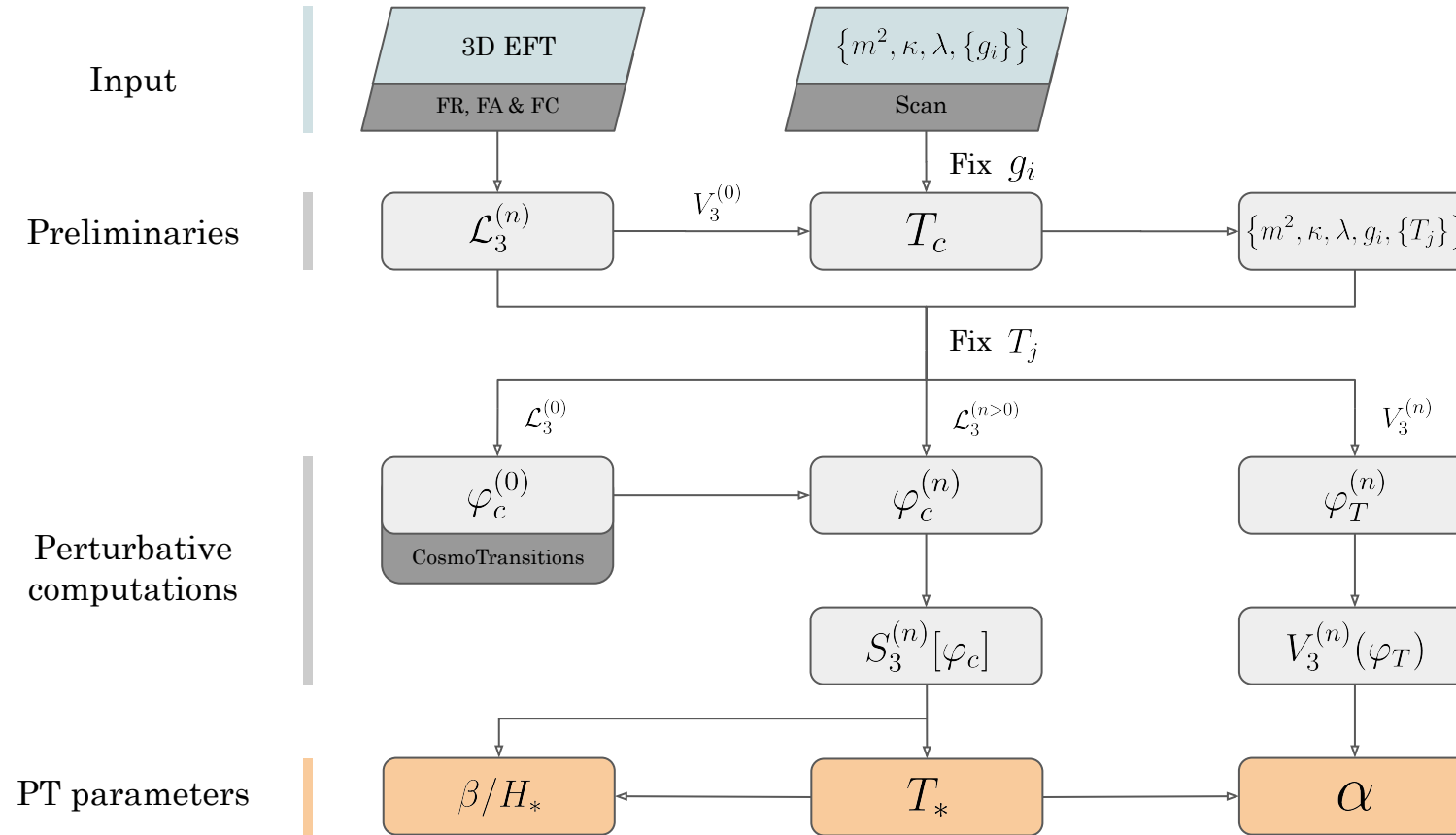
$$\begin{aligned} K_3 &= 1 + \frac{g^2}{12\pi^2}, & m_3^2 &= m^2 + \frac{g^2 T^2}{6}, & \kappa_3 &= \kappa\sqrt{T}, & \lambda_3 &= \lambda T; \\ \alpha_{61} &= -\frac{7\zeta(3)g^6}{192\pi^4}, & \beta_{61} &= -\frac{7\zeta(3)g^2}{384\pi^4 T^2}, & \beta_{62} &= \frac{35\zeta(3)g^4}{576\pi^4 T}; \\ \alpha_{81} &= \frac{31\zeta(5)g^8}{2048\pi^6 T}, & \alpha_{82} &= -\frac{31\zeta(5)g^4}{10240\pi^6 T^3}, & \beta_{81} &= -\frac{31\zeta(5)g^2}{10240\pi^6 T^4}, \\ \beta_{82} &= \frac{217\zeta(5)g^4}{20480\pi^6 T^3}, & \beta_{83} &= \frac{279\zeta(5)g^4}{20480\pi^6 T^3}, & \beta_{84} &= -\frac{217\zeta(5)g^6}{5120\pi^6 T^2}. \end{aligned}$$



3D EFT approach

The algorithm

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Results

GW power spectra and matching scale dependence

